## Magnetic Anisotropy in Magnetic van der Waals FePS ${ }_{3}$

Nauman, M. (Kyungpook Nat. Univ., Physics), Park, J.-G. (Seoul Nat. Univ., Physics), Kang, W. (Ewha Univ., Physics), Jo, Y. (Kyungpook Nat. Univ., Physics)

## Introduction

Layered transition metal tri-chalcogenides $\left(\mathrm{TMPS}_{3}, \mathrm{TM}=\mathrm{Fe}, \mathrm{Co}, \mathrm{Mn}\right.$, and Ni) represents one of the known layered systems in which both magnetic and crystallographic lattices are 2D. The layers are separated by a vdW gap and represent an antiferromagnetic (AFM) order with a transition temperature in the range of $80-150 \mathrm{~K}$ [1, 2]. $\mathrm{FePS}_{3}$ has a monoclinic honeycomb crystal structure with $C 2 / m$ structural space group and performs an Ising-type AFM alignment along the $c$-axis below $T_{N}=118 \mathrm{~K}$. Magnetic anisotropy plays an important role in the understanding of magnetic vdW materials, which exhibit magnetocrystalline anisotropy owing to the layered structure, and a reduced symmetry. We investigated magnetic properties and magnetic anisotropy for both in-plane and out-of-plane directions using torque magnetic measurements on $\mathrm{FePS}_{3}$ single crystals.

## Experimental

$\mathrm{FePS}_{3}$ single crystal was mounted on a piezoelectric resistance cantilever. We measured the angle-dependent torque $\tau(\theta)$ and magnetic field dependent torque $\tau(H)$. We used 18/20 T superconducting magnet (SCM-2). The angular position of the sample was controlled via a rotator, with respect to the applied magnetic field.

## Results and Discussion

The in-plane angle dependence of torque, $\tau(\Phi)$, at different applied magnetic fields at $T=50 \mathrm{~K}$ as shown in Fig. 1(a), indicates a perfect $\sin 2 \Phi$ pattern (red lines) that represents the isotropic behavior of $\mathrm{FePS}_{3}$ along the $a-b$ plane. Fig. 1(b) shows the out-of-plane angle dependent torque measurement $\tau(\theta)$. This curve shows a perfect $\sin 2 \theta$ pattern at low magnetic field, but at 4T or higher it deviates from $\sin 2 \theta$, the positive torque becomes sharper, and the negative torque becomes flatter. The amplitudes, as a function of the applied magnetic field, obtained from simple sinusoidal fitting of the $\tau(\Phi)$ and $\tau(\theta)$ curves are shown in Fig. 1(d). Larger amplitudes in out-of-plane demonstrate the larger anisotropy of $\mathrm{FePS}_{3}$.
 The amplitudes are evaluated using the power-law method $A \propto H^{\alpha}$ where $\alpha=1.99$ for the in-plane and $\alpha=$ 1.89 for the out-of-plane configurations. As shown in Fig. 1(c), out-of-plane $\tau(\theta)$ curves at different temperatures at 5 T show consistent behavior from 5 K to 100 K . The amplitude slightly decreases with increasing temperature. However, the amplitude sharply decreases at 110 K , and the phase is reversed at 120 K above TN and the amplitude again increases. It is noted that a sign change in the $\tau(\theta)$ curves with increasing temperature is caused by the magnetic transition from the antiferromagnetic to paramagnetic state.

Fig. 1 (a) In-plane and (b) out-of-plane angle dependent torque measurements of $\mathrm{FePS}_{3}$, (c) angle dependent torque at different temperatures, (d) Amplitudes of torque vs. magnetic fields for in-plane and out-of-plane rotations.

## Conclusions

A pronounced and explicit difference between out-of-plane and in-plane torque signal was observed that pertains to large anisotropy along these two directions. All the results suggest an imperfect AFM ordering along the c-axis with difference in net magnetization along $+c$ and $-c$ direction. A mixture of Zeeman energy, spin-orbit coupling and single ion anisotropy due to trigonal distortion of $\mathrm{FeS}_{6}$ octahedra contribute to the overall anisotropy of $\mathrm{FePS}_{3}$.

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## References

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