

## Nematic Electronic State and Magnetoelastic Quantum Oscillations in CeRhIn<sub>5</sub>

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### Introduction

CeRhIn<sub>5</sub> crystallizes in the quasi-two-dimensional tetragonal structure and displays antiferromagnetic (AF) ordering at  $T_N \approx 3.8$  K. The evolution of the ground states as a function of applied pressure, including a pressure-induced superconducting transition at 2.1 K, is prototypical of heavy-fermion superconductors and its high  $T_c$  is attributed to the quasi-2D crystal structure [1]. Quantum oscillations in the magnetization [1] and transport [2], observed for  $H//[100]$  and  $H//[001]$  at  $T < T_N$  in magnetic fields to 50T, reveal additional details including an anisotropic Fermi surface and an apparent phase transition to an electronic nematic state of matter at  $H^* \approx 30$  T [2,3].

### Experimental

Fiber Bragg grating (FBG) magneto-strain experiments [4] were carried on *c*-axis oriented single crystals of CeRhIn<sub>5</sub>, in 35T DC and 45T hybrid magnets at the NHMFL DC field facility, in the temperature range 0.35 K to 40 K.

### Results and Discussion

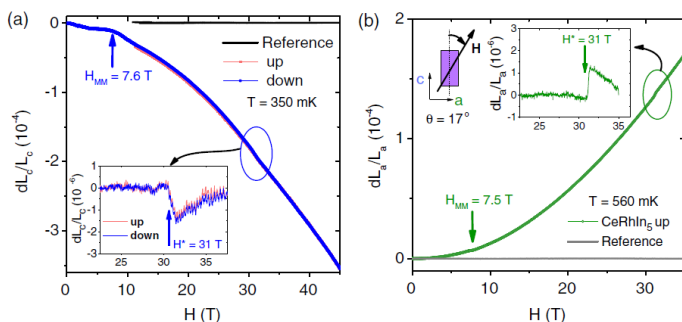
We measured the low temperature thermal expansion and magnetostriction of CeRhIn<sub>5</sub> to magnetic fields exceeding 40T. The axial and transverse magnetostriction, after subtraction of a smooth quadratic in field background, show clear anomalies near 28T (Fig1.). Magnetostriction quantum oscillations confirm a Fermi surface change at  $H^*$  with the emergence of new orbits. By analyzing the field-induced change in the crystal-field ground state, we conclude that the in-plane Ceb4f hybridization is enhanced at  $H^*$ , in agreement with the in-plane lattice expansion. We argue that the nematic behavior observed in this prototypical heavy-fermion material is of electronic origin, and is driven by the hybridization between 4f and conduction electrons which carries the f-electron anisotropy to the Fermi surface.

### Acknowledgements

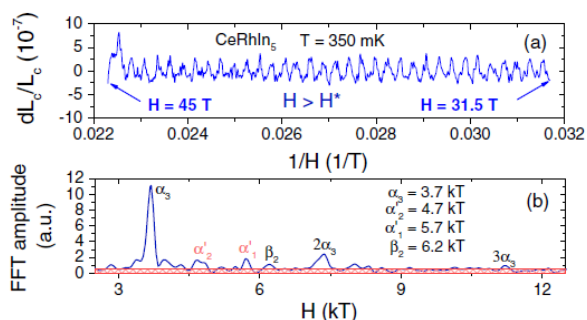
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### References

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**Fig.1** Magnetostriction of CeRhIn<sub>5</sub> along (a) the *c*-axis at  $T = 0.35$  K and (b) the *a* axis at  $T = 0.56$  K for fields applied  $\approx 20^\circ$  off the *c*-axis. Here  $\theta$  is the polar angle between the applied field and the *c*-axis. The azimuthal angle is set to  $\phi = 90^\circ$ . Insets show the data after a background subtraction.



**Fig.2** (a) Magnetostriction of CeRhIn<sub>5</sub> at 350 mK as a function of inverse field for fields applied  $\approx 20^\circ$  off the *c* axis. A high-pass filter was used to remove low-frequency oscillations that likely originate from the background difference below and above  $H^*$ . (b) FFT spectra in the region  $31.5 < H < 45$  T. The dashed area is an estimate of the noise floor.