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Benchmarking transcranial electrical stimulation finite element models: a comparison study

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Abstract

Objective. To compare field measure differences in simulations of transcranial electrical stimulation (tES) generated by variations in finite element (FE) models due to boundary condition specification, use of tissue compartment smoothing filters, and use of free or structured tetrahedral meshes based on magnetic resonance imaging (MRI) data. **Approach.** A structural MRI head volume was acquired at 1 mm³ resolution and segmented into ten tissue compartments. Predicted current densities and electric fields were computed in segmented models using modeling pipelines involving either an in-house (block) or a commercial platform commonly used in previous FE tES studies involving smoothed compartments and free meshing procedures (smooth). The same boundary conditions were used for both block and smooth pipelines. Differences caused by varying boundary conditions were examined using a simple geometry. Percentage differences of median current density values in five cortical structures were compared between the two pipelines for three electrode montages (F3-right supraorbital, T7-T8 and Cz-Oz). **Main results.** Use of boundary conditions commonly used in previous tES FE studies produced asymmetric current density profiles in the simple geometry. In head models, median current density differences produced by the two pipelines, using the same boundary conditions, were up to 6% (isotropic) and 18% (anisotropic) in structures targeted by each montage. Tangential electric field measures calculated via either pipeline were within the range of values reported in the literature, when averaged over cortical surface patches. **Significance.** Apparently equivalent boundary settings may affect predicted current density outcomes and care must be taken in their specification. Smoothing FE model compartments may not be necessary, and directly translated, voxelated tissue boundaries at 1 mm³ resolution may be sufficient for use in tES FE studies, greatly reducing processing times. The findings here may be used to inform future current density modeling studies.

Keywords: tDCS, tES, tACS, finite element simulation, current density, electric field

(Some figures may appear in colour only in the online journal)

Introduction

Finite element (FE) simulations are frequently used to predict current density distributions inside the human head caused

by transcranial electrical stimulation (tES) therapies such as transcranial direct current stimulation (tDCS) or transcranial alternating current stimulation (tACS). Patient-specific FE modeling of tES typically uses structural and diffusion information gathered in MRI scans to generate realistic head models [1–4]. Use of high resolution scans (ca. 1 mm in-plane

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resolution), inclusion of large numbers of tissue types and white matter anisotropy have been considered important factors affecting modeling accuracy [5, 6]. Tissue conductivity values are typically assigned based on literature referenced values [7–10] or direct scaling of diffusion tensor information [11–13]. Segmented head models are then meshed into desired element sizes and types (hexahedral or tetrahedral elements). Commercial platforms [2, 14–16] or specialized software [5, 16–18] have typically been used to solve FE [19–21], finite difference or boundary element method [22] simulations.

Many FE predictions of fields in realistic human head models were initially developed for inverse EEG/MEG applications [13, 23–25]. In these cases, simulations were performed on volume meshes containing either hexahedral or tetrahedral elements. For hexahedral meshes, voxel-based MRI structural data were directly transformed into regular hexahedral meshes [13] or converted into geometry-adapted hexahedral meshes using isoparametric FE approaches [18, 24, 26]. Hexahedral models were determined in studies by Wagner *et al* [27] and Vorwerk *et al* [28] to produce accurate results, but cautioned that the method of generating the hexahedral mesh must be chosen carefully to avoid leakage or artificially closed compartments. Tetrahedral elements have been constructed by compartmentalizing individual voxel-based hexahedral elements into multiple tetrahedral elements [22] or into free meshes using constrained Delaunay tetrahedralization (CDT) approaches [25, 29]. To simulate EEG data, meshed head models and electrodes are typically assigned homogeneous (insulating) Neumann boundary conditions on the head surface, and zero potential at ground electrodes [23]. Internal sources mimicking an active cortex are approximated as dipolar current sources and voltage differences caused by current dipoles are calculated using the reciprocity principle [30]. These forward problems have often been implemented using in-house or freely available FE software such as NeuroFEM, SimBio [13, 23, 24, 31, 32], SCIRun [22, 33].

In tES FE modeling studies, while in-house or freely available FE software [10, 18, 34] have been used, commercial platforms have been much more widely employed than in EEG modeling. COMSOL [2, 4, 14–16, 35–45] has been the most commonly used commercial package used in tES modeling studies. FE software packages such as Abaqus [7, 46, 47], and ANSYS (Ansys Inc., PA, USA) [48] have also been employed. Companion software packages such as ScanFE Simpleware (Synopsys, CA, USA) [15, 16, 36–38, 41–43] or Mimics (Materialise, Belgium) [2, 4, 44, 45] have often been required to mesh the complex geometry of head models for use with COMSOL, Abaqus and ANSYS, because their output formats are compatible with these packages. ScanFE constructs tetrahedral volume meshes using a combination of Delaunay and Advancing front approaches. The conversion process from voxel based MRI volumes into tetrahedral elements that can be successfully imported into commercial software requires a series of smoothing operations to avoid meshing complications [2, 7]. Pre-processing steps are also required to eliminate gaps between subdomains [49].

Previous tES FE studies have used a variety of electrode boundary conditions. Conventional transcranial stimulation protocols use a single pair of electrodes, with one electrode

assigned as the anode and the other the cathode [50]. In tES clinical studies, electrical currents of 2 mA or below [50] are applied between the anode and cathode using a controlled constant current source. Multiple FE simulations of tES studies have been performed with boundary conditions assigned with positive normal current density at the anode, and setting the cathode boundary condition to ground voltage [7, 14, 16, 35–37, 39, 41–43, 46, 47, 51]. Other FE tES studies assigned voltage values at the anode and cathode such that the voltage difference across the electrodes was equivalent to the desired current magnitude [2, 4, 17, 18, 34]. Boundary conditions specifying equal and opposite total current magnitudes applied at the anode and the cathode, respectively, with a reference voltage within the models, have also been employed [10, 21].

Variations in tES FE simulation configurations, such as choices of software, volume meshing types and boundary condition assignments across FE tES studies may cause inconsistencies, rendering comparisons of individual study outcomes difficult. Therefore, there is a need to benchmark existing tDCS FE simulation studies to ensure consistency. In this paper, we present comparisons of results from tES FE models executed using two modeling pipelines. The first pipeline involved a pair of commonly used commercial platforms, ScanFE and COMSOL, while the second pipeline used in-house software [5, 10]. In particular, we investigated the effects of mesh type, segmented model smoothing performed before meshing, and use of different boundary condition settings on outcomes in complex human head geometry simulations. Quadratic basis functions were used in both FE formulations as this is the default used by COMSOL software. The specific outcomes modeled were percentage differences of median current density between block and four levels of smoothing, for both isotropic and anisotropic conductivity distributions. Different levels of smoothing were achieved by applying a recursive Gaussian filter multiple times. White matter tissue anisotropy information was obtained from the diffusion weighted imaging (DWI) scan of the subject and was included in both block and smoothed model pipelines. Three different electrode montages were employed: left frontal-right supraorbital (F3-RS), left-right temporal cortex (T7-T8) and midline central–midline occipital cortex (Cz-Oz), with the first named electrodes in each pair selected as the anode. Distributions formed using either formulation were compared within target, deep and peripheral brain structures. The intended stimulation target structures were the inferior frontal gyrus (IFG) for F3-RS [5, 52], the anterior superior temporal gyrus (ASTG) for T7-T8 [53] and the occipital lobe (OCC) for Cz-Oz [17]. While this study investigated head models derived from a single subject using one set of material properties, we consider our findings indicative of differences that may be introduced using a variety of model workflows. We anticipate that these results will be useful in informing future tES modeling studies.

Materials and methods

All head models were derived from a single subject T1-weighted MRI dataset. White matter conductivity tensor

information was calculated from DWI data collected from the subject in the same imaging session. Block models were constructed directly from the segmented head volumes. Four degrees of recursive Gaussian smoothing were employed in constructing smooth models, and these were compared to block model results for the three electrode montages. Details of modeling and simulation processes are described in the following subsections.

T1-weighted and DWI data acquisition

T1-weighted and high angular resolution diffusion weighted imaging (HARDI) MRI datasets from a healthy human subject were collected using a 3T Achieva Phillips MRI system (the McKnight Brain Institute, University of Florida). HARDI data sets were acquired in 64 at high b -value (1000 s mm^{-2}) and 6 directions at low b -value (100 s mm^{-2}) using a 2D multislice spin echo sequence (TE = 86.0 ms, TR = 9022.8 ms) with a matrix size 112×112 and a total of 70 sagittal slices, each slice being 2 mm thick. T1 data was acquired using a 3D turbo spin echo pulse sequence (TE = 3.69 ms, TR = 8.057 ms) with a matrix size of 256×256 and 0.9375 mm resolution, with 160 axial slices of 1 mm thickness.

Head model construction

Prior to segmentation, the T1-weighted image was resampled to $256 \times 256 \times 256$, 1 mm^3 isotropic resolution using FreeSurfer (Cambridge, MA). Resampled T1-weighted images were segmented into ten tissue types using a combination of manual and automatic segmentation processes. Automatic segmentation of white and gray matter was performed in FreeSurfer, while segmentation for bone, air, and skin were done in SPM (Wellcome Trust Centre for Neuroimaging, London, UK) using an improved tissue probability map developed at CABI [7]. Masks obtained from automatic segmentation software were then imported into ScanIP (Simpleware, Synopsys, Exeter, UK) and corrected. Brainstem, gray matter, and remaining tissue compartments (muscle, fat, eyes, blood and CSF) were completed manually in ScanIP with reference to an anatomical atlas [54]. Figure 1 summarizes the segmentation pipeline used to categorize a single head model into ten final tissue types: white matter, gray matter, CSF, bone, muscle, fat, blood, air and skin. Literature sourced conductivity values were assigned to these ten tissue types, based on available values reported below 1 kHz, as shown in table 1. While conductivities measured at lower frequencies (closer to 10 Hz) would be preferred, existing literature values are scant, and prone to measurement error [55].

The FSL FDT module was used to calculate diffusion tensor (DT) and fractional anisotropy (FA) maps from HARDI DWI data [60]. Anisotropic conductivity tensors were assigned to portions of white matter compartments that had FA values greater than 0.5. White matter conductivity tensors

were calculated from DT principal eigenvector components ($V1_x, V1_y, V1_z$) as described below.

An initial conductivity tensor, \mathbf{S}_w was re-oriented to \mathbf{S}_w^* by pre- and post-multiplying \mathbf{S}_w with a product of rotational matrices, \mathbf{R} and \mathbf{R}^T respectively, such that

$$\mathbf{S}_w^* = \mathbf{R} \mathbf{S}_w \mathbf{R}^T \quad (1)$$

where

$$\mathbf{S}_w = \begin{bmatrix} \sigma_l & 0 & 0 \\ 0 & \sigma_t & 0 \\ 0 & 0 & \sigma_t \end{bmatrix}; l = \text{longitudinal}, t = \text{transverse} \quad (2)$$

$$\text{and } \mathbf{R} = \mathbf{A}_z \mathbf{A}_y \mathbf{A}_x. \quad (3)$$

$\mathbf{A}_x, \mathbf{A}_y,$ and \mathbf{A}_z were rotation matrices about x, y and z axes, respectively, for angles α, β, γ calculated from normalized $V1_x, V1_y$ and $V1_z$ vectors ($V1_{\bar{x}}, V1_{\bar{y}}$ and $V1_{\bar{z}}$) such that

$$\mathbf{A}_x(\alpha) = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha \\ 0 & \sin \alpha & \cos \alpha \end{bmatrix}, \mathbf{A}_y(\beta) = \begin{bmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{bmatrix}, \quad (4)$$

$$\mathbf{A}_z(\gamma) = \begin{bmatrix} \cos \gamma & -\sin \gamma & 0 \\ \sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\alpha = \tan^{-1} \frac{V1_{\bar{z}}}{V1_{\bar{y}}}, \beta = \tan^{-1} \frac{V1_{\bar{z}}}{\sqrt{V1_{\bar{x}}^2 + V1_{\bar{y}}^2}}, \gamma = \tan^{-1} \frac{V1_{\bar{y}}}{V1_{\bar{x}}}. \quad (5)$$

The diffusion free ($S0$) reference image of DWI data was used to co-register diffusion and T1-weighted data, and all T1-weighted data were resampled to the same resolution as diffusion images before performing tensor calculations. A large white matter anisotropy ratio ($\sigma_l/\sigma_t = 10$) was used to illustrate maximal effects. Results from both anisotropic and isotropic white matter models were compared.

TDCS electrode montages

All model results were calculated for three different electrode montages (F3-RS, T7-T8 and Cz-Oz). Current was injected at the anode site, which was the first named electrode in each montage pair. Each electrode had 35 cm^2 surface area and 1 mm thickness. This electrode size was typical of conventional tES electrode sizes [50]. Figure 2 shows locations of the three electrode montages with respect to the head model.

Block model construction

The block model mesh was constructed directly from segmented T1-weighted data. However, hexahedral meshes were not used. Instead, each voxel was transformed into six tetrahedral mesh elements. Each of these tetrahedra was assigned the

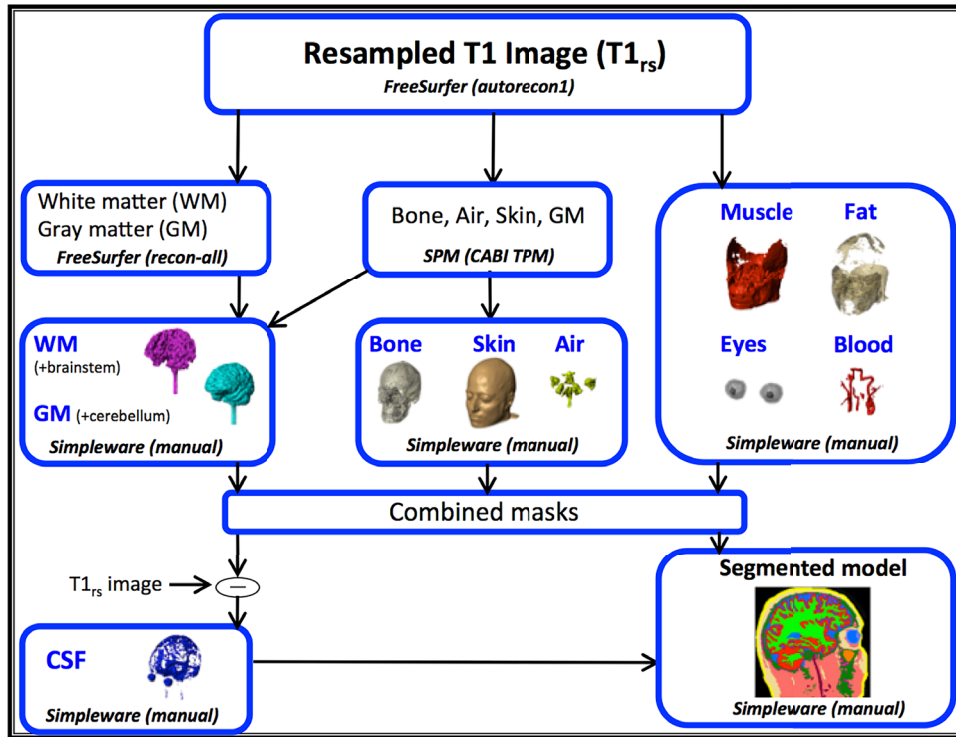


Figure 1. The segmentation pipeline involved both manual and automatic segmentation processes. The pipeline is shown in terms of the three major operation types (Freesurfer, SPM and manual). Results from each operation type were finally combined to produce a single head model containing ten tissue types.

Table 1. Literature-referenced conductivity values for ten tissue types for measurement frequency under 1kHz [10]. Bone conductivity (combined) was computed as $\sigma = \sqrt{\sigma_{can} \cdot \sigma_{cor}}$ where σ_{can} was cancellous bone and σ_{cor} was the average of cortical bone conductivity reported in the reference. Isotropic white matter conductivity was simulated using the formula $\sigma = \sqrt{\sigma_l \cdot \sigma_t}$ where σ_l was longitudinal and σ_t was transverse conductivity.

Tissue types	σ (S m ⁻¹)	Reference
Air	1.0×10^{-9} (Smooth pipeline)	—
Blood	6.7×10^{-1}	Geddes and Baker [56]
Bone	21.4×10^{-3} (cancellous) 5.52×10^{-3} (cortical) 10.9×10^{-3} (combined)	Akhtari et al [57]
Cerebrospinal fluid	1.8	Baumann et al [58]
Fat	2.5×10^{-2}	Gabriel et al [59]
Gray matter	1.0×10^{-1}	Gabriel et al [59]
Muscle	1.6×10^{-1}	Geddes and Baker [56]
Sclera, lens	5.0×10^{-1}	Gabriel et al [59]
Skin	4.3×10^{-1}	Holdefer et al [20]
White matter	1.2×10^{-1} (trans.) 1.2 (long.)	Geddes and Baker [56]

same conductivity as others within the same voxel. We used this configuration because we anticipated that at sufficiently high resolution (e.g. 1 mm³ voxels or below) this approach could be used to avoid any time and possible inaccuracies introduced by meshing steps, while not greatly increasing the size of the problem (each set of six quadratic basis function tetrahedra had 24 nodes, compared to the 20 nodes required to specify a quadratic basis function hexahedral element covering the same region).

Any tissue overlap in the block model was eliminated manually in MATLAB with the following tissue priority: white matter, gray matter, eyes, blood, air, CSF, fat, bone, muscle

and skin. The final block model had approximately four million tetrahedral elements. Air voxels were not included in finite element stiffness matrix assembly, and it was assumed that the slightly different air conductivity values in the two pipelines would not contribute to differences in modeling results.

Smooth model construction

All tissue masks in smooth models were overlapped to prevent gaps in the final mesh, a process referred to as mask solidifying. Smooth models were constructed by applying a recursive Gaussian smoothing filter to the segmented head

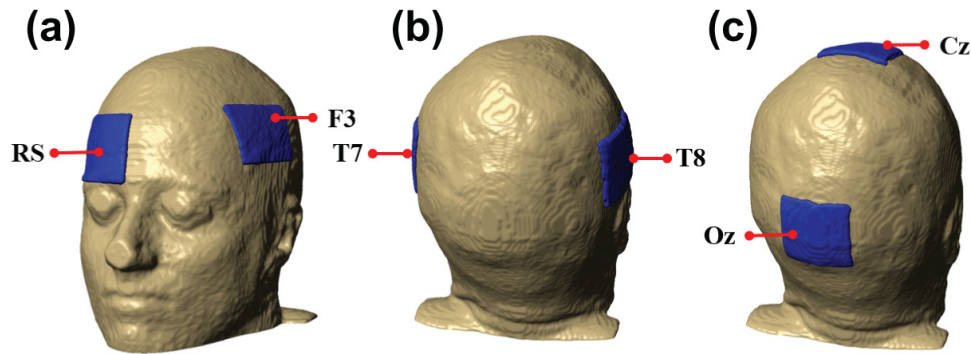


Figure 2. Electrode placements on head models used in this study. From left to right: (a) F3-RS, (b) T7-T8 and (c) Cz-Oz montages.

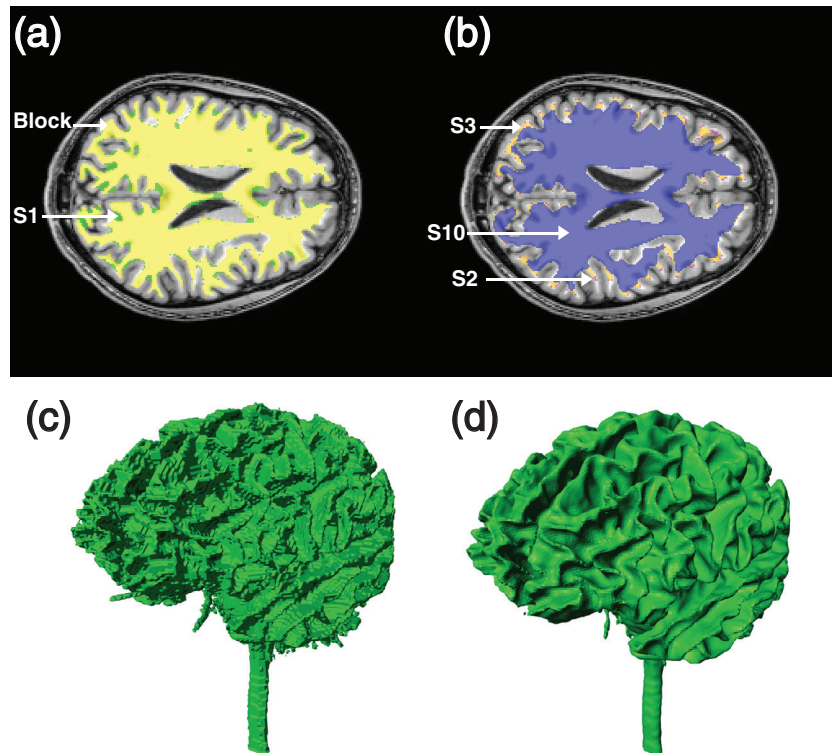


Figure 3. Illustration of differences between block and smoothed models in white matter compartment. (A) S1 white matter compartment (green) and segmented block (yellow) compartments overlaid on one slice of T1 model; (B) Overlay of S2 (purple), S3 (orange) and S10 (blue) models of white matter compartment on T1-weighted model. Differences between white matter compartments are shown in as a volume for (C) block and (D) smoothed models.

model. During this process, individual pixels within the image were multiplied by a Gaussian kernel function and original pixels were replaced by a weighted averaged of the result. The width of the Gaussian kernel was specified as 1 mm in x , y and z directions, to match the underlying data resolution. The smoothing filter only affected the outside boundary of individual tissue compartments. Additional pre-processing was performed on white matter masks in smooth models by applying morphological close and cavity fill operations in ScanIP prior to smoothing, to preserve thin structures [7]. Smooth models were meshed using the Simpleware ScanFE module after all masks were solidified, using the same tissue prioritization as used for the block model. In ScanFE, free tetrahedral mesh type (+FE Free) was selected with compound coarseness of -25 (maximum edge length of 2–5 mm). The

meshing process started by forming a surface mesh of triangular elements then replacing the triangular with tetrahedral elements by using a combination of Delaunay and Advancing Front approaches.

Four levels of smoothing were tested, resulting in four comparisons of block versus smooth models. The first smoothed model was constructed by applying the ScanIP recursive Gaussian smoothing filter to each of the ten tissue masks over a 1 pixel neighborhood in x , y , and z directions, to form model S1. Subsequent smoothed models were obtained by applying the same filter multiple times, to form S2, S3 and S10 i.e. the second level of smoothing (S2) was achieved by applying the smoothing filter twice, the third (S3) by applying the smoothing filter three times, and tenth level (S10) by smoothing ten times. Figure 3 illustrates effects of applying

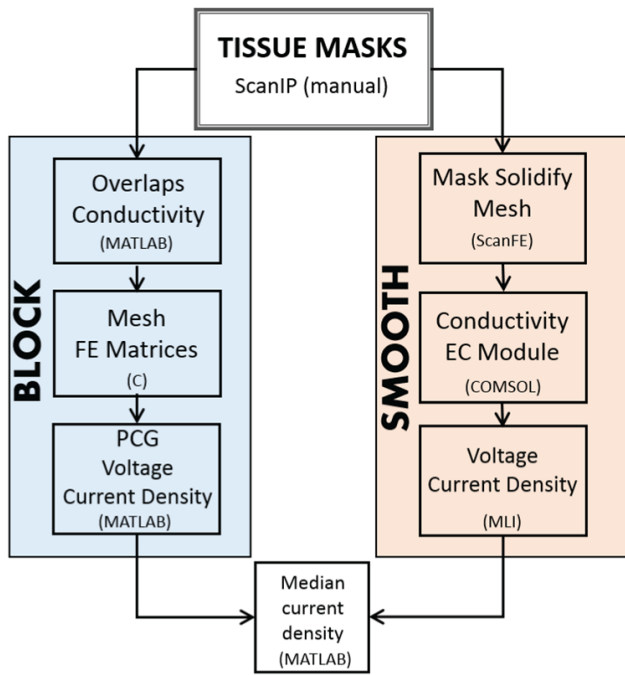


Figure 4. Simulation pipelines for block (left) and smooth (right) head model construction and finite element computation. Block models were processed using C and MATLAB, while smooth models were meshed using ScanFE and solved using COMSOL-MLI. Results from both pipelines were analyzed in MATLAB.

the recursive Gaussian smoothing filter on cross sections of white matter structures. Note that the smoothing filters distorted original boundaries. Smooth model meshes contained approximately 4.5 million tetrahedral elements. Finally, each meshed volume of the smoothed models was exported in COMSOL format. Only results from smoothed models S1, S2, S3 and S10 were used in this study. Figures 3(C) and (D) illustrates the different geometric features of block and smooth models (model S1) at the white matter surface.

Finite element modeling

FE formulations, using quadratic basis functions, were used to solve the Laplace equation inside the head volume with mixed boundary conditions applied on head or electrode surfaces. Block models were meshed and simulated using in-house C software and MATLAB, while smooth model meshes were exported to COMSOL as tetrahedral meshes. A total of 1 mA current magnitude was injected at each anode site in all models. Element numbers, solver types and solver tolerances were matched as closely as possible between block and smooth pipelines. The model processing workflows are summarized in figure 4, and are described below.

Block pipeline FE models. In block models, the stiffness and boundary condition matrices for the block-based tetrahedral mesh were formulated with a Continuous Galerkin finite element framework in C code using quadratic shape functions [61]. For anisotropic models, the three components of V1 in DTI data were exported as three volumes matched to the segmented volume, and anisotropic conductivity tensors were

computed as each element of the global FE stiffness matrix was assembled. Normal current densities were specified at both anode and cathode, such that the required total current was obtained, and a zero-volt reference node was placed near the model center. Voltage solutions for the block formulation were then calculated using the preconditioned conjugate gradient (`pcg`) command in MATLAB and a solution tolerance of 10^{-6} . Current densities, \mathbf{J} were calculated from voltage gradients, $\nabla\phi$, and voxelwise conductivity tensors, \mathbf{D}_w such that

$$\mathbf{J} = -\mathbf{D}_w \nabla\phi. \quad (6)$$

Smooth pipeline FE models. FE simulations for smooth models were performed using the electric current (EC) module with the default quadratic element type, and analyzed using the COMSOL-MATLAB livelink interface (MLI). For anisotropic models, the six unique entries in each white matter \mathbf{S}_w^* matrix were exported to individual volumes matched to the imported mesh, and white matter compartment conductivities in the mesh were specified via an interpolation function based on these components. Equal and opposite total current values were specified on electrode surfaces, with positive values assigned at the anode and negative values at the cathode. A reference voltage was assigned at one internal node of the model. A stationary iterative (conjugate gradient) solver was then used to solve Laplace's equation over the domain using a relative tolerance of 10^{-6} . Voltage and current density results were exported to MATLAB using the function `mphinterp` over a 256×256 (in plane) $\times 216$ (slice) volume with 1 mm^3 isotropic voxel resolution.

Model parameter and focus cortical structures

Median current density values within selected cortical structures were compared to quantify differences between block and smooth model results. Median values were used because we anticipated that block models may be affected by partial volume effects, and medians would be less affected by outlier current density values. Current density distributions were evaluated in five cortical structures. Three structures were target structures for each montage, namely the inferior frontal gyrus (IFR) for F3-RS, anterior superior temporal gyrus (ASTG) for T7-T8 and occipital lobe (OCC) for Cz-Oz configuration. The final two structures: hippocampus (HIP) and pre-central gyrus (PRC) were selected as representative deep and peripheral cortical structures, respectively.

Model verifications and comparisons

A verification study was performed to confirm that block and COMSOL finite element calculations produced identical results, and to compare the results produced by block and smoothed pipeline in a very simple model where the only differences were due to surface voxelation and boundary conditions. This cross-platform comparison was performed using two models. The first confirmatory model (C1) consisted of a $10 \times 10 \times 10 \text{ cm}^3$ box. The second model (C2) was based on the first model with a 5 cm diameter sphere placed at its

center. In both cases, boundary conditions assigning anode and cathode electrodes were applied to opposite faces of the box. Prior to comparisons, a mesh refinement study was performed for each model type to verify that each pipeline's results did not vary with respect to element size.

Both models (C1 and C2) were initially constructed using the COMSOL drawing interface. Model C1 had a uniform conductivity of 1.8 S m^{-1} . In model C2 the sphere had a conductivity of 0.01 S m^{-1} . The models were assigned two types of boundary conditions: normal current densities of an equal and opposite value assigned at the anode and cathode (BC-1) with a zero-voltage reference node placed in the center of the model, or, a normal current density equivalent (normal current density times face area) of 1 mA specified at the anode and zero-voltage (ground) at the cathode (BC-2).

The conductivity distribution of both models C1 and C2 was then exported from COMSOL into MATLAB on a uniform 1 mm^3 grid and used to compute solutions for each model using the block pipeline. The 1 mm stencil used in block models produced 6 million tetrahedral elements. A boundary condition specifying 1 mA current flow through the electrodes was applied in each case. The voltage drop computed across the electrode faces was examined to determine how well calculations agreed between the two approaches for each model. COMSOL results for model C1 were calculated using the surface average derived value output, while C-code average voltages were obtained by averaging voltages computed on electrode nodes on each face. Model C2 results were also compared in terms of voltage distributions, current density profiles in a central slice, and current density distributions within the sphere and surrounding regions.

Comparison of block and smoothed pipeline results

Comparisons of block and smoothed pipeline results for head models were performed as follows. Medians of simulated current density values (J_{median}) were normalized against adjusted resistance values (ΔR_{adj}) prior to comparisons [8], and percentage differences (PD) between median normalized current density values were computed as

$$J_{\text{norm,median}} = \frac{J_{\text{median}}}{\Delta R_{\text{adj}}} \quad (7)$$

$$\text{PD} = \left(\frac{J_{\text{norm,median}}^{\text{block}} - J_{\text{norm,median}}^{\text{smooth}}}{J_{\text{norm,median}}^{\text{block}}} \right) \times 100\%. \quad (8)$$

PDs were calculated within the five focus cortical volumes to assess differences between model results introduced by the different processing pipelines.

Electric field calculation

Tangential electric field components were calculated on the left anterior temporal gyrus of both block and smooth models and the isotropic, T7-T8 montage case, following [3]. In block

models, electric fields \mathbf{E} were computed from samples of local voltages, ϕ , such that

$$\mathbf{E} = -\nabla\phi. \quad (9)$$

Eight small regions (patches) with an area of $\sim 50 \text{ mm}^2$ each were isolated on the surface of the left anterior temporal gyrus and analyzed individually. Normal components of these local electric fields, \mathbf{E}_n , were calculated in each patch as a product of \mathbf{E} and the averaged normal vectors calculated over the patch surface. Finally, the component of the electric field tangent to the surface patches, \mathbf{E}_{tan} , were computed as

$$\mathbf{E}_{\text{tan}} = \mathbf{E} - \mathbf{E}_n. \quad (10)$$

Median values for \mathbf{E}_{tan} magnitudes were calculated in each patch and averaged to obtain a single median value for block and smooth models. The averaged median values were then compared to similar experimentally measured values reported in the literature [3, 62].

In smooth models, electric fields for the entire model computed using BC-1 were exported on the same grid as used for block models using COMSOL-MLI, and the gridded S1 cortical surface was manually registered to the block cortical surface in ScanIP. The normal component of the electric field in S1 was then calculated on the cortical surface using the same method as used for block models, with \mathbf{E}_{tan} magnitudes also calculated using equation (10). Finally, \mathbf{E}_{tan} distributions in each patch were translated into histograms for comparison between the two pipelines.

Results

Model cross-comparisons

The simple box test calculation resulted in negligible differences between results generated by block and COMSOL BC-1 workflows, while results produced by COMSOL BC-2 showed a discrepancy in vertical current density profiles. Figure 5(a) shows the voltage distributions for model C1, while figure 5(b) shows the voltage distribution along a central slice of model C2. The color gradients in block models were discrete while in smooth models were more continuous, because COMSOL displays were interpolated in post-processing steps. The voltage drop across the electrodes in C1 models was found to be 5.55 mV for both block and COMSOL versions (values were less than 0.1% different), and as predicted by analytic calculations. Voltage solutions for the two C2 models were slightly offset due to the necessary differences in locations of zero volt reference points in the two models (figure 5(b)). Current density norm profiles for C2 models are shown in in figure 5(c). In model C2, vertical current density profiles (from anode to cathode) produced by COMSOL BC-2 were asymmetric while the same profiles for COMSOL BC-1 and block models were symmetric. Horizontal current density profiles showed a slight asymmetry in COMSOL and block C2 models. As expected, measures of central tendency were less different. Mean current density norms within the sphere ROI were at most 2% different (medians were at most 1%

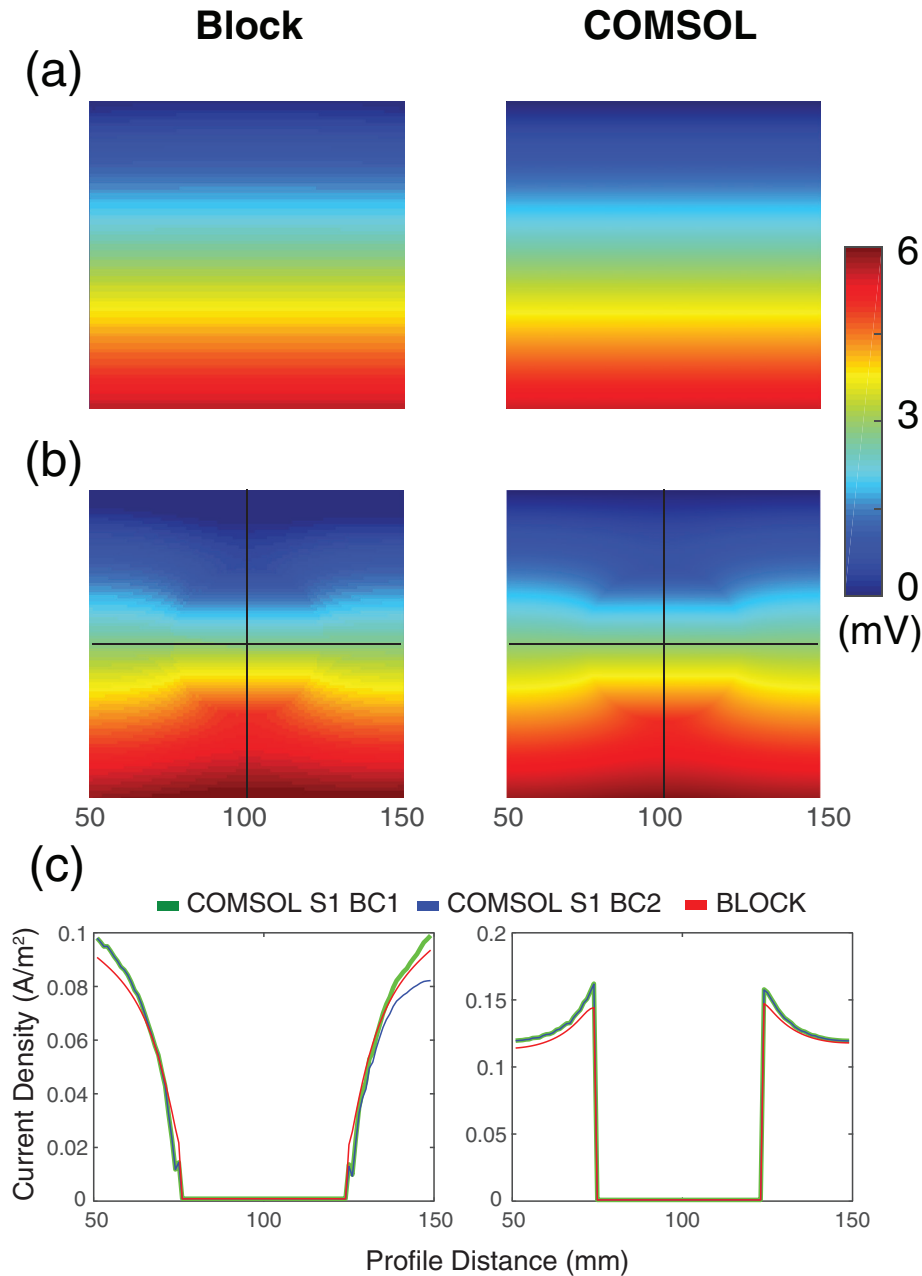


Figure 5. Cross sections of voltage distributions and current density profiles in two of confirmatory models. Block, and COMSOL BC-2 model voltage results for (a) box only model (C1) and (b) box with added sphere (C2). Block, COMSOL BC-1 and COMSOL BC-2 for (c) Vertical (left) and Horizontal (right) current density norm profiles for model C2 along profile lines marked in (b).

different), and at most 0.5% different (medians identical) in the bulk of model C2.

Tissue volumes and modeling time comparisons

PD in volumes and normalized median current densities were calculated between block head models and the four smoothed head models (S1, S2, S3 and S10), producing a total of 24 comparisons over the three electrode montages and white matter anisotropy conditions.

Table 2 shows the volumes for each of the ten tissues included in each model. Total volumes of block and smooth models differed by a maximum of around 2%. The largest difference between focus structure volumes was found to be 22%.

The solidifying steps involved in constructing smoothed models caused compartments to be differently shaped compared to original T1-weighted images. Solidification steps produced large increases in fat and gray matter compartment volumes (58% and 20% respectively), and decreases in bone and CSF (−16% and −35% respectively). While white matter volume did not change greatly, the shape of this compartment was markedly different after multiple smoothing steps (figure M3).

Table 3 shows volumes of focus structures as well as percentage differences between smoothed and block models. While overall gray matter volumes were larger in S1 than for the original block model, solidification and meshing for COMSOL steps resulted in focus structures being smaller by 10%–20%.

Table 2. Volumes (cm³) of each tissue type within the segmented head model for block, S1, S2, S3 and S10 models. Sums of individual tissue volumes are shown in the last row.

Structure	Block volume (cm ³)	S1 volume (cm ³)	S2 volume (cm ³)	S3 volume (cm ³)	S10 volume (cm ³)
White matter	512	554	544	536	497
Gray matter	647	777	786	794	833
Eyes	12	12	11	11	10
Air	68	66	63	62	54
Blood	14	12	11	9	3
CSF	324	210	210	211	213
Fat	218	346	328	315	259
Bone	793	665	654	666	708
Muscle	888	865	869	873	900
Skin	635	635	628	629	629
Total	4111	4140	4105	4106	4105

Table 3. Focus structure volumes for block and S1 model, including percentage differences.

Structure	Block (cm ³)	S1 (cm ³)	Difference (%)
ASTG	7.3	6.2	-15
HIP	7.7	6.9	-10
IFR	12.5	10.4	-17
OCC	41.1	32.0	-22
PRC	28.7	23.9	-17

Computed modeling times were the sum of times required for post-segmentation model processing, meshing and finite element solution and were approximately 200 min and 400 min for isotropic block and smooth models, respectively. In isotropic models, computed modeling times were 40 min for block and 120 min for smooth pipelines. Block models were found to require half the time of smooth models to solve for anisotropic cases and a third of the time for isotropic cases.

Median current density comparisons

Cross-sectional images showing current density values in a single slice of block and smooth models are shown in figure 6. Normalized median current density PD values between block and S1 models in cortical structures for the three electrode montages are summarized in figure 7. PD values in comparisons of anisotropic models were generally larger than for isotropic cases for T7-T8 and F3-RS montages, and smaller than isotropic cases for the Cz-Oz montage. The largest absolute PD values over all three electrode montages and structures were found to be 35.2% for anisotropic and 21.1% for isotropic cases. The corresponding smallest absolute PD values were 0.8% and 1.1% for anisotropic and isotropic cases, respectively. Median current density PD values in structures presumed targeted by each montage were at most around 10% (for the IFR structure with F3-RS).

Effects of additional smoothing

Figure 8 shows current density PD values in focus structures between block models and all smooth models, for each montage. The largest and smallest PD absolute values were

observed in S1 models, and were 35.2% and 0.8%, respectively. Models S2 and S3 showed very similar median current density PD values for almost all structures, as expected, and almost always overlapped. There was no overall clear trend between the current density PD values and smoothing level for all electrode montages and tissue anisotropy assignments. In some cases, PD values were smaller in focus structures in more smoothed models than for S1. As also shown in figure 7, the largest PD values in structures targeted by each montage were less than 6% in isotropic and 18% in anisotropic models.

Tangential electric fields in block and smooth models

Tangential electrical field distributions were calculated on each of the eight sample patches on the left anterior temporal gyrus for the T7-T8 montage. Histograms of tangential electric field magnitudes are shown in figure 9. Averaged median tangential electric field magnitudes were found to be 0.064 mV mm⁻¹ and 0.069 mV mm⁻¹ (7% different) for block and S1 models, respectively.

Discussion

Summary of findings

In this study, structure volumes, current density and electrical field were calculated to quantify the difference produced by block and smooth pipelines. Major differences between different pipeline results could have been caused by any combination of:

- (i) finite element solution differences
- (ii) boundary conditions
- (iii) differences in compartment voxelation, or
- (iv) compartment smoothing and solidification steps

In the sections below, we summarize findings for the different model types in terms of structure shape and volume, current densities and electric fields, first for the cross-platform validation models and then for head models. We assumed that if different pipelines produced differences less than 20% this was an acceptable correspondence.

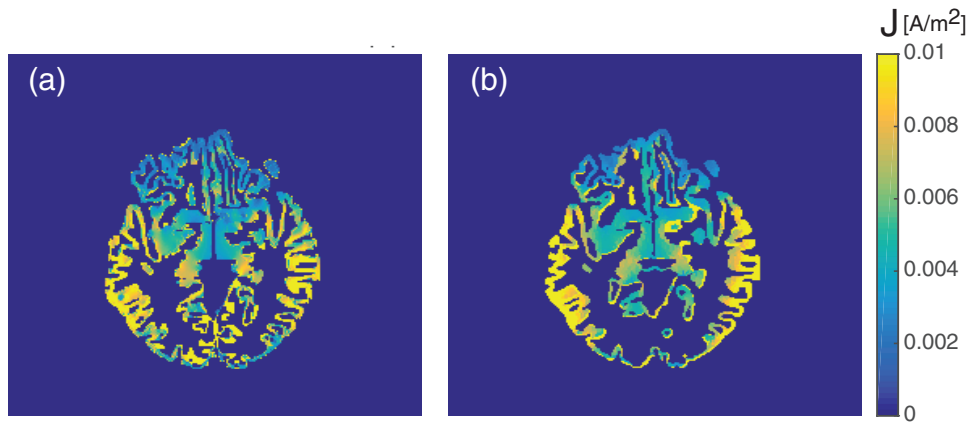


Figure 6. Axial cross-sections of predicted current density distributions in the cortex. Predicted current densities for T7-T8 montage are masked in gray matter structures for isotropic (A) block and (B) S1 models.

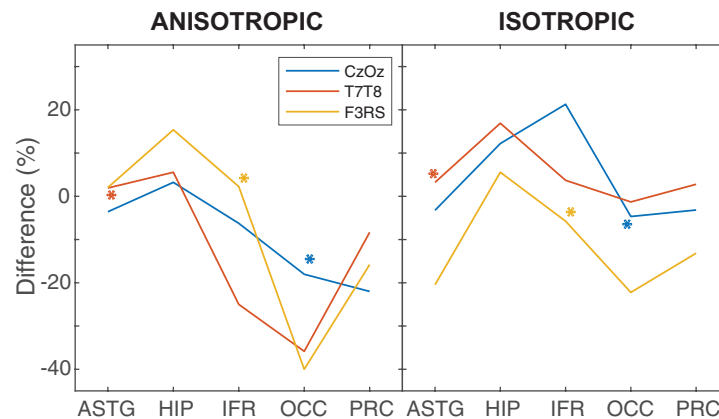


Figure 7. Normalized median current density percentage difference (PD) values in five focus structures for block and S1 BC-1 models. Negative values indicated that normalized current density values in smooth models were larger than those in block models. Presumed target structures for each montage are indicated with asterisks.

Cross-platform validation comparisons

The first cross-platform validation step (C1) showed there were no differences between finite element solutions, since differences were less than 0.1%, and both solutions agreed with analytical values. In the second validation model (C2) it was found that mean and median values were at most 2% or 1% respectively different, in the internal sphere ROI. This indicated that while voxelation of the sphere surface produced some differences, they were small. As expected, there were large (up to 70%) differences between current density measures at the sphere boundary. Current density profiles drawn from the anode to cathode (vertical) were asymmetric for the COMSOL BC-2 model, but the same profiles were symmetric in block and COMSOL BC-1 models. COMSOL BC-2 simulations assigned positive current density at the anode and set the cathode to ground voltage, a setting which is frequently used in tES FE studies [7, 14, 16, 35–37, 39, 41–43, 46, 47, 51]. However, these settings may result in asymmetry, as constant voltage boundary conditions do not necessitate constant normal current density distributions. This observation must be considered carefully in future studies, as it may have been assumed that these settings give the same results as total current or voltage only boundary conditions

without needing to specify potential at a separate point. Use of either equal and opposite normal current density (if compensated to give the correct total current), equal and opposite total current, or voltage only boundary conditions should produce results closer to actual experimental conditions. Specification of boundary conditions using total current or fixed voltage boundary conditions is recommended.

In summary, findings from the confirmation models reflected effects of differences between meshes, sphere voxelation and boundary condition specifications, as well as possible registration errors between the two methods.

Head model comparisons

A total of 24 comparisons of median current density PD values were investigated for four levels of tissue smoothing (S1, S2, S3, S10), three electrode configurations (Cz-Oz, T7-T8, F3-RS), and two types of white matter (isotropic or anisotropic). Relationships between model construction and solution times, structure volumes, calculated current densities, surface electric fields, and effects of additional smoothing levels are discussed relative to electrode placements and tissue anisotropy for both pipelines in the subsections below.

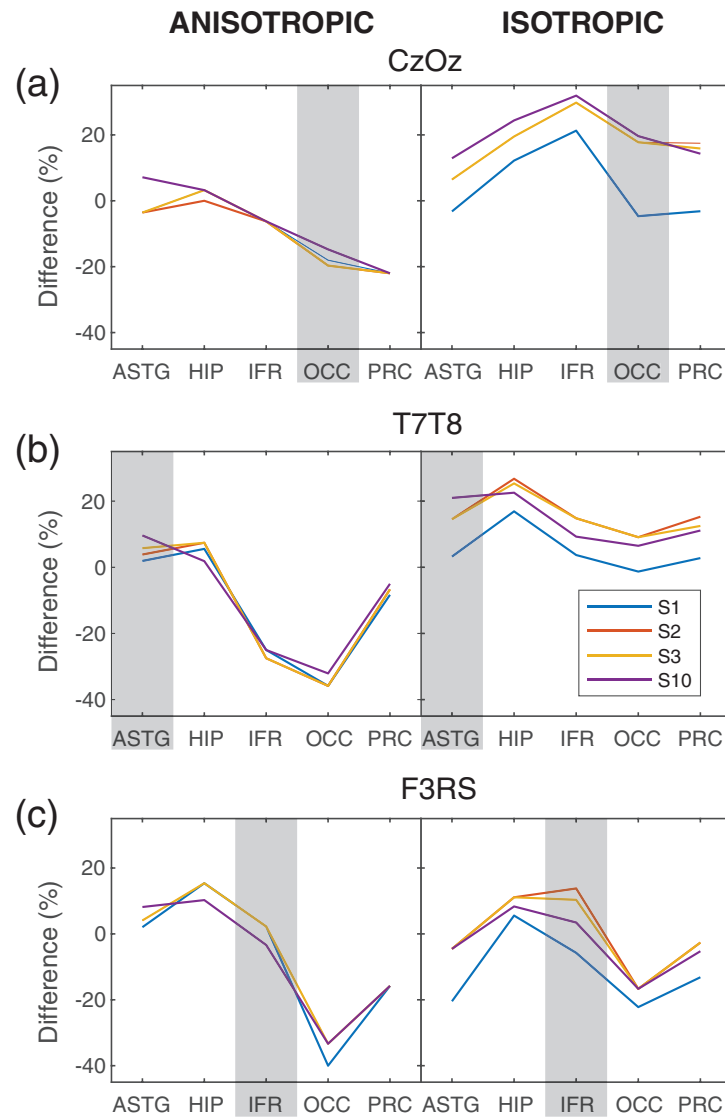


Figure 8. Normalized median current density percentage differences (PD) for block and smooth BC-1 (S1, S2, S3, S10) models. Plots show PD values obtained in (A) Cz-OZ (B) T7-T8 and (C) F3-RS electrode montages for (left) anisotropic and (right) isotropic cases. Negative values indicated that normalized current density values computed in smooth models were larger than in block models.

Overall model construction differences and times. Because of the need to perform smoothing and mask solidification for successful volume meshing in Simpleware for export to the COMSOL platform, the smooth model pipeline took much longer to solve. Since anisotropic smoothed models took three times as long as matching block versions, while isotropic smoothed models took twice as long, block models may have a distinct advantage if rapid modeling is required.

Smoothing distortion effects and cortical structure differences. Overall model volume was preserved after mask solidification and meshing. White matter volumes were also preserved upon smoothing. However, the white matter compartment shape was distorted. Other internal compartments changed volume markedly, with volume increases in fat and gray matter being balanced by decreases in bone and CSF volumes. This indicated that even though the smoothing neighborhood used was minimized, the effect of solidification and smoothing with multiple compartments could produce a

very different tissue distribution than represented in original MRI data. There was no distinct relationship between cortical structure size and volumetric changes on processing block models to smooth models. Instead, the complexity of individual cortical structures (e.g. folds and ridges on the surface of the cortex) and tissue prioritization choice was most likely the major contributor to volumetric differences between the modeling pipelines.

Current density distributions with different electrode montages. While PD values may have been affected by volume changes in target structures (between 10% and 20%), location of structures relative to electrodes was probably the most important determinant of current density PD values. In targets for each montage, median current density PD values for target structures were in the range of $\pm 6\%$ for isotropic and $\pm 18\%$ for anisotropic cases. This indicated that the two pipelines produced similar values for target structures. Comparable observations were obtained for model comparisons

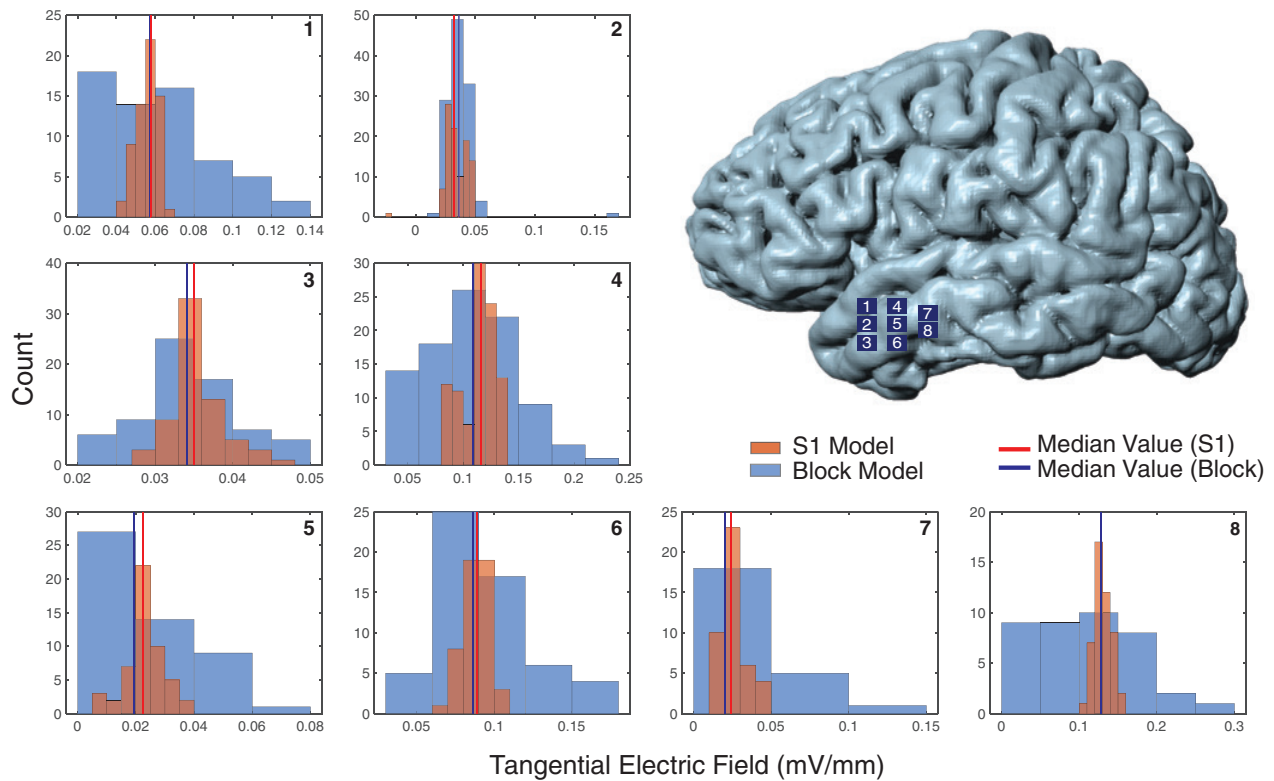


Figure 9. Histograms of tangential electric field (E_{tan}) magnitudes of distributions on the surface of the left anterior temporal gyrus for isotropic block and S1 BC-1 models using the T7-T8 montage. Locations of the eight surface patches are shown on the isosurface model (top right). Plots show E_{tan} magnitude distributions over patches 1–8. Blue lines denote median E_{tan} magnitude values in each patch for the block model while red lines show the same quantity for the S1 BC-1 model.

between block and additional smoothing levels. These findings suggest that the majority of brain regions surrounding target structures were least affected by meshing choice, most likely due to the higher current densities in brain regions nearby electrode locations.

Effects of additional smoothing. The three additional degrees of smoothing (S2, S3 and S10) did not affect changes of normalized median current densities in a predictable manner. For instance, normalized median current density PD values mostly increased from S2 to S10 in isotropic cases for T7-T8 and F3-RS, but decreased in isotropic cases for Cz-Oz. This was most likely because of shape changes produced by each additional smoothing step. Similar current density PD values were observed for S2 and S3 models throughout, as expected, because of the similarities in structure volumes and shapes (table 2).

Tissue anisotropy effects. The anisotropy ratio used here, $\sigma_l/\sigma_t = 10$, was large. We included this value to illustrate maximal effects. Results derived from diffusion tensor imaging studies indicate that white matter anisotropy may be modeled more appropriately with anisotropy ratios closer to 3 [13]. Inclusion of white matter anisotropy altered normalized current density PD values for some electrode configurations as shown in figure 8. Tissue anisotropy reduced PDs in focus structures in T7-T8 and F3-RS montages from $\pm 6\%$ to $\pm 3\%$. OCC current density PD values were larger for anisotropic cases than isotropic cases for the Cz-Oz montage. This

suggests that white matter fiber orientations in and or surrounding OCC contributed to the current flow patterns from Cz to Oz. A similarly large difference was observed for median current density PDs in the PRC for the Cz-Oz montage. The current density PD value was larger in the Cz-Oz anisotropic model than for the Cz-Oz isotropic model. In addition, based on table 3, OCC and PRC were the largest of the five focus structures. The size factor might also contribute to the largest changes in calculated current density PD between anisotropic and isotropic cases with respect to electrode placements. Therefore, a combination of white matter fiber orientation, the large anisotropy ratio used, the smoothed tissue boundary, structure size and electrode locations affected current density distributions in individual cortical structures and resulted in the 12% differences observed between isotropic and anisotropic cases.

Electric field comparisons. Averaged maximum and median local electrical field values on the cortical surface nearby stimulating electrodes calculated for both block and smooth models were within the range of local electric fields observed during TES reported in the literature [3, 62]. For instance, in the study reported by Opitz *et al* [3] the median calculated projected electric field along the brain surface near the anode (T7) ranged between 0.059 and 0.098 $mV\ mm^{-1}$ with a maximum value of 0.37 $mV\ mm^{-1}$. Huang *et al* [62, 63] reported a maximum projected electric field of 0.25 $mV\ mm^{-1}$ in the fronto-lateral location for an M1-SO configuration. These reported values were derived from electric potentials measured

by surface cortical electrode arrays during TES injections of 1 mA. Averaged tangential electric field measures were around 7% different between block and smoothed models. Averaging therefore reduced the differences between voxelated and smoothed results. Therefore, these comparisons served as a model validation for both pipelines, and a confirmation that either pipeline was suitable to predict TES field quantities, as long as surface measures were averaged.

Result accuracy

We can only consider differences between the two modeling pipelines here, not absolute accuracy, although we can compare smoothed compartments with source T1 data in volumetric comparisons. Model accuracy determinations can only be evaluated using independent measurements of electric fields and current densities *in vivo*. This may be possible using (tangential, surface) electric field measurements [3, 62, 63] or measures of internal current density distributions [64–66]. Since conductivities used in models are taken from excised tissue [66], it is also necessary to obtain estimates of *in vivo* conductivity values at frequencies typical of stimulation frequencies to perform full validations.

Finally, it is not clear if tissue conductivities may change during tDCS or tACS administration. In particular, it may be the case that properties of tissues immediately under electrodes may change due to presence of wetting electrolytic materials. This could be explored further using sensitivity studies to reconcile model conductivity against gold standard data obtained using cortical arrays or magnetic resonance phase imaging [9, 44, 62, 63].

Further observations

This study has demonstrated that block-based tetrahedral models provide acceptable results in simulations and that smoothing may not be necessary to obtain good simulation quality. The quadratic element order used here may not be necessary to obtain accurate solutions, particularly at high resolution (voxels smaller than 1 mm³) and linear models may be sufficient. Further study would be necessary to establish this. Differences between meshes may have been responsible for some measure of the differences between results, but this was not rigorously analyzed in this study. Future work may involve consideration of mesh density distribution effects when comparing the two approaches. Use of block-based tetrahedral models, possibly including adjustment of node locations to accommodate tissue features, as in [67–69] may overall improve model formation time, while avoiding leakage effects that may result from usage of hexahedral models [70].

Conclusion

Benchmarking tES FE studies is crucial to ensure accurate and consistent predictions of current density distribution in realistic head models of tES. Many factors need to be considered before performing FE tES studies. Care must be taken to ensure boundary conditions are consistent. For instance,

assigning normal current density at the anode and set the cathode to ground in a simple geometry may produce asymmetric current density profiles across the anode and cathode. Many tES FE studies in the present literature use commercial platforms requiring subdomain smoothing. However, smoothing steps require additional processing times and may distort original structural information. At 1 mm³ image resolution, FE models constructed from direct conversion of segmented MRI volumes may produce comparable averaged field distributions caused by tES, as demonstrated here. Therefore, modeling pipelines involving direct conversion from imaging voxels to FE models may be desirable to use in tES FE studies.

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References

- [1] Laakso I, Tanaka S, Koyama S, De Santis V and Hirata A 2015 Inter-subject variability in electric fields of motor cortical tDCS *Brain Stimul.* **8** 906–13
- [2] Miranda P C, Mekonnen A, Salvador R and Ruffini G 2013 The electric field in the cortex during transcranial current stimulation *NeuroImage* **70** 48–58
- [3] Opitz A *et al* 2016 Spatiotemporal structure of intracranial electric fields induced by transcranial electric stimulation in humans and nonhuman primates *Sci. Rep.* **6** 31236
- [4] Ruffini G, Fox M D, Ripolles O, Miranda P C and Pascual-Leone A 2014 Optimization of multifocal transcranial current stimulation for weighted cortical pattern targeting from realistic modeling of electric fields *NeuroImage* **89** 216–25
- [5] Sadleir R J, Vannorsdall T D, Schretlen D J and Gordon B 2012 Target optimization in transcranial direct current stimulation *Front. Psychiatry* **3** 90
- [6] Datta A, Zhou X, Su Y, Parra L C and Bikson M 2013 Validation of finite element model of transcranial electrical stimulation using scalp potentials: implications for clinical dose *J. Neural Eng.* **10** 036018
- [7] Huang Y, Dmochowski J P, Su Y, Datta A, Rorden C and Parra L C 2013 Automated MRI segmentation for individualized modeling of current flow in the human head *J. Neural Eng.* **10** 066004
- [8] Indahlastari A, Chauhan M, Schwartz B and Sadleir R J 2016 Changing head model extent affects finite element predictions of transcranial direct current stimulation distributions *J. Neural Eng.* **13** 066006
- [9] Parazzini M, Fiocchi S, Rossi E, Paglialonga A and Ravazzani P 2011 Transcranial direct current stimulation:

- estimation of the electric field and of the current density in an anatomical human head model *IEEE Trans. Biomed. Eng.* **58** 1773–80
- [10] Sadleir R J, Vannorsdall T D, Schretlen D J and Gordon B 2010 Transcranial direct current stimulation (tDCS) in a realistic head model *NeuroImage* **51** 1310–8
- [11] Tuch D S, Wedeen V J, Dale A M, George J S and Belliveau J W 2001 Conductivity tensor mapping of the human brain using diffusion tensor MRI *Proc. Natl Acad. Sci. USA* **98** 11697–701
- [12] Kwon O I, Sajib S Z, Sersa I, Oh T I, Jeong W C, Kim H J and Woo E J 2016 Current density imaging during transcranial direct current stimulation using DT-MRI and MREIT: algorithm development and numerical simulations *IEEE Trans. Biomed. Eng.* **63** 168–75
- [13] Rullmann M, Anwander A, Dannhauer M, Warfield S K, Duffy F H and Wolters C H 2009 EEG source analysis of epileptiform activity using a 1 mm anisotropic hexahedra finite element head model *NeuroImage* **44** 399–410
- [14] Datta A, Baker J M, Bikson M and Fridriksson J 2011 Individualized model predicts brain current flow during transcranial direct-current stimulation treatment in responsive stroke patient *Brain Stimul.* **4** 169–74
- [15] Shahid S S, Wen P and Ahfock T 2014 Assessment of electric field distribuion in anisotropic cortical and subcortical regions under the influence of tDCS *Bioelectromagnetics* **35** 41–57
- [16] Datta A, Truong D, Minhas P, Parra L C and Bikson M 2012 Inter-individual variation during transcranial direct current stimulation and normalization of dose using MRI-derived computational models *Front. Psychiatry* **3** 91
- [17] Neuling T, Wagner S, Wolters C H, Zaehle T and Herrmann C S 2012 Finite-element model predicts current density distribution for clinical applications of tDCS and tACS *Front. Psychiatry* **3** 83
- [18] Wagner S, Rampersad S M, Aydin U, Vorwerk J, Oostendorp T F, Neuling T, Herrmann C S, Stegeman D F and Wolters C H 2014 Investigation of tDCS volume conduction effects in a highly realistic head model *J. Neural Eng.* **11** 1–14
- [19] Wagner T, Fregni F, Fecteau S, Grodzinsky A, Zahn M and Pascual-Leone A 2007 Transcranial direct current stimulation: a computer-based human model study *NeuroImage* **35** 1113–24
- [20] Holdefer R N, Sadleir R and Russell M J 2006 Predicted current densities in the brain during transcranial electrical stimulation *Clin. Neurophysiol.* **117** 1388–97
- [21] Suh H S, Lee W H and Kim T S 2012 Influence of anisotropic conductivity in the skull and white matter on transcranial direct current stimulation via an anatomically realistic finite element head model *Phys. Med. Biol.* **57** 6961–80
- [22] Hyde D E, Dannhauer M, Warfield S K, MacLeod R and Brooks D H 2016 Evaluation of numerical techniques for solving the current injection problem in biological tissues *Proc. IEEE Int. Symp. Biomed. Imaging* pp 876–80
- [23] Wolters C H, Anwander A, Tricoche X, Weinstein D, Koch M A and MacLeod R S 2006 Influence of tissue conductivity anisotropy on EEG/MEG field and return current computation in a realistic head model: a simulation and visualization study using high-resolution finite element modeling *NeuroImage* **30** 813–26
- [24] Dannhauer M, Lanfer B, Wolters C H and Knosche T R 2011 Modeling of the human skull in EEG source analysis *Hum. Brain Mapp.* **32** 1383–99
- [25] Vorwerk J, Clerc M, Burger M and Wolters C H 2012 Comparison of boundary element and finite element approaches to the EEG forward problem *Biomed. Tech.* **57** 795–8
- [26] Lau S, Gullmar D, Flemming L, Grayden D B, Cook M J, Wolters C H and Hauelsen J 2016 Skull defects in finite element head models for source reconstruction from magnetoencephalography signals *Front. Neurosci.* **10** 141
- [27] Wagner S, Lucka F, Vorwerk J, Herrmann C S, Nolte G, Burger M and Wolters C H 2016 Using reciprocity for relating the simulation of transcranial current stimulation to the EEG forward problem *NeuroImage* **140** 163–73
- [28] Vorwerk J, Engwer C, Pursiainen S and Wolters C H 2017 A mixed finite element method to solve the EEG forward problem *IEEE Trans. Med. Imaging* **36** 930–41
- [29] Cho J H, Vorwerk J, Wolters C H and Knosche T R 2015 Influence of the head model on EEG and MEG source connectivity analyses *NeuroImage* **110** 60–77
- [30] Ziegler E, Chellappa S L, Gaggioni G, Ly J Q M, Vandewalle G, Andre E, Geuzaine C and Phillips C 2014 A finite-element reciprocity solution for EEG forward modeling with realistic individual head models *NeuroImage* **103** 542–51
- [31] Vorwerk J, Cho J H, Rampp S, Hamer H, Knosche T R and Wolters C H 2014 A guideline for head volume conductor modeling in EEG and MEG *NeuroImage* **100** 590–607
- [32] Fingberg J, Berti G, Hartmann H, Basermann A, Wolters C H, Anwander A, McCarthy A and Woods S 2003 Bio-numerical simulations with SimBio *NEC Res. Dev.* **44** 140–5 (<http://hdl.handle.net/11858/00-001M-0000-0010-A196-8>)
- [33] MacLeod R S, Weinstein O M, de St Germain J D, Brooks D H, Johnson C R and Parker S G 2004 SCIRun/BioPSE: Integrated problem solving environment for bioelectric field problems and visualization *Biomedical Imaging: Nano to Macro, 2004 IEEE Int. Symp.* pp 640–3
- [34] Opitz A, Paulus W, Will S, Antunes A and Thielscher A 2015 Determinants of the electric field during transcranial direct current stimulation *NeuroImage* **109** 140–50
- [35] Datta A, Bansal V, Diaz J, Patel J, Reato D and Bikson M 2009 Gyri-precise head model of transcranial direct current stimulation: improved spatial focality using a ring electrode versus conventional rectangular pad *Brain Stimul.* **2** 201–7, 7 e1
- [36] Truong D Q et al 2013 Computational modeling of transcranial direct current stimulation (tDCS) in obesity: impact of head fat and dose guidelines *NeuroImage Clin.* **2** 759–66
- [37] Kessler S K, Minhas P, Woods A J, Rosen A, Gorman C and Bikson M 2013 Dosage considerations for transcranial direct current stimulation in children: a computational modeling study *PLoS One* **8** e76112
- [38] Bai S W, Dokos S, Ho K A and Loo C 2014 A computational modelling study of transcranial direct current stimulation montages used in depression *NeuroImage* **87** 332–44
- [39] Shahid S, Wen P and Ahfock T 2013 Numerical investigation of white matter anisotropic conductivity in defining current distribution under tDCS *Comput. Methods Programs Biomed.* **109** 48–64
- [40] Xu J, Healy S M, Truong D Q, Datta A, Bikson M and Potenza M N 2015 A feasibility study of bilateral anodal stimulation of the prefrontal cortex using high-definition electrodes in healthy participants *Yale J. Biol. Med.* **88** 219–25
- [41] Seibt O, Brunoni A R, Huang Y and Bikson M 2015 The pursuit of DLPFC: non-neuronavigated methods to target the left dorsolateral pre-frontal cortex with symmetric bicephalic transcranial direct current stimulation (tDCS) *Brain Stimul.* **8** 590–602
- [42] Alam M, Truong D Q, Khadka N and Bikson M 2016 Spatial and polarity precision of concentric high-definition

- transcranial direct current stimulation (HD-tDCS) *Phys. Med. Biol.* **61** 4506–21
- [43] Galletta E E, Cancelli A, Cottone C, Simonelli I, Tecchio F, Bikson M and Marangolo P 2015 Use of computational modeling to inform tDCS electrode montages for the promotion of language recovery in post-stroke aphasia *Brain Stimul.* **8** 1108–15
- [44] Santos L, Martinho M, Salvador R, Wenger C, Fernandes S R, Ripolles O, Ruffini G and Miranda P C 2016 Evaluation of the electric field in the brain during transcranial direct current stimulation: a sensitivity analysis *IEEE-EMBC* (Orlando, FL: IEEE) pp 1778–81
- [45] Fischer D B, Fried P J, Ruffini G, Ripolles O, Salvador R, Banus J, Ketchabaw W T, Santarnecchi E, Pascual-Leone A and Fox M D 2017 Multifocal tDCS targeting the resting state motor network increases cortical excitability beyond traditional tDCS targeting unilateral motor cortex *NeuroImage* **157** 34–44
- [46] Dmochowski J P, Datta A, Bikson M, Su Y and Parra L C 2011 Optimized multi-electrode stimulation increases focality and intensity at target *J. Neural Eng.* **8** 046011
- [47] Haufe S, Huang Y and Parra L C 2015 A highly detailed FEM volume conductor model based on the ICBM152 average head template for EEG source imaging and TCS targeting *Conf. Proc. IEEE Engineering in Medicine and Biology Society* vol **2015** pp 5744–7
- [48] Suh H S, Kim S H, Lee W H and Kim T-S 2009 Realistic simulation of transcranial direct current stimulation via 3-D high-resolution finite element analysis: effect of tissue anisotropy *2009 Annual Int. Conf. of the IEEE Engineering in Medicine and Biology Society* (IEEE) pp 638–41
- [49] Huang Y and Parra L C 2015 Fully automated whole-head segmentation with improved smoothness and continuity, with theory reviewed *PLoS One* **10** e0125477
- [50] Woods A J et al 2016 A technical guide to tDCS, and related non-invasive brain stimulation tools *Clin. Neurophysiol.* **127** 1031–48
- [51] Antal A, Bikson M, Datta A, Lafon B, Dechent P, Parra L C and Paulus W 2014 Imaging artifacts induced by electrical stimulation during conventional fMRI of the brain *NeuroImage* **85** 1040–7
- [52] Pena-Gomez C, Sala-Lonch R, Junque C, Clemente I C, Vidal D, Bargallo N, Falcon C, Valls-Sole J, Pascual-Leone A and Bartres-Faz D 2012 Modulation of large-scale brain networks by transcranial direct current stimulation evidenced by resting-state functional MRI *Brain Stimul.* **5** 252–63
- [53] Turkeltaub P E, Benson J, Hamilton R H, Datta A, Bikson M and Coslett H B 2012 Left lateralizing transcranial direct current stimulation improves reading efficiency *Brain Stimul.* **5** 201–7
- [54] Spitzer V M and Whitlock D G 1997 *National Library of Medicine Atlas of the Visible Human Male: Reverse Engineering of the Human Body* (Burlington, MA: Jones & Bartlett Learning)
- [55] Grimnes S and Martinsen O G 2000 *Bioimpedance & Bioelectricity Basics* (London, San Diego: Academic)
- [56] Geddes L A and Baker L E 1967 The specific resistance of biological material—a compendium of data for the biomedical engineer and physiologist *Med. Biol. Eng.* **5** 271–93
- [57] Akhtari M et al 2002 Conductivities of three-layer live human skull *Brain Topogr.* **14** 151–67
- [58] Baumann S B, Wozny D R, Kelly S K and Meno F M 1997 The electrical conductivity of human cerebrospinal fluid at body temperature *IEEE Trans. Biomed. Eng.* **44** 220–3
- [59] Gabriel C, Gabriel S and Corthout E 1996 The dielectric properties of biological tissues: I. Literature survey *Phys. Med. Biol.* **41** 2231–49
- [60] Ashburner J and Friston K J 2005 Unified segmentation *NeuroImage* **26** 839–51
- [61] Davies A J 1980 *The Finite Element Method: a First Approach* (New York: Oxford University Press)
- [62] Huang Y, Liu A A, Lafon B, Friedman D, Dayan M, Wang X, Bikson M, Doyle W K, Devinsky O and Parra L C 2017 Measurements and models of electric fields in the *in vivo* human brain during transcranial electric stimulation *Elife* **6** e18834
- [63] Huang Y, Liu A A, Lafon B, Friedman D, Dayan M, Wang X, Bikson M, Doyle W K, Devinsky O and Parra L C 2018 Correction: Measurements and models of electric fields in the *in vivo* human brain during transcranial electric stimulation *Elife* **7** e35178
- [64] Jog M V, Smith R X, Jann K, Dunn W, Lafon B, Truong D, Wu A, Parra L C, Bikson M and Wang D J J 2016 *In vivo* imaging of magnetic fields induced by transcranial direct current stimulation (tDCS) in human brain using MRI *Sci. Rep.* **6** 34385
- [65] Kasinadhuni A K, Indahlastari A, Chauhan M, Schär M, Mareci T H and Sadleir R J 2017 Imaging of current flow in the human head during transcranial electrical therapy *Brain Stimul.* **10** 764–72
- [66] Chauhan M, Indahlastari A, Kasinadhuni A K, Schär M, Mareci T H and Sadleir R J 2018 Low-frequency conductivity tensor imaging of the human head *in vivo* using DT-MREIT: first study *IEEE Trans. Med. Imaging* **37** 966–76
- [67] Camacho D L A, Hopper R H, Lin G M and Myers B S 1997 An improved method for finite element mesh generation of geometrically complex structures with application to the skullbase *J. Biomech.* **30** 1067–70
- [68] Wolters C H, Anwander A, Berti G and Hartmann U 2007 Geometry-adapted hexahedral meshes improve accuracy of finite-element-method-based EEG source analysis *IEEE Trans. Biomed. Eng.* **54** 1446–53
- [69] Aydin U et al 2014 Combining EEG and MEG for the reconstruction of epileptic activity using a calibrated realistic volume conductor model *PLoS One* **9** e93154
- [70] Engwer C, Vorwerk J, Ludewic J and Wolters C H 2017 A discontinuous galerkin method to solve the EEG forward problem using the subtraction approach *SIAM J. Sci. Comput.* **39** B138–64