

# A Dimeric Hydride-Bridged Complex with Geometrically Distinct Iron Centers Giving Rise to an $S = 3$ Ground State

Anne K. Hickey,<sup>†</sup> Samuel M. Greer,<sup>1,§</sup> Juan A. Valdez-Moreira,<sup>†</sup> Sean A. Lutz,<sup>†</sup> Maren Pink,<sup>†</sup> Jordan A. DeGayner,<sup>‡</sup> T. David Harris,<sup>‡</sup> Stephen Hill,<sup>1,^</sup> Joshua Telser,<sup>||</sup> and Jeremy M. Smith<sup>\*,†</sup>

<sup>†</sup>Department of Chemistry, Indiana University, 800 East Kirkwood Avenue, Bloomington, Indiana 47405, United States

<sup>1</sup>National High Magnetic Field Laboratory, Florida State University, Tallahassee, Florida 32310, United States

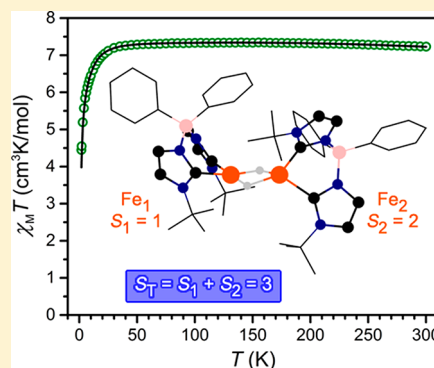
<sup>§</sup>Department of Chemistry & Biochemistry and <sup>^</sup>Department of Physics, Florida State University, Tallahassee, Florida 32306, United States

<sup>‡</sup>Department of Chemistry, Northwestern University, Evanston, Illinois 60208, United States

<sup>||</sup>Department of Biological, Physical and Health Sciences, Roosevelt University, Chicago, Illinois 60605, United States

## Supporting Information

**ABSTRACT:** Structural and spectroscopic characterization of the dimeric iron hydride complex  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  reveals an unusual structure in which a tetrahedral iron(II) site ( $S = 2$ ) is connected to a square planar iron(II) site ( $S = 1$ ) by two bridging hydride ligands. Magnetic susceptibility reveals strong ferromagnetic coupling between iron centers, with a coupling constant of  $J = +110(12) \text{ cm}^{-1}$ , to give an  $S = 3$  ground state. High-frequency and -field electron paramagnetic resonance (HF-EPR) spectroscopy confirms this model. A qualitative molecular orbital analysis of the electronic structure, as supported by electronic structure calculations, reveals that the observed spin configuration results from the orthogonal alignment of two geometrically distinct four-coordinate iron fragments held together by highly covalent hydride ligands.



## INTRODUCTION

Over the past several decades, substantial efforts have been directed toward the study of multinuclear magnetic molecules in which diamagnetic bridging ligands mediate superexchange coupling between paramagnetic metal ions.<sup>1</sup> The nature and magnitude of this coupling is quantified by the exchange constant,  $J$ , which is correlated to the energetic separation of the ground state and excited states. Within the field of molecular magnetism, the development of general strategies to maximize the magnitude of  $J$  represents a central challenge to synthetic chemists, as strong coupling is necessary to maintain collective intramolecular magnetic behavior at high temperature.<sup>2</sup> Because the strength of magnetic exchange between two paramagnetic centers drops precipitously with distance, single-atom bridges provide the ultimate limit in maximizing superexchange interactions. For instance, oxo-bridged complexes often exhibit magnitudes of  $J$  that are well in excess of  $100 \text{ cm}^{-1}$ .<sup>3</sup>

The small size of the hydride ion renders it an ideal bridging ligand for mediating strong superexchange interactions between metal centers. Furthermore, the propensity of hydride to bridge metal ions has long been recognized,<sup>4</sup> although the vast majority of these complexes feature diamagnetic metal centers. It is only recently, with the development of ligands that stabilize low-coordinate metal centers,<sup>5</sup> that paramagnetic

metal hydride complexes, in some cases featuring bridging hydrides, have become readily accessible. Nonetheless, only a limited number of studies on the magnetic properties of these complexes have been reported.<sup>6</sup> Most notably, strong antiferromagnetic coupling was observed for a chromium(II) hydride dimer<sup>7</sup> and for a triangular cobalt(I) hydride trimer,<sup>8</sup> while strong intramolecular, noncompensated antiferromagnetic coupling was reported for a remarkable mixed-valence tetranuclear chromium hydride cluster.<sup>9</sup> In addition, the delocalized bonding of the  $\text{Ni}_2(\mu\text{-H})_2$  core led to an unusual  $S = 1$  ground spin state in  $[\text{Cp}'\text{NiH}]_2$  ( $\text{Cp}' = 1,2,3,4$ -tetraisopropylcyclopentadienyl).<sup>10</sup> Finally, antiferromagnetic coupling mediated through bridging hydrides between cobalt(II) centers in the solid-state compound  $\text{LaSrCoO}_3\text{H}_{0.7}$  was shown to engender a bulk ferrimagnet with a Néel temperature of  $T_N \geq 350 \text{ K}$ .<sup>11</sup>

We have previously reported the ability of bulky bis-(carbene)borate ligands to stabilize coordinatively unsaturated iron complexes, akin to the commonly used  $\beta$ -diketiminato ligands, but with significantly greater donor strength.<sup>12,13</sup> Given that strongly donating ligands are known to stabilize iron(II) in a square planar geometry,<sup>14</sup> we hypothesized that

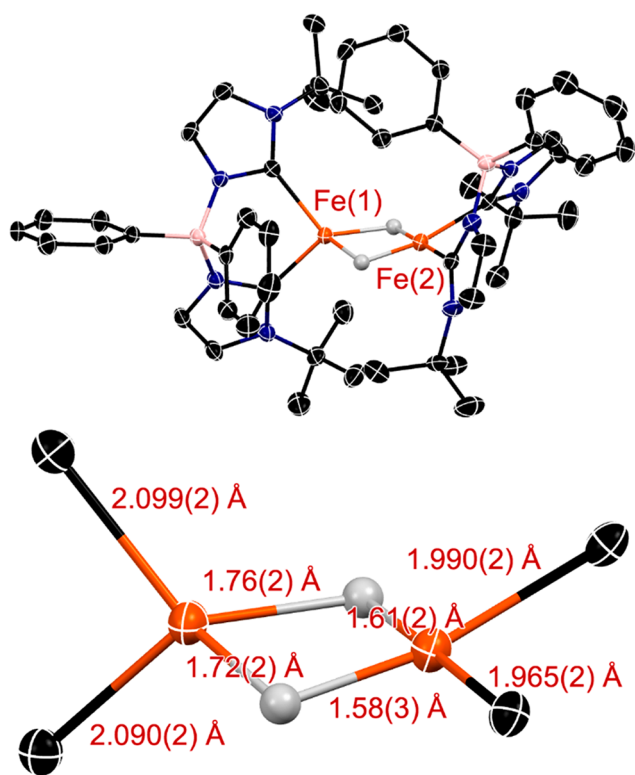
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access to this geometry may similarly occur in a hydride-bridged diiron bis(carbene)borate complex, resulting in unusual magnetic properties. This would be in contrast to the  $\beta$ -diketiminato congeners, where each iron ion resides in a distorted tetrahedral-like geometry.<sup>15</sup> In this report, we describe the synthesis and spectroscopic and magnetic properties and provide a description of the electronic structure of a hydride-bridged diiron bis(carbene)borate complex, wherein strong ferromagnetic coupling between the two iron centers is observed. To our knowledge, this complex provides the first example of ferromagnetic coupling through the hydride ligand.

## RESULTS AND DISCUSSION

The previously reported iron(II) chloride complex  $\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeCl}(\text{THF})$ <sup>12</sup> reacts with equimolar  $\text{NaBEt}_3\text{H}$  to yield the dark blue, hydride-bridged dimeric complex  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  in 88% yield. The molecular structure of the complex, as determined by single-crystal X-ray diffraction, reveals an unusual geometry (Figure 1). As exemplified by  $\beta$ -

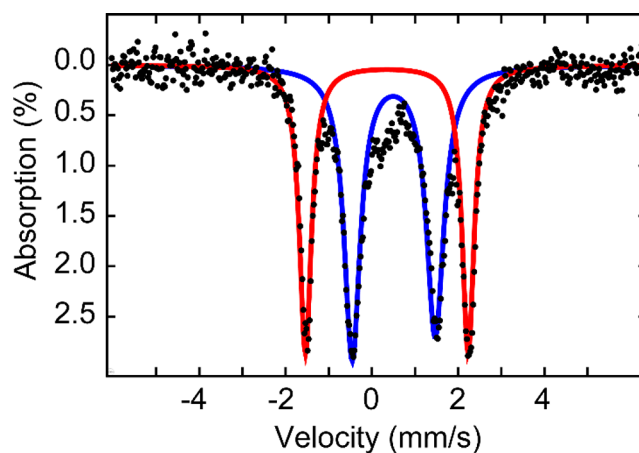


**Figure 1.** Top: X-ray crystal structure of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$ . Thermal ellipsoids shown at 50% probability; peripheral hydrogen atoms omitted for clarity. Bottom: Core structure of the molecule, with selected bond lengths. Black, white, blue, pink, and orange ellipsoids represent carbon, hydrogen, nitrogen, boron, and iron atoms, respectively.

diketiminato complexes,<sup>15</sup> dinuclear complexes with an  $\text{Fe}(\mu\text{-H})_2\text{Fe}$  core generally feature low-coordinate iron centers in equivalent coordination geometries, whereas the two iron centers of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  are distinct, despite each having the same inner coordination sphere. Specifically, Fe(1) resides in a tetrahedral-like environment ( $\tau_4 = 0.7$ ), while Fe(2) is in a square-planar environment ( $\tau_4 = 0.1$ ).<sup>16,17</sup> This complex thus constitutes a rare example of a hydride dimer

comprising geometrically distinct monomers.<sup>18</sup> While we were unable to observe Fe–H vibrations spectroscopically,<sup>19</sup> high resolution mass spectrometry of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  and its deuterated isotopologue,  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeD}]_2$ , combined with the results of preliminary reactivity studies, support the presence of the hydride ligands.<sup>20</sup>

The bond metrics reflect the different coordination environments of the iron centers, with the metal–ligand bond lengths involving Fe(1) being longer than those for Fe(2). While the Fe(1)–H bonds are similar to those in related complexes,<sup>15</sup> the Fe(2)–H bonds are shorter. As a consequence of the shorter bond lengths involving Fe(2), the Fe(1)–Fe(2) distance in  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  (2.5247(4) Å) is slightly shorter than that observed for related iron(II)  $\beta$ -diketiminato complexes (average Fe–Fe distance = 2.56(7) Å).<sup>15</sup> Furthermore, the Fe–C bond lengths at Fe(1) are similar to those observed for four-coordinate high-spin ( $S = 2$ ) iron(II) N-heterocyclic carbene complexes, while those for Fe(2) are similar to those observed for  $S = 1$  square planar iron(II) complexes.<sup>14</sup> These metrics suggest that the core of the complex is composed of a tetrahedral iron(II) ( $S = 2$ ) center that is linked to a square planar iron(II) center ( $S = 1$ ) by two bridging hydride ligands. However, alternative formulations involving a redox change (i.e., disproportionation), such as square planar Fe(I) and tetrahedral Fe(III), both of which have been previously observed, cannot be excluded solely on the basis of the structural data.

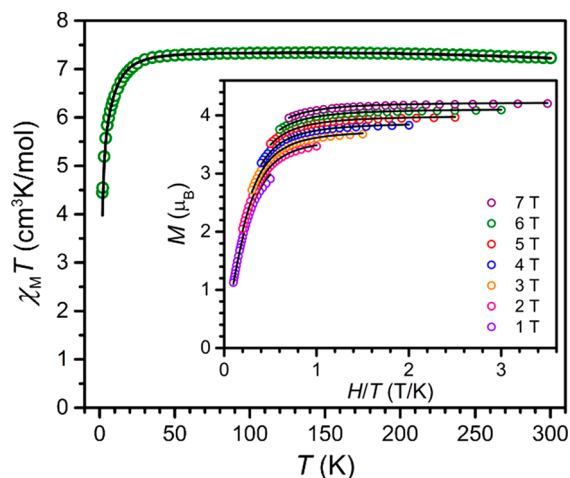


**Figure 2.** Solid-state Mössbauer spectrum of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  at 80 K. Black circles represent experimental data, while red and blue lines correspond to fits of subspectra I and II, respectively (see text for fit parameters).

Given this ambiguity, we used  $^{57}\text{Fe}$  Mössbauer spectroscopy to provide insight into the oxidation states of the two iron centers. As expected for the nonsymmetric structure, the zero-field spectrum of solid  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  indicates two different iron environments (Figure 2). The isomer shift ( $\delta = 0.35$  mm/s) and large quadrupole splitting ( $\Delta E_Q = 3.77$  mm/s) in subspectrum I are both similar to the spectral data of previously reported organometallic  $S = 1$  square planar iron(II) complexes.<sup>14</sup> In contrast, the isomer shift ( $\delta = 0.51$  mm/s) and quadrupole splitting ( $\Delta E_Q = 1.92$  mm/s) of subspectrum II are similar to established examples of tetrahedral  $S = 2$  iron(II) N-heterocyclic carbene complexes.<sup>14b,c</sup> The spectral parameters for this subspectrum are also very similar to those observed for dinuclear  $\beta$ -diketiminato

iron(II) hydride complexes.<sup>15d–f</sup> Alternative fits to the data gave nonreasonable parameters, particularly for the quadrupole splitting.<sup>19</sup> The Mössbauer spectral data therefore confirm that the complex contains two iron(II) centers, one of which is tetrahedral ( $S = 2$ ) and the other square planar ( $S = 1$ ).<sup>21</sup>

To probe the presence of magnetic interactions between iron(II) centers, variable-temperature dc magnetic susceptibility data were recorded for a solid-state sample of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  under an applied field of 5000 Oe (Figure 3). At 300 K,  $\chi_M T = 7.23 \text{ cm}^3 \text{ K/mol}$ , which is much



**Figure 3.** Variable-temperature dc magnetic susceptibility data for  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$ , collected under an applied field of 5000 Oe, with the black line corresponding to a fit of the data. Inset: Low-temperature magnetization data for  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$ , collected at selected fields in the temperature range 2–10 K, with the black lines corresponding to a fit of the data (see text for all fit parameters).

larger than the value of  $\chi_M T = 4.00 \text{ cm}^3 \text{ K/mol}$  expected for magnetically isolated  $S = 2$  and  $S = 1$  iron(II) centers with  $g = 2.00$ . Instead, this value suggests the presence of strong ferromagnetic coupling between iron(II) centers that manifests even up to 300 K. In support of this hypothesis,  $\chi_M T$  continues to increase with decreasing temperature, reaching a maximum value of  $\chi_M T = 7.34 \text{ cm}^3 \text{ K/mol}$  at 155 K. This value falls relatively close to that of  $\chi_M T = 6.00 \text{ cm}^3 \text{ K/mol}$  expected for an  $S = 3$  ground state, with  $g = 2.00$ , arising from ferromagnetic coupling between  $S = 2$  and  $S = 1$  centers. At lower temperatures,  $\chi_M T$  undergoes a precipitous decline, likely stemming from a combination of zero-field splitting (zfs), Zeeman splitting, and/or weak intermolecular antiferromagnetic coupling. To quantify the strength of coupling, the data were modeled with the program PHI<sup>23,24</sup> to provide an exchange constant of  $J = +110(12) \text{ cm}^{-1}$ , with  $g_{\text{iso}} = 2.21(1)$  and  $D = -7.0(4) \text{ cm}^{-1}$ . Note that this axial zfs parameter,  $D$ , acts as an average of the two Fe sites; it is an approximation to avoid complicating this model by overparameterization. In actuality, one would expect such distinct coordination environments at each Fe to give quite different single-ion values of  $D$  (and likely for  $g_{\text{iso}}$  as well).

To assess magnetic anisotropy of the  $S = 3$  ground state of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$ , magnetization data were collected in the temperature range 2–10 K at selected fields from 1 to 7 T. This temperature range should probe only the  $S = 3$  ground state with no convolution from excited states.<sup>25</sup> The splitting of the resulting isofield curves, along with their saturation being

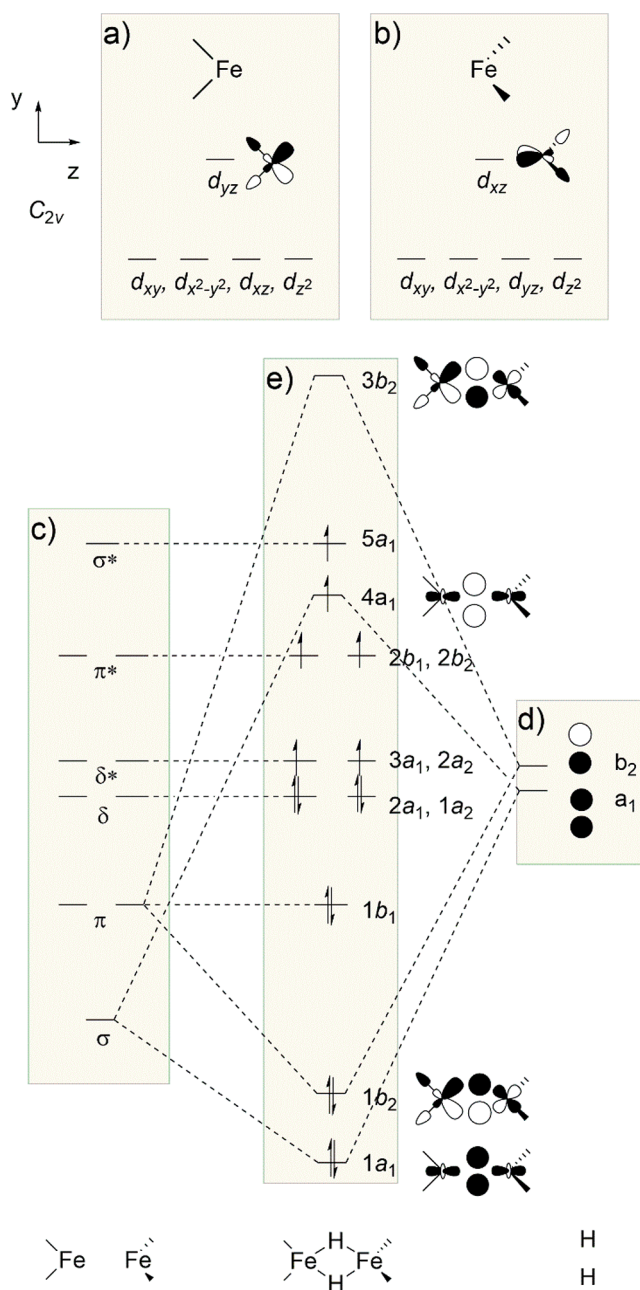
well below the expected  $M = 6.63 \mu_B$  for an  $S = 3$  ground state with  $g = 2.21$ , qualitatively demonstrates the presence of significant magnetic anisotropy (Figure 3, inset). To quantify the anisotropy, the data were fit using the program PHI,<sup>23,26</sup> with a fixed  $g = 2.21$ , to give an axial zero-field splitting parameter of  $D = -10.9(1) \text{ cm}^{-1}$ . This relatively large magnitude of  $D$  for the  $S = 3$  ground state reflects the significant anisotropy associated with iron(II).

High-frequency and -field electron paramagnetic resonance (HFEP) has been widely applied to metal complexes with  $S > 1/2$ , particularly with non-Kramers (integer-spin) ground states.<sup>27</sup> Recognizing that isotopic substitution may change the distance between the two iron centers,<sup>15f,28</sup> and consequently, the extent of magnetic coupling, we applied HFEP to powder  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$ , as well as to  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeD}]_2$ . In principle a species with  $S = 3$  could exhibit a large number of allowed EPR transitions,<sup>29</sup> but if the zfs is large in magnitude relative to the microwave energy, then this might not be the case. Indeed, both isotopologues show identical spectra, which consist of relatively few transitions at low field at frequencies of 600–630 GHz. The observed HFEP spectra, including their temperature dependence, can be satisfactorily simulated using an  $S = 3$  spin Hamiltonian with axial and rhombic second order zfs terms,  $D$  and  $E$ , respectively, and a small, axial fourth-order zfs term,<sup>29</sup>  $B_4^0$ . The effect of both  $E$  and  $B_4^0$  terms on bulk magnetic properties is minimal but in magnetic resonance (HFEP) spectra is apparent. The zfs obtained from HFEP,  $D = -7.1 \text{ cm}^{-1}$ ,  $|E| = 2.1 \text{ cm}^{-1}$ , is in reasonable agreement with that from magnetization, particularly in the sign of  $D$ . The HFEP spectra and their analysis are described in detail in Supporting Information.

The large magnitude of  $J$  observed for  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  is in accord with previous examples of di-,<sup>7,10</sup> tri-,<sup>8</sup> and tetranuclear<sup>9</sup> transition metal complexes supported by bridging hydrides. However, the presence of ferromagnetic coupling is, to our knowledge, unprecedented in hydride-bridged complexes. This coupling scheme likely arises due to orthogonality of spin-bearing d orbitals of the two iron(II) centers, leading to a ferromagnetic alignment of spins via Hund's rules.

The  $S = 3$  spin ground state can be further rationalized via a qualitative molecular orbital diagram (Figure 4).<sup>30</sup> Assuming  $C_{2v}$  symmetry for the molecule, the orbital diagram can be constructed starting from two  $L_2\text{Fe}$  fragments aligned along the Fe–Fe vector (defining the  $z$ -axis) and oriented in an orthogonal fashion. For the square planar site, the  $3d_{yz}$  orbital is destabilized by interaction with the  $\sigma$ -donor orbitals of the  $L_2$  ligand (Figure 4a), but this destabilization is ameliorated by second-order mixing with the higher energy  $4p_y$  orbital. The corresponding interactions for the tetrahedral site involve the  $3d_{xz}$  and  $4p_x$  orbitals because this site is rotated by  $90^\circ$  relative to the square planar site (Figure 4b). Combining these fragments provides the  $L_2\text{Fe}\cdots\text{Fe}L_2$  unit with Fe $\cdots$ Fe orbitals of  $\sigma$ ,  $\pi$ ,  $\delta$  symmetry (Figure 4c). Interestingly, while the  $\pi$ -orbitals are isoenergetic, they are polarized in opposite directions due to the nature of the  $3d$ – $4p$  mixing in each of the two iron sites.<sup>31</sup>

While a number of  $L_2\text{Fe}\cdots\text{Fe}L_2$  fragment orbitals (Figure 4c) have the appropriate symmetry to interact with the  $a_1$  H 1s combination, the most significant interaction will be with the  $\sigma$  ( $a_1$ ) orbital (Figure 4d), providing a low-energy bonding orbital ( $1a_1$ ) and a higher-energy orbital with Fe–H antibonding character ( $4a_1$ ). By contrast, only the  $\pi$  ( $b_2$ ) orbital has the appropriate symmetry to interact with the  $b_2$  H

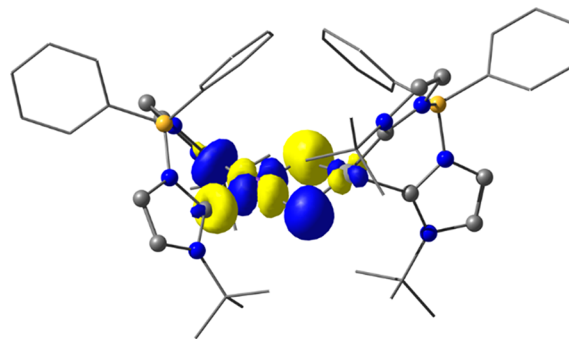


**Figure 4.** Qualitative MO analysis for  $[\text{Ph}_2\text{B}(\text{}^4\text{BuIm})_2\text{FeH}]_2$ , assuming  $C_{2v}$  symmetry and with the  $[\text{Fe}_2\text{H}_2]^{2+}$  core in the  $yz$  plane. Fragment orbitals for (a) square planar and (b) tetrahedral  $L_2\text{Fe}$  iron sites, including  $3d-4p$  mixing, used to construct the fragment orbitals of (c) the  $L_2\text{Fe}\cdots\text{Fe}L_2$  fragment. The final MO diagram obtained from mixing this fragment with the  $2\text{H}$  fragment orbitals (d) is shown in panel (e). Additional orbital mixing in this low symmetry environment has been ignored in the interest of clarity.

$1s$  combination (Figure 4d), providing a low-energy bonding combination ( $1b_2$ ) and a high-energy antibonding combination ( $3b_2$ ). The strong antibonding character of the latter orbital raises it in energy above all of the other  $L_2\text{Fe}\cdots\text{Fe}L_2$  fragment orbitals (Figure 4e). Indeed, due to the strong donor abilities of both the bis(carbene)borate and hydride ligands, this orbital is expected to be significantly destabilized, leading to a  $5a_1$  and  $3b_2$  orbital energy gap that is larger than the spin pairing energy. Accordingly, population of the orbitals with the 16

electrons available to the  $[\text{Fe}_2\text{H}_2]^{2+}$  core is consistent with the experimentally observed  $S = 3$  ground state (Figure 4e).

The results of DFT calculations (B3LYP//def2-TZVP Fe, coordinated C,H/def2-svp non-coordinated C,H,N,B) using the crystallographically determined geometry are generally consistent with this qualitative orbital picture. Most notably, the computed LUMO (Figure 5) shows antibonding interactions between the iron atoms and hydrides, similar to those depicted in the  $3b_2$  orbital (Figure 4e). While lower-energy orbitals of Fe–Fe  $\sigma$ ,  $\pi$ , and  $\delta$  character can be identified, the relatively low actual symmetry of the molecule leads to extensive orbital mixing that complicates the simple qualitative picture (Figure 4e). The calculated HOMO–LUMO gap is 4.8 eV, and the broken symmetry solution provides a coupling constant of  $J = 109 \text{ cm}^{-1}$  (strong coupling formalism),<sup>32</sup> in agreement with the experimentally determined value. Additionally, the computed Mössbauer parameters (B3LYP//CP(PPP) Fe/def2-TZVP coordinated C,H/def2-svp non-coordinated C,H,N,B) for the square planar ( $\delta = 0.15 \text{ mm/s}$  and  $\Delta E_Q = 3.89 \text{ mm/s}$ ) and tetrahedral ( $\delta = 0.37$ ,  $\Delta E_Q = -1.84 \text{ mm/s}$ ) sites show the same relative magnitudes as those observed experimentally.



**Figure 5.** Unrestricted natural orbital representation of the DFT-computed LUMO of  $[\text{Ph}_2\text{B}(\text{}^4\text{BuIm})_2\text{FeH}]_2$  with orbital isodensity at 0.05.

## CONCLUSIONS

The foregoing results demonstrate that the combination of strong donor ability and appropriate steric bulk provided by a bis(carbene)borate ligand stabilizes a unique dinuclear iron(II) hydride-bridged complex exemplifying the two classical four-coordinate geometries: tetrahedral ( $S = 2$ ) and square planar ( $S = 1$ ). Strong ferromagnetic coupling arises from the orthogonal  $d$ -orbital manifolds of the two iron sites, as mediated by the  $\mu$ -dihydride bridge. This ferromagnetic exchange can be further rationalized by a qualitative molecular orbital diagram. More generally, these results demonstrate the potential utility of using metal hydride complexes as building blocks for the assembly of high spin clusters. Specifically, through the appropriate choice of metal complexes, we anticipate that the strong bridging ability of the hydride ligand will facilitate the assembly of multinuclear complexes that feature strong ferromagnetic coupling between the constituent metal ions.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.9b04389.

Full experimental, characterization, and computational details (PDF)

Jmol view of  $[\text{Ph}_2\text{B}(\text{tBuIm})_2\text{FeH}]_2$  (CIF)

## AUTHOR INFORMATION

### Corresponding Author

\*E-mail: [smith962@indiana.edu](mailto:smith962@indiana.edu).

### ORCID

Samuel M. Greer: 0000-0001-8225-3252

T. David Harris: 0000-0003-4144-900X

Stephen Hill: 0000-0001-6742-3620

Joshua Telser: 0000-0003-3307-2556

Jeremy M. Smith: 0000-0002-3206-4725

### Notes

The authors declare no competing financial interest.

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(20) See [Supporting Information](#) for more details.

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(25) The  $S = 3$  ground state results from  $S_1 + S_2 = 2 + 1 = 3$ ; there is also a lower energy excited state with  $S = 2$  and the highest energy state with  $S = 1$ , from  $|S_1 - S_2| = 2 - 1 = 1$ .

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(30) A more extensive analysis is provided in [SI](#).

(31) It should be noted that this qualitative orbital diagram assumes that the two iron sites are isoenergetic, which is unlikely to be correct.

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