

# Momentum-Dark Intervalley Exciton in Monolayer Tungsten Diselenide Brightened via Chiral Phonon

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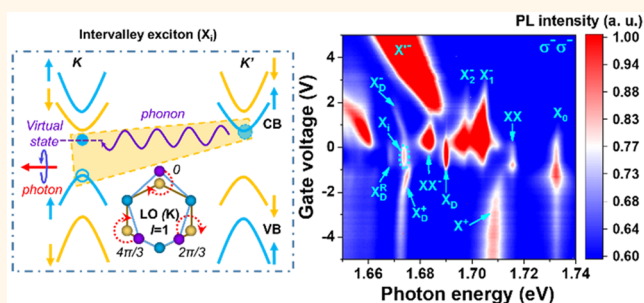
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## Supporting Information

**ABSTRACT:** Inversion symmetry breaking and 3-fold rotation symmetry grant the valley degree of freedom to the robust exciton in monolayer transition-metal dichalcogenides, which can be exploited for valleytronics applications. However, the short lifetime of the exciton significantly constrains the possible applications. In contrast, the dark exciton could be long-lived but does not necessarily possess the valley degree of freedom. In this work, we report the identification of the momentum-dark, intervalley exciton in monolayer WSe<sub>2</sub> through low-temperature magneto-photoluminescence spectra. Interestingly, the intervalley exciton is brightened through the emission of a chiral phonon at the corners of the Brillouin zone (K point), and the pseudoangular momentum of the phonon is transferred to the emitted photon to preserve the valley information. The chiral phonon energy is determined to be ~23 meV, based on the experimentally extracted exchange interaction (~7 meV), in excellent agreement with the theoretical expectation of 24.6 meV. The long-lived intervalley exciton with valley degree of freedom adds an exciting quasiparticle for valleytronics, and the coupling between the chiral phonon and intervalley exciton furnishes a venue for valley spin manipulation.

**KEYWORDS:** intervalley exciton, chiral phonon, magneto-PL, time-resolved PL, tungsten diselenide



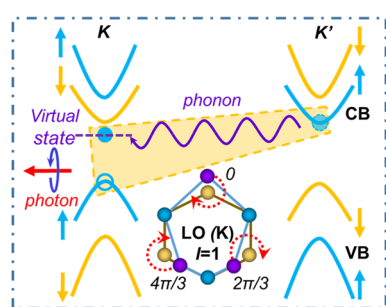
In monolayer transition-metal dichalcogenides (TMDCs), the inversion symmetry breaking and 3-fold rotation symmetry endow the intravalley exciton valley degree of freedom, which can be accessed through circularly polarized

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light<sup>1–9</sup> and ushers in the exciting field of valleytronics.<sup>3–6,10–16</sup> However, the exciton in TMDCs is typically short-lived, with the lifetime of a few to tens of picoseconds,<sup>17–19</sup> limiting the potential applications. The spin-forbidden dark exciton in tungsten-based TMDCs<sup>20–22</sup> possesses a much longer lifetime and caught significant attention for the past few years. However, the out-of-plane radiation of the spin-forbidden dark exciton is not protected by the same valley physics of the in-plane dipole radiation of the bright exciton that is constrained by the 3-fold rotation symmetry and inversion symmetry breaking, and the valley degree of freedom cannot be easily maintained. Intervalley exciton, with the hole residing in the K valley and electron in the K' valley (Figure 1), due to the restricted intervalley

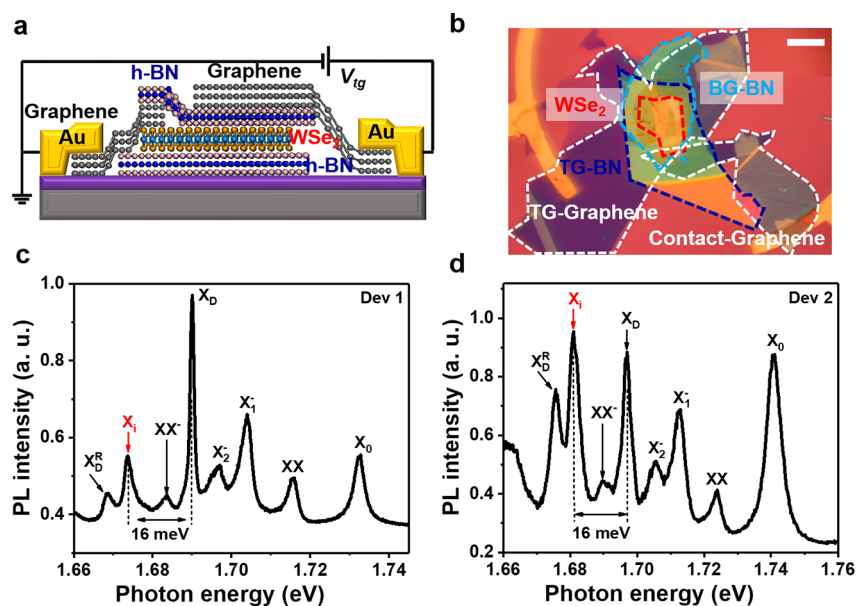


**Figure 1.** Scheme of the intervalley exciton and chiral phonon coupling. Assuming the initial population of excitons in the K valley, the intervalley exciton consists of one electron in the K' valley and one hole in the K valley. The electron transitions to a virtual state in the K valley by emitting a chiral phonon and then recombines with the hole in the valence band of the same valley, emitting a photon with certain helicity. Inset: schematic representation of the chiral phonon mode. Blue sphere: W atom. Yellow and purple spheres are Se atoms in the equilibrium state and the vibration state, respectively.

scattering of the hole, could potentially solve this problem. Intervalley exciton can naturally occur when excitons are selectively generated by optical pumping in the K valley (Figure 1), as the electron can be relatively easy to be scattered to the K' valley, and intervalley exciton thus could be both long-lived due to the momentum mismatch and valley polarized due to the constrained intervalley scattering of the hole<sup>23</sup> because of the large valence band splitting induced by the spin–orbit coupling.

In this work, we identify the intervalley exciton in high-quality monolayer WSe<sub>2</sub> devices through the gate- and magnetic-field-dependent photoluminescence (PL) spectra at low temperature. We also unveil the intervalley exciton coupling with a specialized phonon, the chiral phonon of the WSe<sub>2</sub>. Chiral phonons in monolayer TMDCs, also stemming from 3-fold rotation symmetry, were theoretically predicted at Brillouin-zone corners.<sup>24</sup> The large momentum of the chiral phonon can resolve the dilemma of the momentum mismatch that prohibits the recombination of intervalley excitons. In addition, the pseudoangular momentum (PAM) of the chiral phonon potentially can be transferred to and determine the helicity of the emitted photon through phonon–exciton interaction.

The intervalley exciton and chiral phonon coupling also allows us to directly study the chiral phonon. Chiral phonons play a critical role in intervalley scattering<sup>25</sup> and can be utilized for phononic chirality, phonon-driven topological states, and dissipationless information processing. However, direct probing of the chiral phonon at the Brillouin-zone corners (K and K' valleys) is restrained by the large momentum mismatch between the phonon mode and the photon, which prevents us from using polarization-resolved Raman spectroscopy to investigate the energy and chirality of the phonon mode directly. Very recently, the existence of such chiral phonon modes was verified by transient infrared spectroscopy.<sup>25</sup>



**Figure 2.** Low-temperature PL spectra of BN encapsulated monolayer WSe<sub>2</sub>. (a) Schematic of the BN encapsulated monolayer WSe<sub>2</sub> device, with the carrier density controlled by a top gate. (b) Optical microscope image of a fabricated device, scale bar: 20  $\mu\text{m}$ . (c) PL spectra of device 1 at 4.2 K with the gate voltage of  $-0.44$  V. A CW laser centered at 1.959 eV was applied as the excitation source, with the excitation power of 40  $\mu\text{W}$ . (d) PL spectra of device 2 at 4.2 K with no gate voltage applied. CW lasers centered at 1.959 eV for (c) and 1.879 eV for (d) were applied as the excitation source, with excitation power of 40  $\mu\text{W}$  and 60  $\mu\text{W}$  for (c) and (d), respectively.

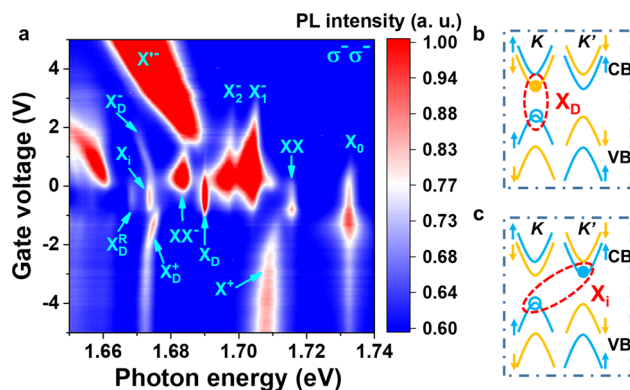
However, the experimental complexity for chiral phonon creation and detection gives rise to uncertainty in determining the chiral phonon energy. In contrast, the energy conservation of the intervalley exciton recombination allows us to determine the phonon energy accurately by analyzing the emitted photon energy through PL spectra, and we experimentally extracted the phonon energy to be  $\sim 23$  meV, in excellent agreement with the prediction from first principle calculations (24.6 meV, Supplementary Table S3).<sup>25</sup> The angular momentum conservation of the process also allows us to determine the phonon PAM to be 1 (inset of Figure 1) by analyzing the emitted photon chirality in the helicity-resolved PL spectra.

Through the valley-resolved PL spectra, we found that the emitted photon retains the valley information due to the PAM transferred from the chiral phonon, and the time-resolved PL (TRPL) spectra show that the lifetime of the momentum-dark intervalley exciton is  $\sim 200$  ps, comparable to that of the spin-forbidden dark exciton ( $\sim 250$  ps) but much longer than that of the bright exciton ( $\sim 5$  ps).<sup>26</sup> The valley-polarized, long-lived intervalley exciton thus adds to the excitonic quasiparticles for the valleytronics, and the coupling between the chiral phonon and the intervalley exciton can be utilized as a tunable knob for manipulating valley degree of freedom in TMDCs.

## RESULTS

We fabricated a BN-encapsulated monolayer WSe<sub>2</sub> device using a dry pickup method described in our previous studies.<sup>26,27</sup> The device structure is shown schematically in Figure 2a, in which both few-layer graphene flakes were employed to work as the contact electrode to the WSe<sub>2</sub> and the semitransparent top-gate electrode, respectively, with the top BN flake working as the gate dielectric. The top-gate voltage allows efficient control of carriers in WSe<sub>2</sub>. The optical microscope image for a typical device is shown in Figure 2b. The PL spectra for two different devices at 4.2 K are shown in Figure 2c,d, in which most of the features have been identified in the previous work, including the bright exciton ( $X_0$ ), the biexciton (XX), and charge biexciton (XX<sup>-</sup>),<sup>27–34</sup> the two negative trions ( $X_1^-$  and  $X_2^-$ ),<sup>4,35–38</sup> the dark exciton ( $X_D$ ),<sup>19–22,39,40</sup> and the dark exciton phonon replica ( $X_D^R$ ).<sup>26</sup> However, the origin of the PL peak at 1.674 eV in Figure 2c and 1.681 eV in Figure 2d remains unknown. Despite the variation of the PL peak position in the two different devices, the energy difference between this emerging peak and that of the dark exciton remains constant,  $\sim 16$  meV for both device 1 (Figure 2c) and device 2 (Figure 2d). We thus believe that the emerging PL peak in the two different devices shares the same origin and we label it as  $X_i$ . This hypothesis is also confirmed by the results from two other devices (see Figure S1).

To explore the origin of the PL peak  $X_i$ , we performed PL spectra as a function of the top gate voltage, which effectively controls the density and type of the charge carriers in the monolayer WSe<sub>2</sub>. As shown in Figure 3a, the intensity of the emerging PL peak as a function of the gate voltage closely follows the gate dependence of the dark exciton peak  $X_D$ . We found the PL intensity of the peak  $X_i$  is maximized between gate voltages  $-1.0$  to  $+0.35$  V, when WSe<sub>2</sub> is charge neutral, which suggests that the  $X_i$  stems from the photon emission of a charge-neutral quasiparticle. The gate dependence rules out the possibility of the Q–K exciton,<sup>41</sup> whose PL will have different gate dependence from that of the spin-forbidden dark exciton ( $X_D$ ). It is worth noting that at the n doping (gate voltage  $> 0.35$  V) and p doping (gate voltage  $< -1.0$  V) sides,



**Figure 3.** Gate-voltage-dependent PL spectra at low temperature. (a) Color plot of PL spectra of device 1 at 4.2 K as a function of the top gate voltage. The excitonic complexes are all well-resolved, including the bright exciton ( $X_0$ ), the biexciton (XX), and charge biexciton (XX<sup>-</sup>), the two negative trions ( $X_1^-$  and  $X_2^-$ ), the dark exciton ( $X_D$ ) and the dark trions ( $X_D^+$  and  $X_D^-$ ), and the dark exciton phonon replica ( $X_D^R$ ). The color represents the PL intensity. The spectra were obtained by photoexcitation with a CW laser centered at 1.879 eV, with an excitation power of 100  $\mu$ W. (b, c) Schematic representations of intravalley spin-forbidden dark exciton ( $X_D$ ) and intervalley exciton ( $X_i$ ).

PL peaks of negative dark trion ( $X_D^-$ ) and positive dark trion ( $X_D^+$ ) emerge (Figure 3a).<sup>22,42,43</sup>

To further illustrate the nature of  $X_i$ , we performed PL spectra measurements as a function of the out-of-plane magnetic field. The PL spectra measurements were performed in a valley-resolved configuration in which we excited the monolayer WSe<sub>2</sub> with either right circularly polarized light or left circularly polarized light ( $\sigma^-$ ), and we detected PL with the same polarization, i.e., ( $\sigma^+\sigma^+$ ) or ( $\sigma^-\sigma^-$ ) configuration.<sup>1,7,8</sup> The PL peaks exhibit a Zeeman splitting as defined by  $E = E_0 \pm \frac{1}{2}g\mu_B B$ , where  $g$  is the Landé  $g$ -factor and  $\mu_B$  is the Bohr magneton. The “+” and “-” signs are for the PL peak energies from the K and K’ valleys, respectively. The PL peaks in the magneto-PL spectra of the ( $\sigma^-\sigma^-$ ) configuration exhibits a linear blue shift due to the valley Zeeman effects,<sup>20,44–53</sup> and the size of the shift in a fixed B field is determined by the  $g$ -factor. It is evident from Figure 4a that the  $X_i$  peak has the largest  $g$ -factor because the shift of the  $X_i$  peak has the steepest slope as a function of the magnetic field. Quantitatively, we can extract the  $g$ -factor from the experimental data of ( $\sigma^-\sigma^-$ ) and ( $\sigma^+\sigma^+$ ) by calculating the Zeeman shift difference between the K and K’ valleys (see Supplementary Note 1), and the results are shown in Figure 4b. The experimentally extracted  $g$ -factor for  $X_i$  is  $\sim -12.5$ , significantly larger than that of the bright exciton ( $\sim -3.6$ ), trions ( $\sim -4.5$  and  $\sim -4.0$  for intervalley trion and intravalley trions, respectively), and the dark exciton ( $\sim -9.3$ ).<sup>19,27,32,33,45</sup> Theoretically, the  $g$ -factor can be determined by the overall contribution of the spin, valley, and orbital momentum components combined, and the theoretically expected value can be used to illustrate the nature of the quasiparticle.<sup>46–53</sup> Considering the charge-neutral nature of the quasiparticle associated with PL peak  $X_i$  and the large  $g$ -factor, the only possibility for the quasiparticle is the intervalley exciton, as schematically shown in Figure 3c. Based on a noninteracting picture (see the SI), the intervalley exciton is supposed to have a  $g$ -factor of  $-12$ , much larger in magnitude than the

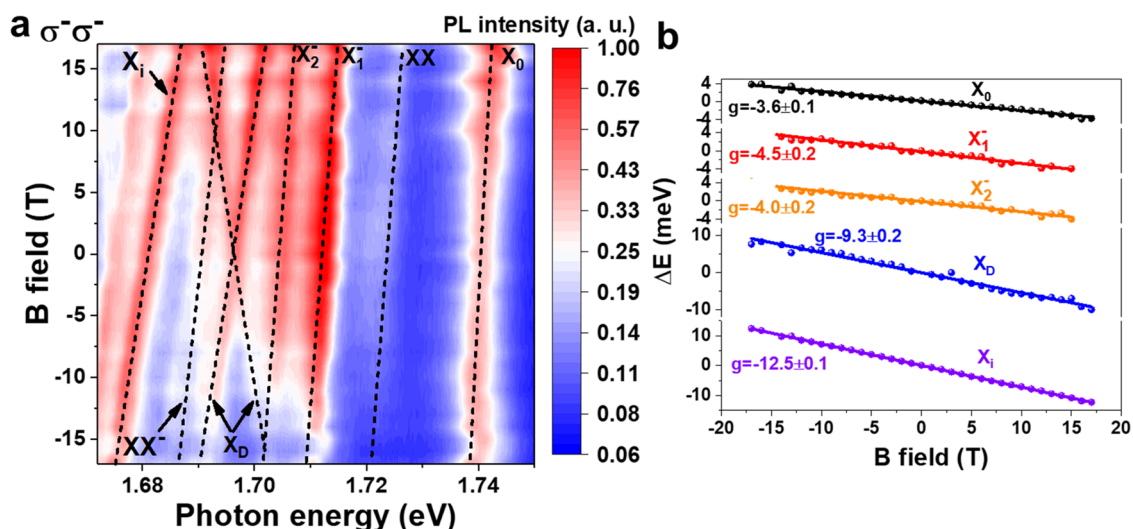


Figure 4. Magnetic-field-dependent PL spectra at low temperature. (a) Color plot of the PL spectra of device 2 at 4.2 K as a function of the out-of-plane magnetic field for ( $\sigma^- \sigma^-$ ) configuration. The spectra were obtained with photoexcitation of a CW laser centered at 1.959 eV, with excitation power of 40  $\mu$ W. (b) Extracted  $g$ -factor from a linear fit of the Zeeman splitting obtained from  $\sigma^- \sigma^-$  and  $\sigma^+ \sigma^+$  configurations.

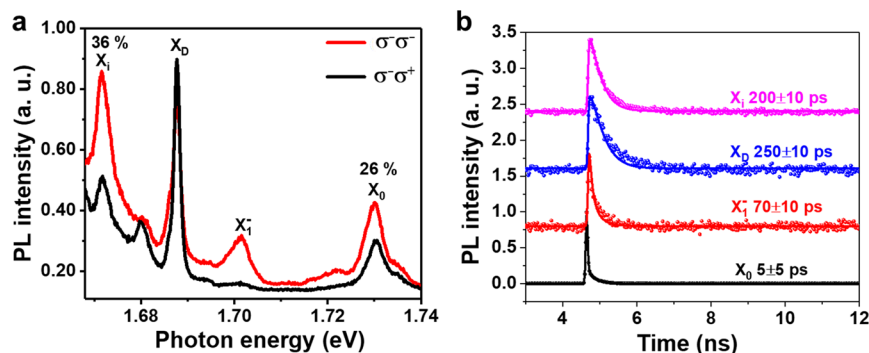


Figure 5. Valley polarization and time-resolved PL. (a) Valley-resolved PL spectra at 43 K with left circularly polarized excitation ( $\sigma^-$ ) and left ( $\sigma^-$ ) or right ( $\sigma^+$ ) circularly polarized detection. The spectra were obtained with the photoexcitation of a CW laser centered at 1.797 eV, with the excitation power of 100  $\mu$ W. The valley polarization, defined as  $P = \frac{I_{\sigma^-} - I_{\sigma^+}}{I_{\sigma^-} + I_{\sigma^+}}$ , where  $I$  stands for PL intensity and is  $\sim 26\%$  for  $X_0$  and  $\sim 36\%$  for  $X_1$ . (b) Time-resolved PL spectra of the distinct PL peaks shown in (a). The spectra were obtained with the pulsed laser excitation (pulse width  $\sim 120$  fs) centered at 2.756 eV, with the excitation power of 50  $\mu$ W. The lifetime of each excitonic complex is obtained through the convolution with the response from the laser pulse (see the SI).

theoretically expected value for the bright exciton ( $g$ -factor of  $-4$ ) and the dark exciton ( $g$ -factor of  $-8$ ). Our experimentally extracted  $g$ -factor for the  $X_i$  peak is in excellent agreement with the theoretically expected value of the intervalley exciton. In contrast, we expect significantly different  $g$ -factor for the Q–K exciton, as the orbital component of the conduction band at Q point is different from that of the K valley, and the Berry phase of the Q point is zero.<sup>54</sup>

Based on the results from the gate and magnetic-field dependent PL spectra, we identify the  $X_i$  peak to be associated with the intervalley exciton as shown in Figure 3c. The intervalley exciton consists of an electron and a hole residing in different valleys, and the direct recombination is forbidden due to their momentum mismatch. Therefore, additional quasi-particles with momentum  $\mathbf{K}$  must be involved in the emission process to satisfy momentum conservation. Interestingly, the  $X_i$  peak exhibits a valley polarization as large as  $\sim 36\%$ , as observed in valley-resolved PL spectroscopy (Figure 5a). The valley polarization of the  $X_i$  requires the quasi-particle to comply with 3-fold rotation symmetry and have well-defined PAM. Furthermore, the energy of peak  $X_i$  is highly

reproducible in all four samples (with respect to the bright exciton). We thus conclude that the involved quasi-particle should be intrinsic to monolayer WSe<sub>2</sub>, and it is likely to be a phonon with momentum  $\mathbf{K}$ . Such phonon-assisted emission of the intervalley exciton is already illustrated in Figure 1, where the intervalley exciton scatters to a virtual intravalley state through emitting a phonon.

To gain further knowledge of this phonon mode, we analyze the behavior of  $X_i$  peak in detail. The intervalley dark exciton is theoretically expected to have an energy about 6 meV higher than that of the intravalley dark exciton due to the exchange interaction,<sup>4,20,35,37,45,55</sup> and our measurement of the splitting between the intervalley and intravalley n-trions reveals a value of  $\sim 7$  meV (energy difference between  $X_1^-$  and  $X_2^-$ ). Combined with the experimentally determined separation of 16 meV between the intravalley dark exciton ( $X_D$ ) and the  $X_i$  peak, we determine the energy of the phonon mode to be around 23 meV. It is worth noting that the temperature-dependent PL spectra (Figure S1e) show that, despite the shift of the  $X_D$  and  $X_i$  as a function of the temperature, the energy separation between these two peaks remains a constant,  $\sim 16$  meV,

consistent with the phonon-assisted recombination picture (assuming that the exchange interaction is not a sensitive function of temperature).

In addition, the peculiar valley polarization of the  $X_i$  peak bears important information about the symmetry properties of the phonon mode. Because holes have much longer valley lifetime than electrons in  $\text{WSe}_2$ ,<sup>23</sup> the bright excitons in the K valley created by  $\sigma^+$  excitation will mainly generate intervalley excitons with holes maintained in the K valley and electrons scattered into the  $K'$  valley. As a result, if the phonon mode does not have PAM, the intervalley dark exciton emission is expected to show the opposite circular helicity as the excitation light,<sup>21</sup> similar to the case of interlayer exciton emission in AB-stacked  $\text{WSe}_2/\text{MoSe}_2$  heterostructure.<sup>56</sup> In contrast, a large positive circular helicity is observed experimentally for the  $X_i$  peak, suggesting that the phonon mode with momentum  $\mathbf{K}$  should have additional PAM of  $-2$  (or equivalently,  $+1$ ).

Theoretical analysis<sup>24</sup> reveals that at the corner of the Brillouin zone (K or  $K'$  point) the 3-fold rotational symmetry endows phonon eigenmodes with a PAM which includes both orbital and spin parts. The orbital PAM for sublattice W and Se at K point can be determined through phase change under counterclockwise  $120^\circ$  rotations, and the total PAM of each phonon mode can be obtained after further analyzing the spin PAM of each phonon eigenmode (see Figure S6 and Table S3). Since the system is symmetric with regard to a mirror operation with respect to the monolayer  $\text{WSe}_2$  plane ( $M = 1$ ), we determine the phonon mode involved in the intervalley exciton recombination process to be the chiral phonon mode  $\text{LO}(E')$  at K, with the energy calculated to be 24.6 meV (Table S3). This chiral phonon mode features a unidirectional circular rotation as illustrated in the inset of Figure 1, and it has been previously observed experimentally by investigating the transition between A and B exciton states with transient infrared spectroscopy.<sup>25</sup> The observation of intervalley exciton PL here provides another direct evidence of the chiral phonon and its intriguing capability in manipulating the symmetry and valley selection rules of exciton states. In addition, owing to the simplicity of the intervalley exciton emission process without involving an additional infrared photon, the energy of the chiral phonon mode ( $\sim 23$  meV) can be accurately determined, which is in excellent agreement with the theoretical prediction of 24.6 meV. It is worth noting that the dark exciton emission in Figure 5a is not valley polarized, consistent with our expectation, since the observed  $X_D$  emission arises from out-of-plane dipole radiation, which is  $p$ -polarized and has equal intensity from the left-polarized and right-polarized PL in our detection scheme.<sup>27</sup>

Finally, we performed time-resolved PL measurements using the time-correlated single photon counting (TCSPC) technique. The second harmonic generation (SHG) signal from a Ti:sapphire oscillator (80 MHz,  $\sim 120$  fs pulse width), centered at 2.756 eV, was used as the excitation source. The direct observation of the sharp intravalley dark exciton PL peak (Figure 5a) enables time-resolved PL measurements, which reveals the lifetime of the dark exciton. As shown in Figure 5b, the lifetime of the  $X_i$  peak is determined to be around 200 ps. It is significantly longer than that of the bright exciton ( $\sim 5$  ps) but comparable to the lifetime of the spin-forbidden intravalley dark exciton ( $\sim 250$  ps),<sup>17,19,20,28,36</sup> which rules out the possibility of defects, since they usually possess even longer lifetime.<sup>29,57</sup> The long lifetime of the  $X_i$  peak also suggest that the measured high valley polarization (36%, Figure 5a) arises

from decreased channels of valley depolarization, which also rules out the possibility of defects and strongly suggests the angular momentum conservation in the recombination process, consistent with our interpretation involving the chiral phonon.

## CONCLUSIONS

In summary, we have identified the PL peak associated with the momentum-dark, intervalley exciton through gate and magnetic-field dependent PL spectroscopy at low temperature. The energy and momentum conservation demand a  $\mathbf{K}$  phonon with particular energy to be involved to brighten the intervalley exciton. The robust valley polarization also requires the angular momentum conservation in the phonon assisted recombination process, which can be satisfied with the chiral  $\mathbf{K}$  phonon possessing a well-defined PAM. The long-lived intervalley exciton with valley information coded sheds light on promising routes of realizing excitonic valleytronics, and the coupling of the intervalley exciton with the chiral phonon could inspire future endeavor of exploiting chiral phonon for valleyspin manipulation.

## METHODS

### Fabrication of h-BN-Encapsulated Monolayer $\text{WSe}_2$ Device.

We fabricated the h-BN-encapsulated monolayer  $\text{WSe}_2$  device by the well-known dry transfer technique.<sup>26,27</sup> The monolayer  $\text{WSe}_2$ , few-layer graphene, and few-layer h-BN were first exfoliated onto 285 nm  $\text{SiO}_2/\text{Si}$  substrate, followed by the inspection through optical microscope. Afterward, the PPC (polypropylene carbonate) stamp sequentially picks up the few-layer h-BN, monolayer  $\text{WSe}_2$ , few-layer graphene, and another few-layer h-BN. The whole stack of the constructed vdW structure was placed onto prepatterned Au electrodes, and the PPC is removed by heating up to  $90^\circ\text{C}$ . Then the whole chip was soaked in chloroform for 2 h to remove the PPC polymer residue. Finally, we transferred another few-layer graphene flake to work as the top-gate electrode, on top of the top h-BN flake.

**Magneto-PL and Time-Resolved PL Measurements.**<sup>26,27</sup> We applied a confocal micro-PL setup to measure the magneto-PL with the out-of-plane magnetic field. A linear polarized light is converted into circularly polarized light by a quarter waveplate, and the laser is focused by a 50 $\times$  objective (NA:  $\sim 0.65$ ) to excite the  $\text{WSe}_2$  sample, with a spot size of  $\sim 2\ \mu\text{m}$ . The PL is collected with the same objective and goes through the quarter waveplate to be converted into linear light. The assembly of a half waveplate and a linear polarizer is used to distinguish the  $\sigma^+\sigma^+$  and  $\sigma^-\sigma^-$  configurations and the signal is detected by the CCD camera attached to the spectrometer.

We measured the TRPL using the time-correlated single photon counting (TCSPC) technique, with the pulsed laser excitation centered at 2.756 eV and the excitation power of 50  $\mu\text{W}$ . For each PL peak, we integrated the PL intensity around the center with an uncertainty of  $\pm 1$  meV for the time-resolved information. The pulsed laser is the second harmonic generation (SHG) signal from a Ti:sapphire oscillator (80 MHz,  $\sim 120$  fs pulse width). We performed convolution and extracted the lifetime of different excitonic complexes by a single exponential function  $I = Ae^{-t/\tau}$  convolved with the response of the laser as a kernel.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsnano.9b06682>.

Additional information on the g-factor calculation, lifetime measurement, calculations of the K-point phonon modes, and data from other devices (PDF)

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## Notes

The authors declare no competing financial interest.

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