

LETTER

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## Letter

# In-field critical current performance of 4.0 $\mu\text{m}$ thick film REBCO conductor with Hf addition at 4.2 K and fields up to 31.2 T

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## Abstract

We present results on the in-field critical current ( $I_c$ ) performance of 4.0  $\mu\text{m}$  thick REBCO film with 15% Hf addition with fields up to 31.2 T and field orientations in the B||ab plane and B||c axis. Unlike the behavior at B||c, the critical current at B||ab is only very weakly dependent on field, decreasing from self-field to 31.2 T by only 22%, i.e. from the self-field value of  $\sim 7700$  A/4 mm width to  $\sim 6300$  and 5812 A/4 mm width at 14 and 30 T, respectively. These values are remarkably 3 and 5.7x higher than the corresponding critical currents at B||c. The in-field behavior of the present 15% Hf sample at field orientation B||c axis is nearly identical to the previously reported record values found in 4.3 and 4.6  $\mu\text{m}$  thick 15% Zr samples in terms of critical current density. In contrast to the pinning force behavior in the B||c orientation, which saturates to a constant value of  $1.7 \text{ TN m}^{-3}$  above  $\sim 5\text{--}6$  T, the pinning force in the B||ab orientation increases near-linearly, reaching a remarkable value of over  $11.5 \text{ TN m}^{-3}$  at 31.2 T. These results demonstrate the potential of thick REBCO conductors at 4.2 K for high field and energy density applications, in particular where the magnetic field is contained near the ab-plane.

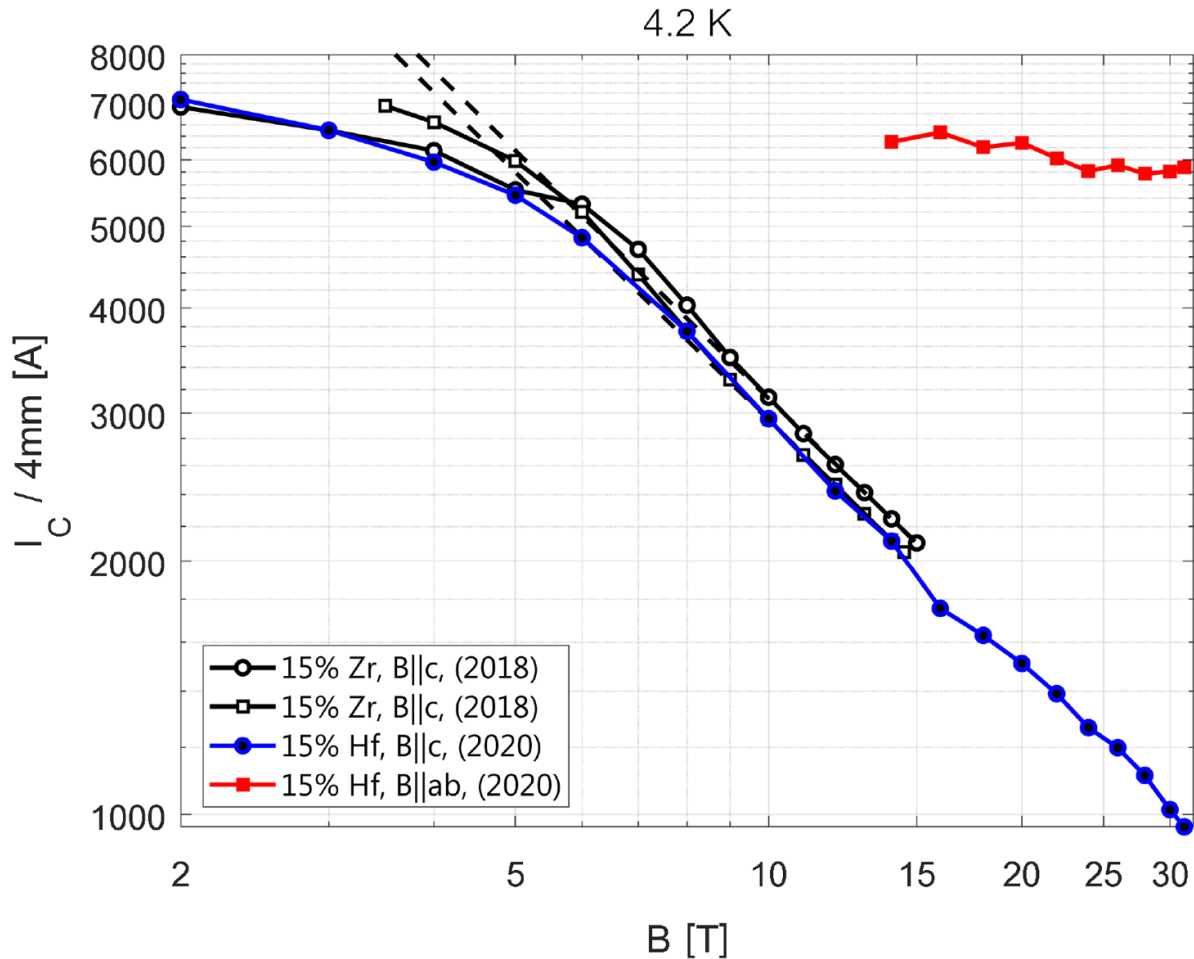
Keywords: — High temperature superconductors, MOCVD, superconducting films, flux pinning, critical current

(Some figures may appear in colour only in the online journal)

## 1. Introduction

With continuous progress in the development of REBCO coated conductors, they have become increasingly attractive for low temperature (4.2–20 K) and intermediate to very high

field applications, due to their high current density and the ability to operate at much higher fields compared to LTS conductors, opening possibilities for attaining very high magnetic fields and/or high magnetic energy density. Recently, an exciting fusion reactor project has been announced that



**Figure 1.** Critical current vs. magnetic field at 4.2 K for field orientations  $B\parallel ab$  and  $B\parallel c$  for the present 15% Hf-containing GdYBCO tape and  $I_c(B\parallel c)$  of the previously reported 15% Zr containing GdYBCO tape [16]. The dashed lines are log-linear fit to the data above 9 T, with alpha value of 1.02.

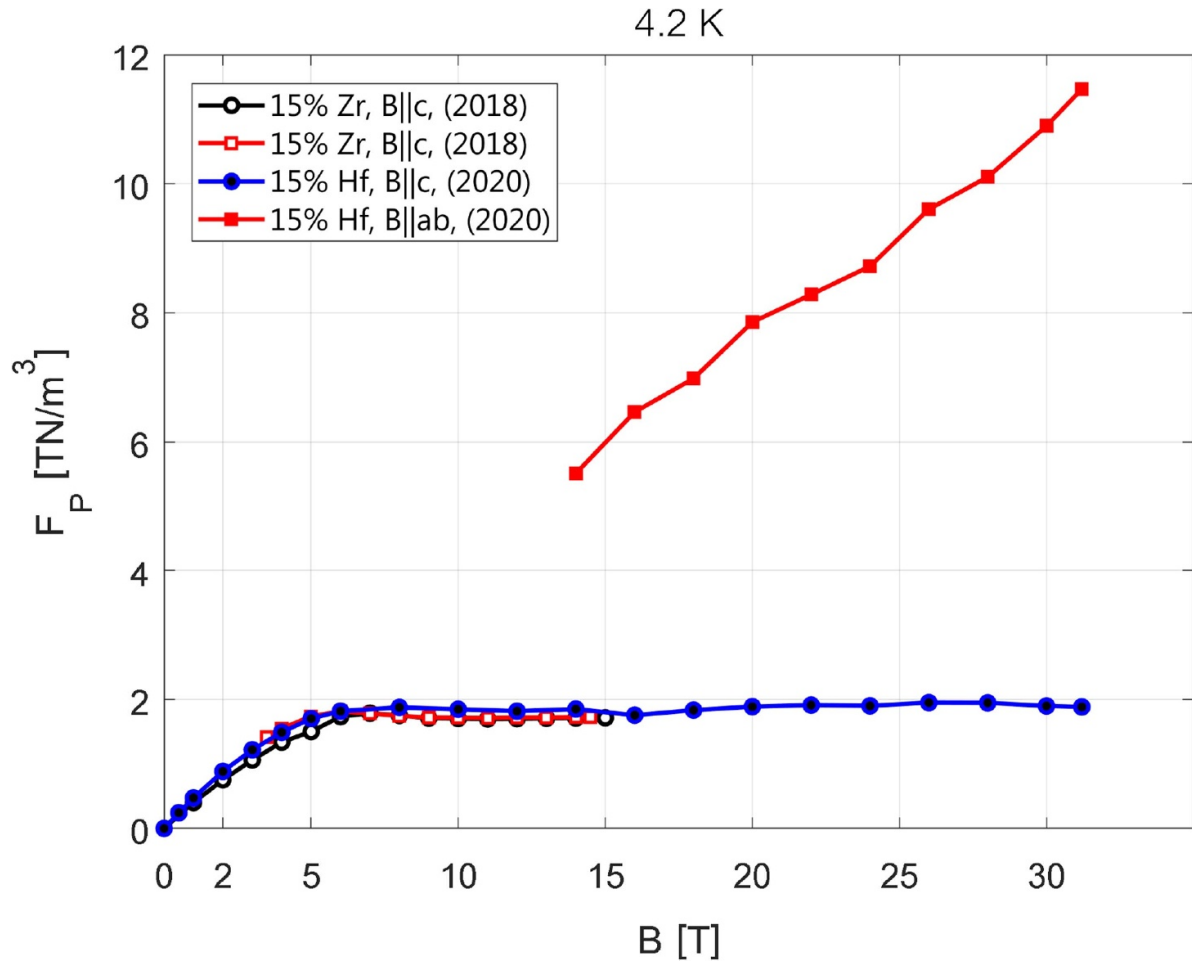
would provide a net energy gain with a footprint of just a fraction of the ITER reactor, where the enabling technology is the REBCO conductor producing a field of 12 T (maximum field at conductor 21 T) at 20 K [1]. One can foresee the potential for significant further performance increase by increasing the critical current of the conductor and/or lowering the operating temperature towards 4.2 K. Other examples include the record 42.5 T magnet with a 13.5 T REBCO coil insert operating at 4.2 K, all-HTS magnet demonstrations and/or designs with fields of 25–32 T, a conceptual study of a steady state 100 T HTS magnet, various SMES devices operating at temperatures of 4.2–20 K, MRI coils and superconducting motors and generators [2–15].

Significant progress in understanding and overcoming obstacles to growing thick REBCO films with high current densities has been made in recent years [16–32]. Additionally, the understanding and control of the parameters influencing the effectiveness of artificial pinning centers (APC) in increasing in-field critical currents has significantly progressed from both fundamental and processing aspects [16–18, 20, 33–61]. Among the most studied systems are  $BaZrO_3$  (BZO) and  $BaHfO_3$  (BHO), with varying reports regarding whether either of the two systems has any clear advantage over the other.

We have recently reported the performance of  $4+\mu m$  thick REBCO films at 20–40 K and 4.2 K [16, 17]. The engineering current density ( $J_e$ ) of the conductor at 4.2 K, 14 T at the ‘worst’ orientation  $B\parallel c$  axis was measured at  $5\text{ kA mm}^{-2}$  which constitutes an over five-fold increase over  $Nb_3Sn$ . However, the remaining unanswered question on the performance of these thick films is their performance at  $B\parallel ab$ . For toroid-based designs (fusion tokamak, SMES), this field orientation is the most relevant, as the field is concentrated near the ab plane direction. In this letter, we report on both  $B\parallel ab$  and  $B\parallel c$  performance of a  $4.0\mu m$  thick film REBCO tape with 15% Hf addition at 4.2 K and fields up to 31.2 T. The main objectives are: 1.) to compare the optimized BHO sample to the performance of the previously reported optimized BZO sample at  $B\parallel c$  where nanorods are effective pinning centers, and 2.) to compare the performance of the same tape in the  $B\parallel ab$  plane relative to  $B\parallel c$ .

## 2. Experimental

A (GdY)-Ba-Cu-O (GdY)BCO film containing 15% Hf was deposited using the advanced MOCVD (A-MOCVD) system



**Figure 2.** Pinning force vs. magnetic field at 4.2 K for field orientations  $B||ab$  and  $B||c$  for the present 15% Hf-containing GdYBCO tape and  $F_p(B||c)$  of the previously reported 15% Zr-containing GdYBCO tape [16].

[16–19], under deposition conditions described in [16]. The film was deposited on  $\text{LaMnO}_3$  capped Hastelloy C276 tapes 12 mm wide over a deposition length of 30 cm. A micro-bridge with dimensions  $80 \times 500 \mu\text{m}$  was made using a focused ion beam (FIB) in order to enable measurement of very high currents at  $B||ab$ . In addition to the small size of the bridge that enabled measuring currents within the limits of the experimental setup, the FIB bridge also offers the advantage over etched bridges of the capability to sustain higher currents without delamination. The critical currents at  $B||ab$  and  $B||c$  orientations were measured on the same sample, as the FIB bridge survived measurements down to 0 T (self-field). The 2D-XRD measurements were performed in a new state-of-the-art facility at the University of Houston utilizing a microfocus source, Goebel mirror, 0.5 mm collimator and a high-resolution Vantec 500 detector. Transmission electron microscopy (TEM) was performed on a JEOL 2000FX microscope.

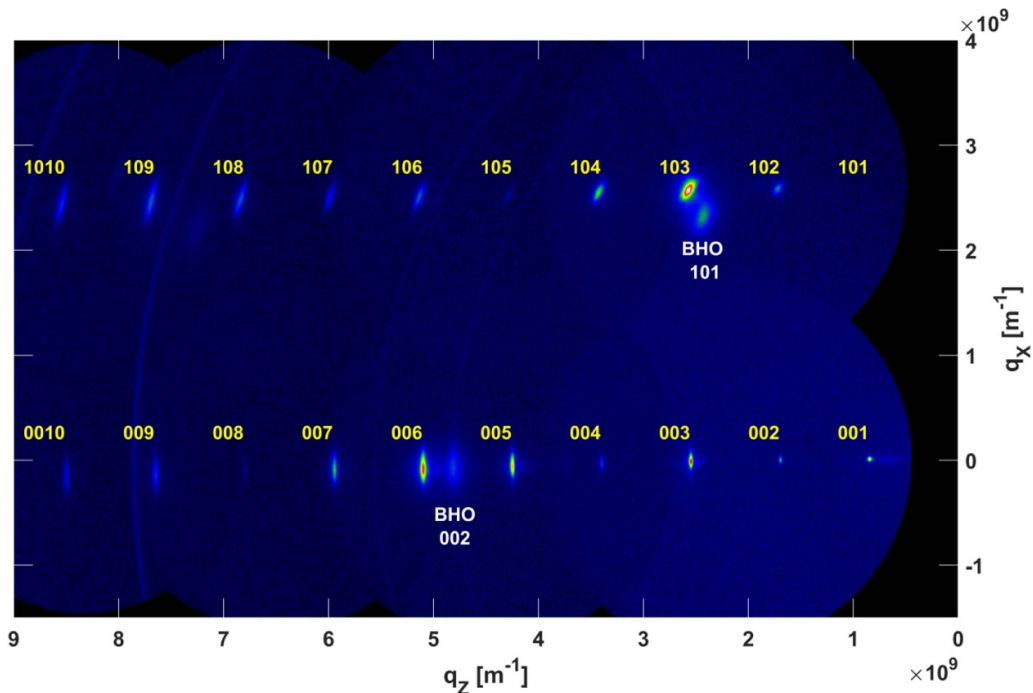
### 3. Results and discussion

Figure 1 shows the critical current at 4.2 K of the measured 15% Hf-containing GdYBCO sample at both orientations

$B||ab$  and  $B||c$ . For convenience, the previously-reported data on the 15% Zr samples at field orientation  $B||c$  are also superimposed on the graph. Two remarkable observations can be made from the data:

- The critical current at  $B||ab$  reaches 6300 A/4 mm-width at 14 T and 5812 A/4 mm-width at 30 T. These extraordinarily high values are 3 and 5.7x higher than the values obtained from the same sample at  $B||c$  with the same fields (2114 and 1013 A/4 mm, respectively).
- The critical current of the 15% Hf sample at  $B||c$  is nearly identical or practically identical to the previously reported 15% Zr samples, as well as the alpha value of 1.02 ( $J \sim B^{-\alpha}$ ).

The corresponding critical current density values ( $J_c$ ) at  $B||ab$  are 40.5 and 36.3  $\text{MA cm}^{-2}$  at 14 and 30 T, respectively. In order to obtain an estimate of engineering current density, we use the same values as for the Zr sample [16] of 50  $\mu\text{m}$  thick Hastelloy, 200 nm buffer, 40  $\mu\text{m}$  Cu stabilizer and 3  $\mu\text{m}$  thick silver, which leads to  $J_e$  values of 16.2 and 14.9  $\text{kA mm}^{-2}$  at  $B||ab$ , 14 and 30 T, respectively. Compared to  $J_e \sim 1 \text{ kA mm}^{-2}$  of  $\text{Nb}_3\text{Sn}$  at 15 T [62], the performance is higher by a factor of



**Figure 3.** 2D-XRD diffraction pattern of the 4.0  $\mu\text{m}$  thick GdYBCO film with 15% Hf addition. The pattern reveals sharp out of plane texture of GdYBCO, strong BZO (101) and (002) peaks, as well as absence of REO (004) and (222) peaks or any other secondary phase, including a-axis oriented GdYBCO grains.

over 16 times. The corresponding values for  $B\parallel c$  are 5.4 and 2.6  $\text{kA mm}^{-2}$  at 14 and 30 T.

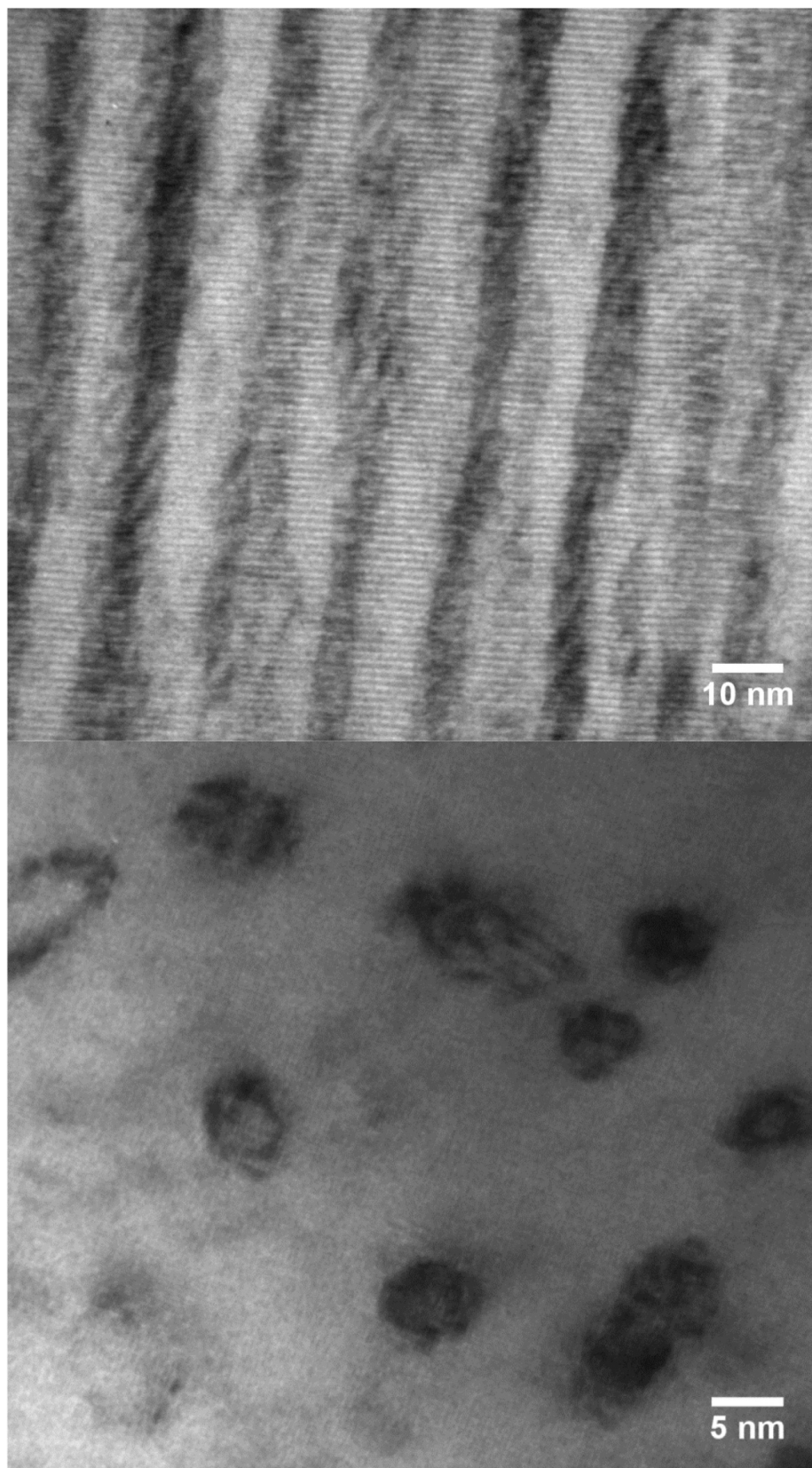
Measurements of critical current of REBCO at 4.2 K,  $B\parallel ab$  are relatively scarce (e.g. [55, 63]). The measurements on REBCO tapes from four commercial manufacturers at 12 and 14 T indicate about six-fold anisotropy of  $J_c$  between  $B\parallel ab$  and  $B\parallel c$  [63, 64]. From the measurements provided on a high  $J_e$  REBCO tape reported in [55, 62], we estimate  $J_e$  at  $B\parallel ab$  and  $B\parallel c$  of 3.42 and 0.53  $\text{kA mm}^{-2}$  at 15 T and 2.72 and 0.33  $\text{kA mm}^{-2}$  at 30 T, respectively. This constitutes  $J_c$  anisotropy of 6.5 at 15 T and 8.3 at 30 T. Regarding the  $J_e$  values at  $B\parallel ab$ , the present data are higher by 4.7 and 5.5x at 15 and 30 T, respectively.

Figure 2 shows the pinning force for the same samples—the present 15% Hf-added GdYBCO at  $B\parallel ab$  and  $B\parallel c$ , as well as the previously reported 15% Zr-added GdYBCO at  $B\parallel c$ . The pinning force in the  $B\parallel c$  direction reaches a saturation point at  $\sim 5\text{--}6$  T, identical for both Hf and Zr containing samples, after which it is constant up to 31.2 T at a level of 1.7  $\text{TN mm}^{-2}$ . Remarkably, unlike the  $B\parallel c$  direction, the pinning force in the  $B\parallel ab$  direction near-linearly increases, reaching near 6  $\text{TN m}^{-3}$  at 15 T and exceeding 10  $\text{TN m}^{-3}$  at 30 T. The shape of the curve indicates no signs of saturation of  $F_p$  at fields immediately above 31.2 T, where it reaches 11.5  $\text{TN m}^{-3}$ .

The magnitudes of the pinning force found in the present sample have strong implications in terms of opportunities that these results present for applications. For applications where field magnitude is of primary interest, the induced magnetic field  $B$  is directly proportional to  $I_c$  and the number of turns per unit length,  $n$ ,  $B \sim knI$ , where  $k$  is a geometric constant

dependent on the shape of the device. For a device operating at  $B\parallel ab$ , this would constitute an 8–12 fold increase in peak magnitude for the same amount of tape, or alternatively, a 8–12 fold reduction of the amount of tape needed for the same field. However, for applications where magnetic energy density is of primary concern,  $E \sim B^2 = knBI = knF_p$ , therefore directly linearly proportional to the pinning force. This suggests that the magnetic energy density can be drastically increased by operating the device at 4.2 K and with very high fields, e.g. 30 T. On the other hand, for  $B\parallel c$ ,  $F_p$  is constant above 5–6 T, indicating that for energy density, there are no benefits to operating at fields higher than this saturation point. The calculated maximum compressive stress during the measurement at  $B\parallel ab$ , 30 T was  $\sim 40$  MPa, which the sample survived, as the measurements were completed in the sequence 15–31.2 T, followed by 31.2–0 T. The  $I_c$  values of repeated measurements between 15–31.2 T were near identical.

Figure 3 shows the 2D-XRD scan of the present Hf-containing sample, revealing sharp texture akin to that of the previously reported Zr sample [16], with no signs of unwanted secondary phases other than the  $\text{BaHfO}_3$ . The sample is purely c-axis oriented, with no trace of a-axis oriented grains, as also confirmed by scanning electron microscopy (SEM). One notable difference compared to the Zr-containing sample [16] is a complete lack of any of the rare earth oxide ( $\text{RE}_2\text{O}_3$ —REO) peaks in the present Hf sample, as can be seen by comparing the strong 004 and 222 REO peaks reported in [16] compared to the present sample. The fact that the Hf and Zr samples show near-identical  $B\parallel c$  performance, despite a large difference in REO signal intensity, suggests that REO precipitates play no significant role in  $B\parallel c$  pinning at 4.2 K. Further, by comparing



**Figure 4.** Cross-section and plane-view TEM micrographs of the 4.0  $\mu\text{m}$  thick, 15% Hf-containing GdYBCO sample. The microstructure reveals continuous  $\text{BaHfO}_3$  nanorods aligned along c-axis, and almost complete absence of  $\text{RE}_2\text{O}_3$  precipitates. The average nanorod spacing and diameter are 18 and 5.4 nm, respectively.

the 101 BZO peaks, there are no obvious differences between the current Hf and the previously reported Zr samples.

Figure 4 shows a cross-section TEM micrograph of the present Hf sample. The average BHO nanorod diameter is 5.4 nm, which is larger than the 3.7 nm diameter found in the Zr sample [16]. On the other hand, the average spacing between nanorods has been estimated to be 18 nm, which is near-identical to the 19 nm average spacing found in the 15% Zr sample. The 18 nm spacing corresponds to a matching field of  $\sim 6.4$  T, which agrees well with the observed saturation field for the pinning force shown in figure 2. Regarding the REO precipitates, unlike the previously reported Zr sample and in accordance with the 2D-XRD results, there are no observable REO precipitates in the shown micrograph. We did observe some REO precipitates in this sample by TEM, but they have been found to be scarce compared to the TEM observations for the Zr sample [16].

#### 4. Summary

An optimized REBCO sample with 15% mol. Hf addition and 4.0  $\mu\text{m}$  thickness was produced using A-MOCVD and measured for  $I_c$  performance at 4.2 K both in  $B\parallel ab$  and  $B\parallel c$  orientations up to a maximum field of 31.2 T. The critical current at  $B\parallel ab$  reached remarkable values of 6300 and 5812 A/4 mm-width at 14 and 30 T, respectively, which are 3 and 5.7 times higher than the corresponding measured  $I_c(B\parallel c)$  values of 2114 and 1013 A/4 mm, respectively. The corresponding critical current density values are 40.5 and 36.3 MA  $\text{cm}^{-2}$  at 14 and 30 T, respectively. The critical current density of the present 15% Hf sample in the  $B\parallel c$  field orientation has been found to be near-identical to the values of the previously reported 4.3 and 4.6  $\mu\text{m}$  thick REBCO films with 15% Zr addition [16]. The engineering current density performance at  $B\parallel ab$ , 15 T is over 16 times higher than that of  $\text{Nb}_3\text{Sn}$ . The pinning force at the  $B\parallel ab$  orientation exceeds 10 TN  $\text{m}^{-3}$  above 30 T, reaching 11.5 TN  $\text{m}^{-3}$  at 31.2 T.

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