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Application of High-Resolution ⁹⁵Mo Solid State NMR Spectroscopy in Structural Studies of Complex Alkali Molybdate Crystals and Glasses

Sabyasachi Sen,* Ivan Hung, Jacob M. Lovi, and Zhehong Gan



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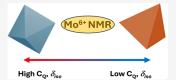


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ABSTRACT: The Mo–O coordination environment is probed in a series of simple and complex crystalline alkali molybdates as well as in mixed-alkali molybdate glasses using high-field (18.8 and 20.0 T) 95 Mo magic-angle-spinning (MAS) and multiple-quantum MAS (MQMAS) nuclear magnetic resonance (NMR) spectroscopy. When taken together, the 95 Mo NMR spectroscopic results indicate that somewhat contrary to the conventional wisdom the corner- and edge-shared MoO₆ octahedral sites in these alkali molybdates are characterized by higher values of the isotropic



shift (δ_{iso}) and quadrupolar coupling constant (C_Q) compared to the MoO₄ tetrahedral sites. These trends are hypothesized to be related to the unusually strong distortion of MoO₆ octahedra in corner- and edge-shared configurations and the resulting increase in the paramagnetic component of the chemical shift. While the ⁹⁵Mo C_Q of the MoO₄ sites displays an approximately linear positive correlation with the degree of tetrahedral distortion, no such correlation is observed for the MoO₆ sites. High-resolution ⁹⁵Mo NMR spectra show the coexistence of tetrahedral and octahedral Mo-O environments in the structure of alkali molybdate glasses, with the relative fraction of the latter environment increasing with Mo content. The results presented in this study indicate that high-resolution ⁹⁵Mo NMR spectroscopy at high magnetic fields (\sim 20 T or higher) may prove to be a promising tool for investigating the Mo-O coordination environments in nuclear waste glasses.

1. INTRODUCTION

Molybdenum (Mo) is an early transition metal that has found a wide range of technological applications. Mo is used as an alloying agent in steel and other metals to improve mechanical strength and resistance to corrosion and wear. On the other hand, several crystalline oxide and chalcogenide compounds of Mo with high catalytic activity are utilized in various chemical processes.^{2,3} It may be noted here that glasses can also be made with MoO₃ as a major constituent where, however, the glassformation requires addition of one or more alkali oxides.⁴⁻⁷ Mo is present as Mo⁶⁺ in these alkali molybdate glasses, where recent Raman spectroscopic studies have indicated that the structural network consists of a mixture of corner and edgeshared MoO₄ and MoO₆ polyhedra. In contrast to molybdate glasses where it serves as a major component, MoO₃ is often incorporated as a minor component in silicate, borate and phosphate glasses as a nucleating agent for synthesis of glassceramics as well for photochromic, electrochromic and gas sensing applications where the redox behavior of Mo provides the desired functionality.8-11

Inorganic molybdates display rich structural chemistry where Mo is often present in multiple coordination states, predominantly as a mixture of MoO_4 tetrahedra and MoO_6 octahedra with a wide range of connectivity. Despite its unique ability to quantitatively investigate multiple coordination states of various elements in crystal and glass structure, the application of solid-state nuclear magnetic resonance (NMR)

spectroscopy of Mo has remained quite limited in the literature, owing to the rather low gyromagnetic ratio γ (-1.751×10^7 rad T⁻¹ s⁻¹) and thus the low sensitivity, as well as the quadrupolar nature (I=5/2) of the NMR active ⁹⁵Mo and ⁹⁷Mo nuclides. Besides, the low γ of these two nuclides also poses a problem with efficient generation of multiple quantum coherences needed for high-resolution multiple quantum magic angle spinning (MQMAS) NMR spectroscopy, ¹⁸ often desirable for acquiring high-resolution spectra of quadrupolar nuclides free of quadrupolar broadening interactions.

Performing NMR spectroscopy at high magnetic fields (\geq 14.1 T) is often the most effective way to address the problems associated with low- γ nuclides such as 95 Mo. A few such studies have indeed been reported in the literature over the last two decades. Most interestingly, these studies indicated the lack of any clear correlation between the 95 Mo isotropic chemical shift $\delta_{\rm iso}$ and the various structural parameters associated with the Mo–O coordination environments such as the Mo–O bond length and O–Mo–O bond

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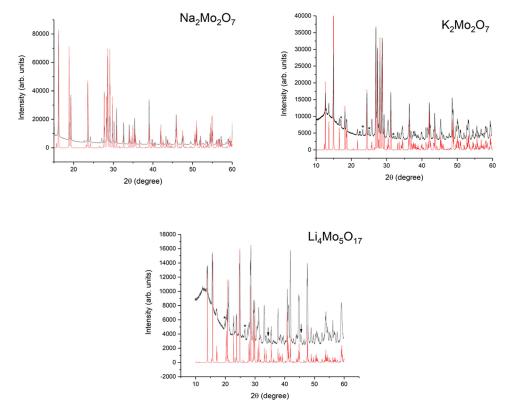


Figure 1. Experimental (black line) PXRD patterns for $Na_2Mo_2O_7$, $K_2Mo_2O_7$, and $Li_4Mo_5O_{17}$ crystals synthesized in this study. Corresponding theoretical patterns are shown as red lines. Peaks in the experimental patterns marked by asterisks and arrows are from secondary phases.

angles and their distributions as well as the Mo–O coordination numbers in these compounds, even within a given chemical series, namely the alkali molybdates. 18,19 On the other hand, a positive correlation was reported between the quadrupolar coupling constant $C_{\rm Q}$ of $^{95}{\rm Mo}$ and the bond angle distortion factor σ for the isolated $[{\rm MoO_4}]^{-2}$ tetrahedra in alkali orthomolybdates. 19 The distortion factor represents the root-mean-square deviation of these tetrahedra from $T_{\rm d}$ symmetry in terms of the root-mean-square departure of the O–Mo–O bond angles from the ideal value of 109.47° and is defined as: 19

$$\sigma = \sqrt{\frac{1}{6} \sum_{i \neq j=1}^{6} \left[(\angle O_i - Mo - O_j) - 109.47^{\circ} \right]^2}$$

A similar definition for σ can be obtained for MoO₆ octahedra in terms of the departure of the 12 O-Mo-O angles between the cis oxygen atoms from 90°. 18 However, these pioneering studies primarily focused on alkali and alkaline-earth orthomolybdate compounds, where Mo is present in the structure as isolated and somewhat regular tetrahedral $[MoO_4]^{-2}$ anions. Only a handful of structures with Mo solely in the octahedral environment $[MoO_6]^{-6}$ have been investigated with 95Mo NMR and practically no systematic NMR study of molybdates with mixed coordination states exist in the literature. 18,19 Here we utilize high-field (18.8 and 20.0 T) 95Mo MAS and MOMAS NMR spectroscopy in combination with density functional theory (DFT) based calculation of 95Mo quadrupolar parameters of a series of alkali molybdate crystalline compounds with Mo in coexisting tetrahedral and octahedral coordination environments with complex connectivity patterns and consequently large distortion parameters to address this gap in knowledge. It may be noted that for central transition spectra of half-integer quadrupolar nuclides such as ⁹⁵Mo, broadened by second-order quadrupolar effect, an increase in magnetic field results in a lowering of quadrupolar line broadening effects and an increase in resolution and sensitivity proportional to the square of the magnetic field in ppm scale. We also report preliminary ⁹⁵Mo NMR results for select alkali molybdate glasses to show the feasibility of this spectroscopic technique in investigating the Mo–O coordination environments in disordered solids.

2. EXPERIMENTAL DETAILS

2.1. Sample Synthesis. The crystalline orthomolybdate phases of composition A₂MoO₄ (A = Li, Na, and K) were purchased from Aldrich (99.9%) and their phase purity was confirmed using powder X-ray diffraction (PXRD). These materials were used as received for 95Mo measurements. The pyromolybdates Na₂Mo₂O₇ and K₂Mo₂O₇ as well as the polymolybdate phase Li₄Mo₅O₁₇ were synthesized from the corresponding melts. Stoichiometric mixtures of alkali carbonate and MoO₃ reagents (Li₂CO₃, Alfa Aesar, 99.999%; Na₂CO₃, EMD, 99.5%; K₂CO₃, Alfa Aesar, 99.997%; MoO₃, Thermo Scientific, 99.5%) were taken in silica crucible and melted at temperatures ranging between 730 and 800 °C for 20 min. The melts were transferred to another furnace and held at 420, 450, and 550 °C for 20 h, respectively, for Li₄Mo₅O₁₇, K₂Mo₂O₇, and Na₂Mo₂O₇ followed by cooling down to ambient temperature by turning off the furnace. The orangish yellow melts were observed to turn to opaque white crystalline material. The phase purity of the resulting materials was checked with PXRD (Figure 1). Previous studies of the

phase relations in the $\text{Li}_2\text{MoO}_4\text{--MoO}_3$ system indicate the presence of only two compounds: $\text{Li}_4\text{Mo}_5\text{O}_{17}$ and $\text{Li}_2\text{Mo}_4\text{O}_{13}$. The $\text{Na}_2\text{Mo}_2\text{O}_7$ sample was found to be phase-pure, while the PXRD patterns of the $\text{K}_2\text{Mo}_2\text{O}_7$ and $\text{Li}_4\text{Mo}_5\text{O}_{17}$ phases showed the presence of small amounts of additional phases that were formed during the synthesis (Figure 1). The identity/chemistry of these secondary phases, including a small amount of residual amorphous phase, remains unclear at this point. However, the relative concentration of these secondary phases, as estimated from the PXRD patterns turns out to be rather small (5–10%) and, as discussed below, their presence does not affect the interpretation of the ^{95}Mo NMR spectra of the primary phases.

The two mixed-alkali molybdate glasses studied here were prepared using the conventional melt-quenching method. The compositions of these two glasses are 60 and 62.5 mol % MoO₃ with the remainder being a mixture of a Li₂O, Na₂O, and K₂O in a 1:1:2 molar ratio. These glasses are referred to as Mo-60 and Mo-62.5 in the subsequent discussion. Details of their synthesis can be found in a recent publication. Appropriate stoichiometric mixtures of the above-mentioned alkali carbonate and MoO3 reagents were mixed together and melted in quartz crucible in air at 730 °C for 20 min. The melts were poured onto a graphite plate and quenched by pressing with an aluminum plate. PXRD on a part of the quenched glass samples indicated the lack of any detectable crystallinity. The rest of the samples were immediately transferred and stored in a vacuum desiccator until further measurements.

2.2. 95 Mo NMR Spectroscopy. 2.2.1. 95 Mo MAS NMR at 18.8 T with Cryoprobe. The 95Mo MAS NMR spectra of the Na₂Mo₂O₇ crystal, the Na₂MoO₄ crystal and the Mo-60 glass were acquired at a magnetic field of 18.8 T using a Bruker Avance NEO console operating at a 95Mo Larmor frequency of 52.1 MHz, and a 3.2 mm double-resonance Bruker CPMAS cryoprobe. The use of the cryoprobe resulted in excellent signal:noise ratio within a relatively short acquisition time. Crushed samples were packed into ZrO2 rotors and spun at 6 kHz for the Na₂MoO₄ crystal and at 15 kHz for the Na₂Mo₂O₇ crystal and the Mo-60 glass samples. The single-pulse MAS NMR spectrum of the Na₂MoO₄ crystal was acquired using only a single scan with a $\pi/2$ pulse of duration 1.5 μ s to obtain a signal:noise ratio of 250:1. Additionally, a saturation recovery experiment on Na₂MoO₄ yielded a relatively long ⁹⁵Mo spinlattice relaxation time T_1 of ~68 s, consistent with the small quadrupole moment and the cubic environment (see below) of this nuclide in this crystal. The ⁹⁵Mo MAS NMR spectra of the Na₂Mo₂O₇ crystal and the Mo-60 glass were acquired using a Hahn echo pulse sequence $(\pi/2-\tau-\pi$ -acquisition) with a $\pi/2$ pulse length of 5.5 μ s and τ = 63 μ s and recycle delays of 20 and 5 s for the crystal and the glass, respectively. A total of 16 scans for the crystal and 256 scans for the glass were averaged and Fourier-transformed to obtain each spectrum.

2.2.2. ⁹⁵Mo MAS and MQMAS at 20.0 T. The ⁹⁵Mo MAS and MQMAS NMR spectra of the Li₂MoO₄, K₂MoO₄, K₂Mo₂O₇, and Li₄Mo₅O₁₇ crystals and the Mo-62.5 glass were acquired at the National High Magnetic Field Laboratory (NHMFL, Tallahassee, Florida) using a Bruker Avance NEO console operating at a ⁹⁵Mo Larmor frequency of 55.42 MHz and a Low-E 3.2 mm HX MAS probe designed and constructed at the NHMFL. Crushed glass samples were packed into 3.2 mm ZrO₂ rotors and spun at 16 kHz. The 1D spectra of Li₂MoO₄, K₂MoO₄, and K₂Mo₂O₇ were acquired

using a spin–echo experiment with satellite-transition saturation/inversion, 24 employing a 2 ms WURST pulse at an offset of +272,000 kHz and rf amplitude of 12.5 kHz. The QCPMG pulse sequence was used instead of the spin–echo for the Li₄Mo₅O₁₇ crystal and Mo-62.5 glass samples for additional signal enhancement, along with the same WURST enhancement used for the other samples. $\pi/2$ - and π -pulses of 4.5 and 9.0 μs were used at a rf amplitude of 18.5 kHz. The recycle delay and the number of averaged transients for each sample were: (K₂MoO₄, 65 s, 704), (Li₂MoO₄, 65 s, 1024), (K₂Mo₂O₇, 40 s, 2048), and (Li₄Mo₅O₁₇, 5 s, 2320).

The 2D spectra of $K_2Mo_2O_7$, $Li_2Mo_2O_7$ and Mo-62.5 employed a combination of triple-quantum (3Q) shifted-echo split- t_1 SPAM MQMAS^{26–29} and QCPMG (Figure 2). Pulses

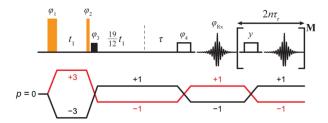


Figure 2. MQMAS/QCPMG pulse sequence used in this study with cogwheel phase cycle [26]: $\varphi_1 = 17k \cdot (2\pi/66)$, $\varphi_2 = \varphi_3 = 12k \cdot (2\pi/66)$, $\varphi_4 = 15k \cdot (2\pi/66)$, $\varphi_{Rx} = k \cdot \pi$, where k = 0, 1, 2,..., 65.

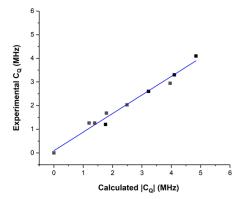
of 4.95 and 1.65 μ s were used for excitation and conversion of 3Q coherence with a rf amplitude of ~80 kHz. The indirect spectral window was rotor-synchronized by incrementing the first spli- t_1 delay by the rotor period, 62.5 μ s. Hypercomplex acquisition was accomplished by incrementing the phase of the 3Q excitation pulse by 30°, as described by States et al. 30 The recycle delay, number of averaged transients, number of complex points acquired in the indirect dimension, and total experiment time for each sample were: ($K_2Mo_2O_7$, 20 s, 66, 32, 23.5 h), ($Li_4Mo_3O_{17}$, 5 s, 660, 32, 58.7 h), and (Mo-62.5, 3 s, 1518, 8, 20.2 h). In all cases using QCPMG acquisition, the echoes in the echo trains were summed prior to processing of the spectra. The ^{95}Mo chemical shift for all spectra reported here were referenced to the isotropic shift of crystalline Na_2MoO_4 at -33.5 ppm. 18

2.3. DFT-Based Calculation of 95Mo Quadrupolar Parameters. The 95Mo NMR quadrupolar parameters, i.e., C_0 and the asymmetry parameter of the electric field gradient (EFG) tensor η for crystalline orthomolybdates (Li₂MoO₄, K₂MoO₄, Na₂MoO₄, BaMoO₄, PbMoO₄, and CaMoO₄); alkali pyromolybdates (Na₂Mo₂O₇ and K₂Mo₂O₇) and Li polymolybdate (Li₄Mo₅O₁₇) were calculated using the gaugeincluding projector augmented wave (GIPAW) method based on DFT. 31-33 The GIPAW method allows the reconstruction of the all-electron magnetic response from the pseudowave functions and is ideally suited for extended systems with periodic boundary conditions. Several past studies have demonstrated the excellent accuracy of the GIPAW method in calculating the NMR parameters for a wide variety of nuclides in crystalline and amorphous systems.³³ The NMR calculations in the present study were carried out using the code CASTEP-NMR (Biovia Inc.) within the generalized gradient approximation with Perdew-Burke-Ernzerhof (PBE) XC functionals using on-the-fly generated ultrasoft pseudopotentials and plane wave basis sets with an energy cutoff of 610 eV, where the total energy converged within 10^{-7} eV/atom.

Table 1. Comparison between Measured and Calculated (by DFT) ⁹⁵Mo Quadrupolar Parameters for Crystalline Orthomolybdate and Pyromolybdate Compounds

composition	Mo site	$C_{\rm Q}$ (measured) \pm 0.05 MHz	$C_{\rm Q}$ (MHz) (calculated)	η (measured) \pm 0.05	η (calculated)	structural reference
Na_2MoO_4	tetrahedral	0.0	0.0			
Li_2MoO_4	tetrahedral	1.25	1.39	0.65	0.72	
K_2MoO_4	tetrahedral	1.26	1.20	0.65	0.66	
BaMoO ₄ ^a	tetrahedral	1.68	1.79	0.10	0.00	
PbMoO ₄ ^a	tetrahedral	2.03	2.49	0.00	0.00	
CaMoO ₄ ^a	tetrahedral	2.94	3.96	0.00	0.00	
$Na_2Mo_2O_7$	tetrahedral	1.20	1.76	0.70	0.60	
	octahedral	4.10	4.84	0.80	0.75	
$K_2Mo_2O_7$	tetrahedral	2.60	3.20	0.30	0.20	
	octahedral	3.30	4.10	0.65	0.56	
$\text{Li}_4\text{Mo}_5\text{O}_{17}$	octahedral		5.76		0.39	
	octahedral		5.61		0.41	
	octahedral		5.66		0.35	
	octahedral		4.31		0.95	
	octahedral		4.35		0.49	
a Evnorimental		rof 18	4.33		0.49	

^aExperimental data are from ref 18.



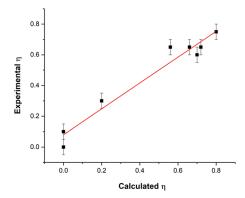


Figure 3. Comparison between Experimental and DFT-calculated values of 95 Mo C_Q (left) and η (right) of MoO₄ and MoO₆ sites in various molybdate compounds listed in Table 1. Straight lines through data points represent linear least-squares fits given by eqs 3 and 4 in the text.

The Brillouin zone was sampled using the Monkhorst–Pack scheme and a $4 \times 4 \times 4$ k-point grid. The EFG tensor was calculated from the second spatial derivative of the electrostatic potential resulting from the charge density in the crystal. The principal components of this tensor $V_{zz} \geq V_{yy} \geq V_{xx}$ yield the C_Q and η parameters as follows:

$$C_{Q} = eQV_{zz}/h \tag{1}$$

$$\eta = \frac{V_{xx} - V_{yy}}{V_{zz}} \tag{2}$$

In eq 1 e is the electronic charge, h is Planck's constant, and Q is the quadrupole moment of ^{95}Mo (-22 mB). Calculations were carried out on molybdate crystal structures reported in the literature, with and without geometry optimization. In all cases geometry optimization did not seem to have a significant effect on the resulting quadrupolar parameters. Therefore, here we report the DFT results on unoptimized structures (Table 1). Preliminary attempts of calculating ^{95}Mo magnetic shielding tensors for these crystalline compounds yielded unreliable results at this level of DFT without the consideration of relativistic effects and were not pursued further in the present study. It is to be noted that previous studies in the literature 34 demonstrated the importance of the inclusion of relativistic effects in the DFT calculations for

accurate estimation of the large paramagnetic component of ⁹⁵Mo chemical shift that arises from a symmetry-allowed mixing of occupied and virtual d-orbitals of this transition metal.

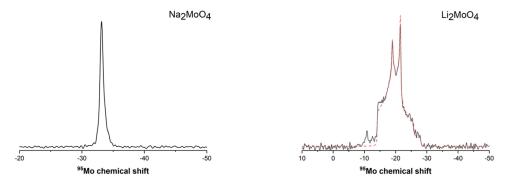
The $C_{\rm Q}$ and η values obtained for the compounds BaMoO₄, PbMoO₄, and CaMoO₄ are found to be in good agreement with those reported in a previous study where the calculations were carried out employing the linearized augmented plane wave + local orbitals (L/APW+lo) method using the code WIEN2k. A comparison between the experimental and DFT-calculated $|C_{\rm Q}|$ and η values indicates approximate linear correlation with 15–20% overprediction (Figure 3) that can be described using the following relationships ($R^2 \sim 0.95$):

$$C_{\rm Q}^{\rm expt} = 0.80 \times C_{\rm Q}^{\rm calc} + 0.09$$
 (3)

and

$$\eta^{\text{exp}t} = 0.84 \times \eta^{\text{calc}} + 0.08 \tag{4}$$

The superscripts expt and calc in eqs 3 and 4 correspond, respectively, to experimental and calculated values of $|C_Q|$ and η . It is interesting to note that quite similar functional relationships as in eqs 3 and 4 were also obtained in a recent study of ⁹⁵Mo NMR parameters on a wide range of molybdates, where the authors utilized a higher cutoff energy of 700 eV for the plane wave basis sets and a more stringent



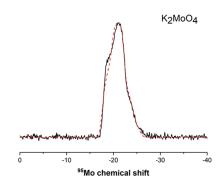


Figure 4. Experimental ⁹⁵Mo MAS NMR spectra (black solid lines) of alkali orthomolybdate crystals. Spectrum of Na₂MoO₄ was collected at 18.8 T while the spectra of Li₂MoO₄ and K₂MoO₄ were collected at 20.0 T. Simulated spectra of K₂MoO₄ and Li₂MoO₄ are shown as red dashed lines.

Table 2. ⁹⁵Mo NMR Parameters for Octahedral Mo Environments in Crystalline Alkali Orthomolybdate and Pyromolybdate Compounds Obtained from Simulation of ⁹⁵Mo MAS NMR Spectra Acquired in This Study

Mo site	Mo−O range (average) Å	σ (°)	$\delta_{\rm iso}~(\pm 0.5~{ m ppm})$	$C_{\rm Q}~(\pm 0.05~{ m MHz})$	$\eta \ (\pm 0.05)$
tetrahedral	1.788	0.00	-33.5	0.0	
tetrahedral	1.759-1.769 (1.76)	1.40	-26.0	1.25	0.65
tetrahedral	1.742-1.785 (1.76)	1.68	-17.5	1.26	0.65
tetrahedral	1.708-1.784 (1.76)	1.46	-25.0	1.20	0.70
octahedral	1.685-2.267 (1.95)	9.76	-11.0	4.10	0.80
tetrahedral	1.701-1.837 (1.76)	2.78	-20.0	2.60	0.30
octahedral	1.725-2.262 (1.97)	10.58	47.0	3.30	0.65
	tetrahedral tetrahedral tetrahedral tetrahedral octahedral tetrahedral	tetrahedral 1.788 tetrahedral 1.759–1.769 (1.76) tetrahedral 1.742–1.785 (1.76) tetrahedral 1.708–1.784 (1.76) octahedral 1.685–2.267 (1.95) tetrahedral 1.701–1.837 (1.76)	tetrahedral 1.788 0.00 tetrahedral 1.759–1.769 (1.76) 1.40 tetrahedral 1.742–1.785 (1.76) 1.68 tetrahedral 1.708–1.784 (1.76) 1.46 octahedral 1.685–2.267 (1.95) 9.76 tetrahedral 1.701–1.837 (1.76) 2.78	tetrahedral 1.788 0.00 -33.5 tetrahedral 1.759-1.769 (1.76) 1.40 -26.0 tetrahedral 1.742-1.785 (1.76) 1.68 -17.5 tetrahedral 1.708-1.784 (1.76) 1.46 -25.0 octahedral 1.685-2.267 (1.95) 9.76 -11.0 tetrahedral 1.701-1.837 (1.76) 2.78 -20.0	tetrahedral 1.788 0.00 -33.5 0.0 tetrahedral 1.759-1.769 (1.76) 1.40 -26.0 1.25 tetrahedral 1.742-1.785 (1.76) 1.68 -17.5 1.26 tetrahedral 1.708-1.784 (1.76) 1.46 -25.0 1.20 octahedral 1.685-2.267 (1.95) 9.76 -11.0 4.10 tetrahedral 1.701-1.837 (1.76) 2.78 -20.0 2.60



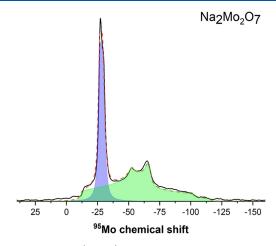
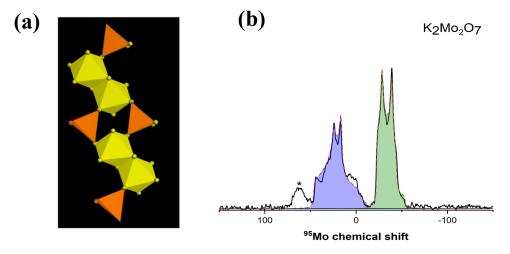


Figure 5. Left: Connectivity of MoO_6 octahedra (teal) and MoO_4 tetrahedra (brown) in the crystal structure of $Na_2Mo_2O_7$. Right: Experimental (black solid line) and simulated (red dashed lines) ⁹⁵Mo MAS NMR spectra of $Na_2Mo_2O_7$ crystal. Individual simulation components are shown as blue and green shaded peaks.



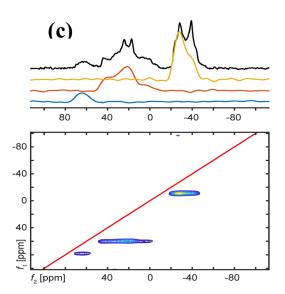


Figure 6. (a) Connectivity of MoO_6 octahedra (yellow) and MoO_4 tetrahedra (orange) in the crystal structure of $K_2Mo_2O_7$. (b) Experimental (black solid line) and simulated (red dashed lines) ^{95}Mo MAS NMR spectra of $K_2Mo_2O_7$ crystal. Individual simulation components are shown as blue and green shaded peaks. Peak with asterisk corresponds to a secondary phase. (c) Contour plot of ^{95}Mo MQMAS NMR spectrum of $K_2Mo_2O_7$ crystal. A comparison between MAS spectrum (black line) and MAS projections of the individual resonances in the MQMAS spectrum is shown on top.

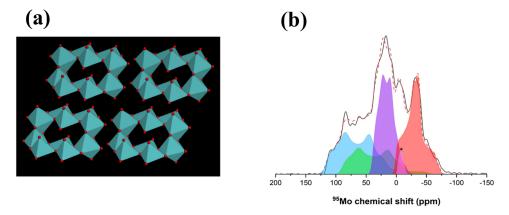
energy convergence criterion of 10^{-8} eV/atom as implemented in the code VASP. 35,36

3. RESULTS AND DISCUSSION

3.1. Alkali Orthomolybdates. The ^{95}Mo MAS NMR spectra of three $A_2\text{MoO}_4$ phases (A = Li, Na and K) were collected to check for consistency with previously reported data in the literature (Figure 4). The values of δ_{iso} , C_{Q} , and η obtained from simulation of these spectral line shapes carried out using the software Dmfit (Figure 4) are reported in Table 2. These values are found to be consistent with those reported in previous studies based on simulation of spectra collected at magnetic fields ranging between and 11.7 and 19.6 T. $^{18-20}$ As noted earlier, the MoO $_4$ tetrahedra in these compounds are rather regular with a small spread in Mo–O bond distances and Mo–O–Mo angles. The perfect $T_{\rm d}$ symmetry of the MoO $_4$ tetrahedron in Na $_2$ MoO $_4$ resulted in a symmetric line

shape with negligible quadrupolar broadening, i.e., $C_{\rm Q}\approx 0$ MHz. ^{18,19} The MoO₄ tetrahedra in Li₂MoO₄ and K₂MoO₄ show small departure from tetrahedral symmetry ^{38,39} with similar bond angle distortion parameters ($\sigma\sim 1.40-1.68$), which results in relatively low $C_{\rm Q}$ values for these Mo sites (1.25–1.26 MHz, see Table 2). The average Mo–O bond distances for the Na, Li, and K phases are 1.788, 1.764, and 1.762 Å, respectively. On the other hand, the $\delta_{\rm iso}$ for the Na, Li, and K phases are –33.5, –26.0, and –17.5 ppm, respectively, implying the lack of any clear correlation between $\delta_{\rm iso}$ and Mo–O bond distance (Table 2), as was noted in previous studies. ^{18,19}

3.2. Na and K Pyromolybdates. The structure of the Na₂Mo₂O₇ phase is characterized by chains of corner-shared MoO₆ octahedra along the *a*-axis where pairs of adjacent octahedra share corners with a MoO₄ tetrahedron, ⁴⁰ resulting in octahedral and tetrahedral Mo sites being present in 1:1



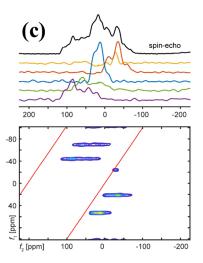


Figure 7. (a) Connectivity of MoO_6 octahedra in the structure of $Li_4Mo_5O_{17}$ crystal showing groups of $10\ MoO_6$ octahedra sharing edges to form ribbons along b direction. (b) Experimental (black solid line) and simulated (red dashed lines) ^{95}Mo MAS NMR spectra of $Li_4Mo_5O_{17}$ crystal. Individual simulation components are shown as colored shaded peaks. Peak with asterisk corresponds to a secondary phase. (c) Contour plot of ^{95}Mo MQMAS NMR spectrum of $Li_4Mo_5O_{17}$ crystal. A comparison between MAS spectrum (black line) and MAS projections of the individual resonances in the MQMAS spectrum is shown on top.

ratio (Figure 5). The Mo-O bond distances (1.71-1.78 Å with an average of 1.76 Å) and the angular distortion parameter σ (~1.46) for the MoO₄ tetrahedra in this structure are quite similar to those observed in orthomolybdates (Table 2). In contrast, the MoO₆ octahedra in this structure are highly distorted with Mo-O bond distances ranging between 1.69 and 2.27 Å (average Mo–O of 1.95 Å) and a rather large σ (~9.8). The experimental ⁹⁵Mo MAS NMR spectrum (Figure 5) displays two partially overlapping but clear quadrupolar powder patterns corresponding to these two environments. Simulation of this spectrum (Figure 5) yields for the narrower resonance: $\delta_{\rm iso}$ = -25 ppm, $C_{\rm Q}$ = 1.2 MHz, and η = 0.70, while for the broader resonance we obtained: $\delta_{iso} = -11$ ppm, $C_{iso} = -11$ 4.1 MHz, and $\eta = 0.80$ (Table 2). The peak areas suggest that these two Mo environments are indeed present in equal abundance in the structure of Na₂Mo₂O₇. It may be noted here that the narrower resonance was also observed in a previous study by Wren et al., 19 however, the relatively low magnetic fields of 9.4 and 14.1 T used in that study failed to clearly resolve the broader second resonance. The equal abundance of the two environments makes their identification somewhat challenging in the ⁹⁵Mo NMR spectrum (Figure 5). However, it may seem logical to assign the resonance at $\delta_{iso} = -11$ ppm,

characterized by the rather large $C_{\rm Q}$ to the MoO $_6$ environment, which is highly distorted with a large spread in the Mo–O distances and a high value of σ (Table 2). Moreover, the Mo–O spread and the σ value of the MoO $_4$ environment of Na $_2$ Mo $_2$ O $_7$ is consistent with the $C_{\rm Q}$ of the resonance centered at $\delta_{\rm iso}=-25$ ppm. Therefore, we assign the resonance at $\delta_{\rm iso}=-25$ ppm to the tetrahedral Mo site and the resonance at $\delta_{\rm iso}=-11$ ppm to the octahedral Mo site. This assignment is completely consistent with the results of the DFT calculations that yield a significantly higher $C_{\rm Q}$ value (3.3 MHz) for the octahedral environment compared to that (1.2 MHz) for the tetrahedral environment (Table 1).

The structure of the $K_2Mo_2O_7$ phase is characterized by pairs of edge-shared MoO_6 octahedra connected via cornersharing pairs of opposing MoO_4 tetrahedra forming chains along the b-axis, ⁴¹ again resulting in an equal abundance of octahedral and tetrahedral Mo sites (Figure 6a). The Mo–O bond distances in the MoO_4 tetrahedra vary over a wider range (\sim 1.70–1.84 Å) in this structure and display a significantly larger σ (\sim 2.78), compared to those characteristic of its Na analogue. The edge-shared octahedral environment in $K_2Mo_2O_7$ on the other hand, is characterized by Mo–O distances ranging between 1.73 and 2.26 Å and σ \sim 10.6, both

Table 3. ⁹⁵Mo NMR Parameters for Octahedral Mo Environments in Crystalline Li-Polymolybdate Compound Li₄Mo₅O₁₇, Obtained from Simulation of ⁹⁵Mo MAS and MQMAS Spectra Acquired in This Study

composition	Mo site	Mo-O range (average) Å	σ	$\delta_{\rm iso}~(\pm 0.5~{\rm ppm})$	$C_{\rm Q}~(\pm 0.05~{ m MHz})$	$\eta \ (\pm 0.05)$
Li ₄ Mo ₅ O ₁₇	octahedral octahedral octahedral octahedral	1.688-2.497 (1.95-1.97)	~9.5	124.0 99.4 40.3 2.1	5.38 5.55 3.37 3.73	0.39 0.31 0.50 0.79

being comparable to those observed for the MoO₆ site in the structure of Na₂Mo₂O₇. The ⁹⁵Mo MAS NMR spectrum of K₂Mo₂O₇ (Figure 6b) displays two well-resolved quadrupolar powder patterns corresponding to these two environments. Simulation of this spectrum yields for the narrower resonance: $\delta_{\rm iso}$ = -20 ppm, $C_{\rm Q}$ = 2.6 MHz, and η = 0.30, while the broader resonance can be simulated using $\delta_{\rm iso}$ = 47 ppm, $C_{\rm Q}$ = 3.3 MHz, and $\eta = 0.65$ (Table 2). The $C_{\rm O}$ of the narrower resonance is consistent with the σ for the MoO₄ tetrahedral sites in K₂Mo₂O₇ and falls on the extrapolated linear trend between C_Q and σ for the MoO₄ tetrahedral sites that was reported by Wren et al. ¹⁹ in a previous ⁹⁵Mo MAS NMR study of alkali molybdates. Similar to the spectrum of Na₂Mo₂O₇, we assign the resonance with smaller C_Q and η in the ⁹⁵Mo MAS NMR spectrum of $K_2Mo_2O_7$ to the tetrahedral site and the broader resonance to the octahedral site. Again, the DFT calculations are observed to be in complete agreement with this assignment, predicting a lower $C_{\rm O}$ (2.2 MHz) for the tetrahedral site, compared to that (2.8 MHz) for the octahedral site (Table 1). Finally, we note the presence of a relatively weak (~8%) signal centered at ~63 ppm in the ⁹⁵Mo MAS NMR spectrum of K₂Mo₂O₇ associated with the presence of a small amount of an additional phase formed during its synthesis, as was also observed in the X-ray diffraction pattern. The ⁹⁵Mo MQMAS/QCPMG NMR spectrum of K₂Mo₂O₇ is shown in Figure 6c. The Mo resonances of the tetrahedral and octahedral Mo sites are clearly resolved in the isotropic dimension, implying excellent feasibility of such experiments in future for high-resolution ⁹⁵Mo NMR in solids.

The result of a high-coordinated cation site in these alkali pyromolybdates being characterized by a higher value of δ_{iso} and C_O compared to a low-coordinated site may seem somewhat counterintuitive at first, since an opposite trend is typically observed for nearly all other nuclides in oxides (e.g., ²³Na, ²⁷Al, ¹¹B, ⁷Li, ²⁵ Mg, ²⁹Si, ³¹P, ⁴³Ca, ⁴⁵Sc, and even the heavy element ⁸⁹Y) in oxides. ⁴² A similar observation was also made in recent DFT calculations by Yasui et al. of NMR parameters of tetrahedral and octahedral Mo sites in the perovskite Ba7Nb4MoO20, where the latter sites are characterized by higher values of $\delta_{\rm iso}$ and $C_{\rm Q}$. However, before considering these trends in the ⁹⁵Mo NMR parameters to be anomalous, one needs to consider the unusually large spread of Mo-O distances with a bimodal distribution for the cornerand edge-shared MoO6 environments in these alkali pyromolybdate compounds. Typically, two of the six Mo-O bonds in these MoO6 octahedra are longer than 2.2 Å while the other four range between \sim 1.7 and 1.9 Å. 40,41 Therefore, these MoO₆ octahedra can be considered as MoO₄ tetrahedra, if an Mo-O distance of >2.0 Å is considered as the cutoff for bonding interactions. A similar situation exists for the crystal structure of molybdite (MoO₃), which can be described in terms of layers of corner- and edge- shared MoO₆ octahedra or alternatively, upon considering a cutoff distance of ≥ 2.25 Å for the Mo-O bonding interaction, this structure can also be

described as parallel chains of corner-shared MoO₄ tetrahedra along the *c*-axis. A similar cutoff distance leads to a description of the structure of Na₂Mo₂O₇ in terms of isolated MoO₄ tetrahedra and corner-sharing MoO₄ tetrahedra, and of the structure of K₂Mo₂O₇ in terms of isolated tetrahedra and dimers. The possible effect of such distortion on the 95 Mo $\delta_{\rm iso}$ is discussed in more detail in the following section.

3.3. Li Polymolybdate Li₄Mo₅O₁₇. The structure of the polymolybdate phase Li₄Mo₅O₁₇ is described as a complex network of highly distorted edge-shared MoO₆ octahedra that are interlinked via LiO₄ and LiO₆ octahedra. ^{23,44,45} In particular, the structure can be considered to consist of groups of 10 MoO₆ octahedra sharing edges to form ribbons (Figure 7a). These ribbons are linked together by LiO₄ and LiO₆ polyhedra. The 95Mo MAS NMR spectrum of Li₄Mo₅O₁₇ (Figure 7b) displays a line shape consisting of strongly overlapping signals from multiple Mo sites. The 95Mo MQMAS/QCPMG NMR spectrum of this compound (Figure 7c) shows clear resolution of at least four Mo sites in the isotropic dimension along with a relatively weak resonance corresponding likely to small amounts of a secondary phase in this sample. The δ_{iso} and the quadrupolar product $P_Q = C_Q$ (1 + $\eta^2/3$) for these resonances for a spin I = 5/2 nuclide such as 95 Mo can be estimated from their isotropic peak position δ_{F_1} and their center of gravity $\delta_{F_2}^{\rm CG}$ along the MAS dimension in the MQMAS spectrum using the following relations:

$$\delta_{\rm iso} = \frac{17}{27} \delta_{F_1} + \frac{10}{27} \delta_{F_2}^{\rm CG} \tag{5}$$

$$P_{Q} = \left[\left(\delta_{F_{1}} - \delta_{F_{2}}^{CG} \right)^{1/2} \right] \left(\frac{\nu_{0}}{1000} \right) \sqrt{\frac{680[2I(2I-1)]^{2}}{81[4I(I+1)-3]}}$$
 (6)

where ν_0 is the resonance frequency (55.4 MHz) and I=5/2 is the spin quantum number of the nuclide. Using these relations, one can obtain the $\delta_{\rm iso}$ and $P_{\rm Q}$ values for the four resolved 95 Mo resonances in the MQMAS spectrum corresponding to the ${\rm Li_4Mo_5O_{17}}$ phase, which are then used to simulate the MAS line shape as shown in Figure 7b. The resulting 95 Mo NMR parameters are listed in Table 3.

Structural refinement based on X-ray diffraction suggests that Mo atoms in this compound are present in five crystallographically different environments with Mo–O distances ranging between ~1.688 and 2.497 Å and average Mo–O distances in these octahedra ranging between 1.96 and 1.98 Å. These values are similar to those reported for the MoO₆ octahedra in Na and K-pyromolybdates (Table 2). It may be noted that the Mo–O distances for at least two of these 5 Mo sites (Mo2 and Mo5) are rather similar. Moreover, an Mo–O bonding cutoff distance of ~2.2 would render several of these sites 5-fold or even 4-fold coordinated to oxygen. Considering this complexity in the bonding Mo–O environment, we do not attempt any structural assignment of the four resonances in the ⁹⁵Mo NMR spectra. Rather we

simply note that 95 Mo MQMAS NMR can be highly effective in resolving strongly overlapping signal in regular MAS spectra from multiple crystallographically similar local environments in a structurally complex compound. Nevertheless, the quadrupolar parameters $C_{\rm Q}$ and η for these Mo sites as calculated using DFT (Table 1) indeed agree reasonably well with those obtained from the simulation of the 95 Mo MAS NMR line shape (Figure 7b, Table3).

As noted earlier, previous high resolution 95Mo MAS NMR spectroscopic studies of molybdate compounds failed to observe any clear correlation between the 95 Mo $\delta_{\rm iso}$ and the Mo–O bond length or coordination number. However, within a limited set of compounds such as the single-alkali orthomolybdates, the $^{95}\text{Mo}^{1}\delta_{\text{iso}}$ for the MoO₄ tetrahedral environment was found to vary within a relatively narrow range of -17 to -34 ppm. The $\delta_{\rm iso}$ for the MoO₄ sites in the Na and K pyromolybdates studied here is observed to fall within this range as well. Interestingly, this chemical shift range for tetrahedral Mo extends further to lower frequencies when one includes the reported 95 Mo δ_{iso} of -125 ppm for the MoO₄ site in the mixed-alkali orthomolybdate CsLiMoO₄. 46 This chemical shift range for tetrahedral Mo is consistent the observed ⁹⁵Mo δ_{iso} of -73 ppm for molybdite (MoO₃), if one considers the Mo to be four-coordinated in this compound. On the other hand, the 95 Mo δ_{iso} for the MoO₆ and Mo_{5+ δ} sites with similar Mo-O bond length ranges (~1.7 to 2.3 Å) and average Mo-O distances (1.95-1.98 Å) in Na- and Kpyromolybdates and in the Li-polymolybdate phases are higher and range between -11 and 124 ppm (Tables 2 and 3). When taken together, these NMR results on a variety of alkali molybdates suggest that the $Mo_{5+\delta}$ and MoO_6 sites are characterized by higher values of $\delta_{\rm iso}$ than those characteristic of the MoO₄ sites. A similar observation has also been made in recent ⁹⁵Mo NMR spectroscopic studies and DFT calculations by Yasui et al. and Tansho et al. on Ba-Nb-W-molybdate perovskites. 35,36 The shielding of a nuclear spin and hence, its chemical shift is a sum of a diamagnetic and a paramagnetic component. The diamagnetic component is typically larger and decreases the chemical shift but is controlled by the core electrons, while the paramagnetic component increases the chemical shift. As the latter is controlled by the valence electrons, the paramagnetic shift is most sensitive to the bonding environment of the nuclide. In the vast majority of cases, an increase in the coordination number results in a lengthening of the metal-ligand bond (e.g., metal-oxygen M-O bonds in an oxide), which lowers the electron density at the metal nuclide and consequently the chemical shift decreases (or shielding increases). The opposite trend observed in the present study for 95Mo in alkali molybdates may be explained as follows. The Mo⁶⁺ (d⁰ electronic state) in MoO₆ octahedra is known to be somewhat unique among transition metals in the sense that these octahedra undergo large distortions, often characterized by a trimodal Mo-O bond length distribution centered at ~1.7, 1.9, and 2.3 Å, when Mo-O-Mo linking is present via corner or edge- sharing between MoO₆ polyhedra.⁴⁷ Such large distortions have been shown to arise from both bond topological constraints associated with a lowering of crystallographic symmetry as well as from electronic effects such as the pseudo- or secondorder Jahn—Teller distortion [ref]. The latter leads to a strong overlap and mixing between oxygen 2p orbitals at the top of the valence band (HOMO) and the transition metal t_{2g} orbitals at the bottom of the conduction band (LUMO), $^{47-49}$

which is expected to amplify the paramagnetic component of the chemical shift and thus may increase the $\delta_{\rm iso}$ of MoO₆ sites. This increase in $\delta_{\rm iso}$ could be further amplified due to a distortion-induced increase in the chemical shift anisotropy if the effect of the paramagnetic deshielding is the strongest for the least shielded component. Finally, we note that similar to the observation made by Wren et al.¹⁹ in a previous study on a limited set of alkali molybdates, an approximately linear increase in $C_{\rm Q}$ with σ is observed in the present study for the MoO₄ tetrahedral sites in both alkali ortho- and pyromolybdates (Figure 8). Moreover, MoO₄ sites in other

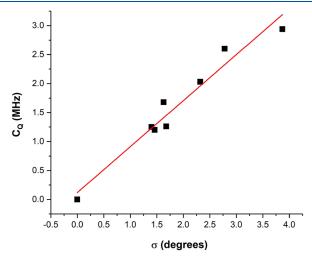
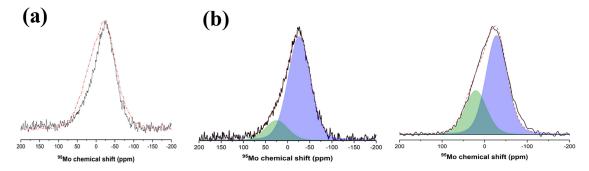


Figure 8. Variation of ⁹⁵Mo $C_{\rm Q}$ with tetrahedral distortion parameter σ for MoO₄ sites in ortho- and pyro-molybdate compounds listed in Table 1. Straight line through the data points is a linear least-squares fit given by the relation: $C_{\rm Q} = 0.79 \times \sigma + 0.12$.

orthomolybdates such as BaMoO₄, PbMoO₄, and CaMoO₄ appear to obey the same functional relationship between $C_{\rm Q}$ and σ (Figure 8). However, no clear correlation between these two quantities is observed for the high-coordinated Mo sites in alkali pyro- and poly-molybdates studied here (Tables 2 and 3)

3.4. Mixed-Alkali Molybdate Glasses. The 95Mo MAS NMR spectra of the two alkali molybdate glasses Mo-60 and Mo-62.5 are shown in Figure 9a,b. These spectra display a rather broad and asymmetric line shape with a tail on the high frequency side (higher ppm values). The 95Mo MQMAS/ QCPMG NMR of the Mo-62.5 glass clearly indicate the presence of at least two resonances (Figure 9c). The orientation of the contours of the stronger resonance centered near -25 ppm implies a Mo site with a rather small P_Q , that is spectrally broadened predominantly from a chemical shift distribution. On the other hand, the weaker resonance appears to have a higher degree of quadrupolar broadening and its location in the MQMAS spectrum yields $\delta_{\rm iso}$ ~ 28 ppm and $P_{\rm O}$ \sim 2.5 MHz from eqs 5 and 6. The corresponding quadrupolar broadening of the line shape for this resonance at this magnetic field would be on the order of ~25 ppm, which justifies the simulation of the asymmetric line shapes of the ⁹⁵Mo MAS NMR spectrum with two broad Gaussian peaks with $\delta_{iso} = -26$ \pm 2 and 28 \pm 2 ppm and full-width-at-half-maximum (fwhm) of 60 ppm, as shown in Figure 9b. This simulation yields a ratio of ~85:15 for the two Mo sites with $\delta_{iso} = -26 \pm 2$ and 28 \pm 2 ppm, respectively, in the Mo-60 glass. The 95 Mo MAS NMR spectrum of the Mo-62.5 glass can similarly be simulated



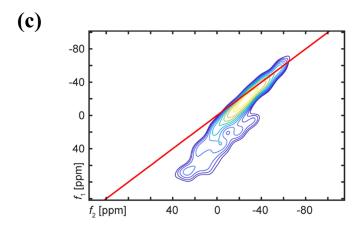


Figure 9. (a) Overlay plot of ⁹⁵Mo MAS NMR spectra of Mo-60 (solid line) and Mo-62.5 (dot-dashed line) glasses. (b) Experimental (black solid line) and simulated (red dashed lines) ⁹⁵Mo MAS NMR spectra of Mo-60 (left) and Mo-62.5 (right) glasses. Individual simulation components corresponding to MoO₄ and MoO₆ sites are shown as blue and green shaded peaks, respectively. (c) Contour plot of ⁹⁵Mo MQMAS/QCPMG NMR spectra of Mo-62.5 glass displaying a strong and a weak resonance corresponding, respectively, to MoO₄ and MoO₆ sites.

with two Gaussian peaks with the same fwhm and similar $\delta_{\rm iso}$ = -27 ± 2 and 23 ± 2 ppm with a relative ratio of 70:30 (Figure 9b). The presence of more than one type of Mo-O coordination environment in these glasses was also reported in a recent study based on Raman spectroscopy, which suggested a coexistence of MoO₄ tetrahedra and MoO₆ octahedra in the structure of these alkali molybdate glasses. More specifically, the structure of glasses with \leq 55% MO₃ was found in this study to be dominated by isolated $[MO_4]^{2-}$ tetrahedra, while the Raman spectra of glasses with higher MO₃ content such as the ones studied here were interpreted to indicate the formation of chain-like moieties of corner- and edge- sharing MO₆ octahedra and MO₄ tetrahedra. Considering the lack of a clear chemical shift systematics, a structural assignment of the two 95Mo resonances in Figure 9 cannot be made solely on the basis of their differences in $P_{\rm O}$. However, the δ_{iso} and small P_{O} of the stronger resonance are indeed similar to those that are characteristic of the MoO₄ tetrahedral sites in a variety of alkali ortho- and pyro-molybdates, as shown in the present study (Table 2). On the other hand, the δ_{iso} of the weaker resonance (23-28 ppm) falls within the range (-11 to 124 ppm) observed in this study for the MoO₆ and $Mo_{5+\delta}$ sites sites in Na- and K-pyromolybdate and in the Lipolymolybdate compounds. Therefore, we tentatively assign the stronger resonance to the MoO₄ tetrahedral environment and the weaker resonance to the high-coordinated MoO₆/ $Mo_{5+\delta}$ environment in the structure of alkali molybdate glasses. Simulations presented in Figure 9b indicate that the relative

fraction of the latter environment increases from ~ 15 to $\sim 30\%$ upon increasing the Mo content in these glasses. This trend is consistent with the results of a previous Raman spectroscopic study, which indicated the tetrahedral environment to be more abundant than the octahedral environment in these glasses and the relative fraction of the latter increased with Mo content.

3.5. Implications for the Local Structure around Mo in Nuclear Waste Glasses. Molybdenum serves as an important constituent in various structural components in nuclear reactors where its oxidation results in significant concentration of MoO₃ in the high-level liquid waste (HLW) produced by the reprocessing of the spent fuel.⁵⁰ The HLW is typically immobilized and stored via vitrification in a borosilicate glass matrix. However, the low solubility of Mo in the borosilicate glass network may result in phase separation and crystallization of various alkali molybdates in the so-called yellow phase that can negatively impact the durability of the encapsulated waste form. SO-52 Majority of the applications of ⁹⁵Mo NMR spectroscopy therefore focused on the characterization of these crystalline alkali molybdate phases in typical Na-Ca-borosilicate glass matrices. 51,52 There has, however, been one systematic study in the literature of the Mo environment in borosilicate glasses by Magnin et al. that utilized 95Mo MAS NMR at a magnetic field of 18.8 T.53 These spectra are characterized by broad and unresolved line shapes where the broadening cannot be conclusively attributed to quadrupolar or chemical shift distribution effects. The center

of gravity of the 95Mo MAS NMR line shape was found to systematically shift upfield to more negative chemical shift values on progressive replacement of Na with Ca as the network modifying cation. These spectra were a priori assigned to [MoO₄]⁻² tetrahedral species on the basis of previous a Mo K-edge EXAFS study⁵⁴ and the authors argued that these tetrahedra, when charge balanced by Ca cations resulted in a more negative chemical shift (e.g., -50 ppm) compared to when charge balanced by Na cations (e.g., -5 ppm) in the glass structure.⁵³ At first glance this chemical shift trend is at odds with the fact that the 95 Mo δ_{iso} of CaMoO₄ is \sim 80 ppm higher than that of Na₂MoO₄. However, the ⁹⁵Mo δ_{iso} has been shown to be highly sensitive to the number and type of charge balancing cations that share the oxygen with the Mo-O polyhedra. The low solubility of Mo in the borosilicate glass has been associated in the literature with the formation of isolated [MoO₄]⁻² tetrahedra, which would scavenge the alkali and alkaline-earth cations from the structure of the host glass to phase separate and eventually form the so-called yellow phase. 51,52 However, an increased solubility of Mo on replacement of Si with P in such glasses has been hypothesized to be related to the formation of MoO6 octahedra that can be connected to the PO₄ tetrahedra via Mo-O-P linkages in the glass structure. 52 The results presented in this study indicate that such a mechanism of Mo dissolution may be investigated in future using high-field and high-resolution 95Mo NMR spectroscopy.

4. SUMMARY

High-resolution 95Mo MAS and MQMAS NMR spectra of crystalline anhydrous alkali ortho-, pyro-, and poly-molybdates in combination with DFT-based calculations provide distinct 95 Mo δ_{iso} ranges of -17 to -125 ppm for MoO₄ and of -11 to 124 ppm for MoO₆ coordination environments, respectively. The corresponding C_0 values are found to be ≤ 2.6 MHz for MoO_4 and ≥ 3.3 MHz for MoO_6 sites. These apparently anomalous trends of higher $\delta_{\rm iso}$ and $C_{\rm Q}$ for a high-coordinated site are tentatively linked to the highly distorted geometry of MoO₆ octahedra in corner- and edge-shared configuration in these molybdates that results in an increase in the paramagnetic component of the 95Mo chemical shift. However, detailed DFT based calculations will need to be carried out in future to critically test this hypothesis. The MoO₄ tetrahedral sites display an approximately linear increase in C_0 with σ , while no such correlation is observed for the high-coordinated Mo sites. The 95Mo MAS and MQMAS NMR spectra indicate that tetrahedral and octahedral Mo-O environments coexist in alkali molybdate glasses and the MoO₆:MoO₄ ratio increases with increasing MoO₃ content. The results presented in this study may have important implications in the future applications of high-resolution 95Mo NMR spectroscopy in probing the Mo coordination environment and its dissolution mechanism in nuclear waste glasses.

AUTHOR INFORMATION

Corresponding Author

Sabyasachi Sen — Department of Materials Science and Engineering, University of California at Davis, Davis, California 95616, United States; orcid.org/0000-0002-4504-3632; Email: sbsen@ucdavis.edu

Authors

Ivan Hung − National High Magnetic Field Laboratory, Tallahassee, Florida 32310, United States; ⑤ orcid.org/ 0000-0001-8916-739X

Jacob M. Lovi – Department of Materials Science and Engineering, University of California at Davis, Davis, California 95616, United States

Zhehong Gan — National High Magnetic Field Laboratory, Tallahassee, Florida 32310, United States; orcid.org/ 0000-0002-9855-5113

Complete contact information is available at: https://pubs.acs.org/10.1021/acs.jpcc.5c00614

Notes

The authors declare no competing financial interest.

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