

$\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$: Triple-Anion Glassy Superionic Conductors for High-Performance Solid-State Batteries

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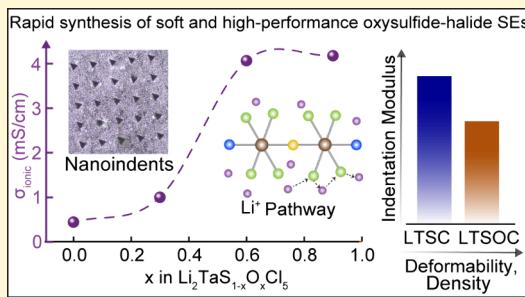
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ABSTRACT: The development of stable, high-performance solid electrolytes is critical for advancing all-solid-state lithium batteries (ASSLBs). We report a series of triple-anion electrolytes, $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ ($0 \leq x \leq 0.9$), synthesized via a rapid 2-hour mechanochemical process. The optimal composition, $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ (LTSC), achieves a room-temperature ionic conductivity of $\sim 4.2 \text{ mS cm}^{-1}$, over 15 times that of $\text{Li}_2\text{TaS}\text{Cl}_5$. XRD confirms its amorphous nature, while ^{67}Li NMR reveals one magnetically equivalent lithium environment due to fast ion-exchange dynamics. Raman spectroscopy shows extensive anion mixing within Ta-centered octahedra, where O^{2-} and S^{2-} occupy axial positions, linking $\text{Ta}-\text{O}-\text{S}-\text{Cl}$ units, while Li^+ primarily migrates along equatorial Cl^- -lined pathways. Nanoindentation reveals a reduced elastic modulus with oxygen incorporation. When employed with commercial NMC811, LTSC delivers an initial capacity of 187 mAh g^{-1} at 0.1C , with $\sim 81.8\%$ retention after 100 cycles and Coulombic efficiency exceeding 99%. These results demonstrate the promise of amorphous, mixed-anion solid electrolytes for scalable, high-performance ASSLBs.



Meeting global decarbonization goals requires energy storage systems that are high-performing, safe, scalable, and sustainable. Lithium-ion batteries are the current standard for portable electronics and are central in the transition toward electric vehicles and renewable energy integration. However, reliance on flammable liquid electrolytes and critical materials like cobalt and nickel raises concerns around safety, cost, and long-term resource availability.^{1–3} All-solid-state lithium batteries (ASSLBs), which replace the flammable liquid with an inorganic solid electrolyte (SE), offer a promising alternative with potential for enhanced energy and power density, thermal stability, and improved safety profiles.^{4–7}

A key challenge in advancing ASSLBs is developing solid electrolytes (SEs) that combine high lithium-ion conductivity, broad electrochemical stability, mechanical deformability, and scalable synthesis.^{8,9} Oxide-based SEs, such as $\text{Li}_2\text{La}_3\text{Zr}_2\text{O}_{12}$ and perovskite-type $\text{Li}_{0.33}\text{La}_{0.55}\text{TiO}_3$, exhibit excellent oxidative stability but typically require high-temperature sintering and suffer from interfacial impedance and mechanical rigidity.^{10–12} Thiophosphate-based conductors, including $\text{Li}_{10}\text{GeP}_2\text{S}_{12}$ (LGPS) and argyrodite-type $\text{Li}_6\text{PS}_5\text{Cl}$, offer superionic conductivities exceeding 10 mS cm^{-1} and low interfacial resistance due to their ductility but are limited by poor

electrochemical and chemical stability.^{13–15} Halide electrolytes such as Li_3InCl_6 , Li_3HoCl_6 , Li_3ZrCl_6 , and Li_3YCl_6 have shown promise with high oxidation stability (up to 4.5 V vs Li^+/Li) and good processability, yet they often exhibit lower conductivities compared to thiophosphates and susceptibility to instability against Li metal anodes.^{16–20}

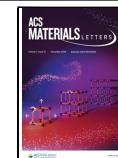
To reconcile the trade-offs among these distinct electrolyte families, mixed-anion strategies have emerged as a promising approach to enhance ionic transport and stability. By combining multiple anionic species with different polarizabilities and bonding preferences, mixed-anion systems introduce local structural disorder, soften the lattice, and enable low-barrier Li-ion conduction pathways.^{21–23} Recent work has demonstrated the efficacy of this approach in various binary systems. In LiNbOCl_4 , corner-sharing oxygen linkages form a flexible framework with low-energy diffusion path-

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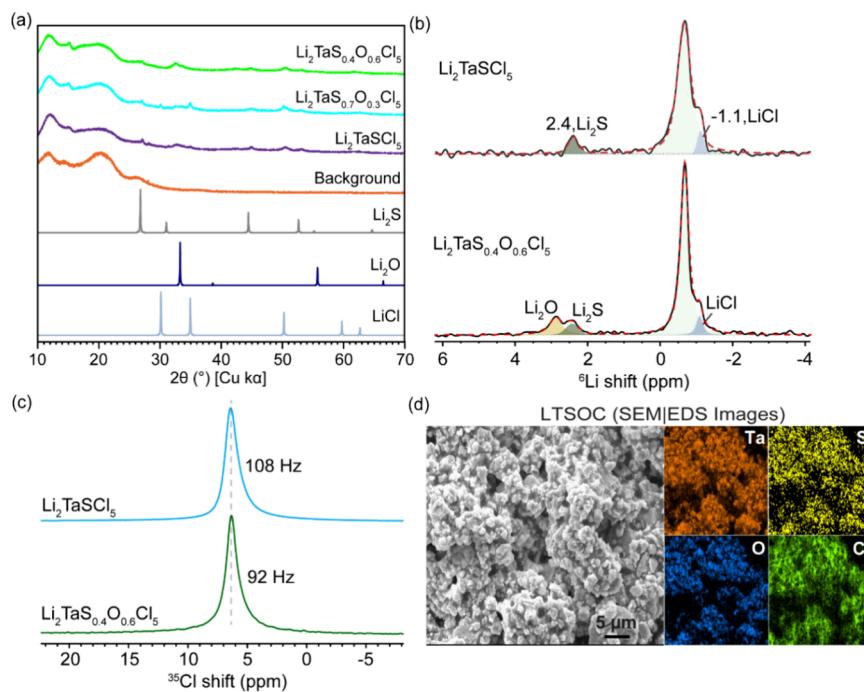


Figure 1. Average structure, local environments, and morphology of the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ series. (a) Powder XRD patterns. (b) ⁶Li MAS NMR for $\text{Li}_2\text{TaS}\text{Cl}_5$ and $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$. (c) ³⁵Cl MAS NMR spectra showing the halide local environment. (d) Scanning electron microscopy (SEM) images and energy dispersive spectroscopy (EDS) mapping of as-milled $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ sample.

ways.^{24–26} LiTaOCl_4 , a highly amorphous material, exhibits superionic properties, though its transport mechanism remains poorly understood.^{27–29} Achieving high ionic conductivity in these SEs often requires prolonged high-energy ball milling (10–40 h) or high-temperature sintering to induce disorder or to improve crystallinity.^{30,31} These energy- and time-intensive processes hinder scalability. Thus, fast, low-energy routes to produce high-conductivity amorphous halide electrolytes are urgently needed to advance practical ASSLBs.

Here, we investigate a sulfide-halide compound, $\text{Li}_2\text{TaS}\text{Cl}_5$, and introduce a triple mixed-anion system, $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ ($0 \leq x \leq 0.9$), designed to combine the key advantages of sulfide, oxide, and halide solid electrolytes within a single disordered framework. The series is synthesized via a rapid, scalable 2 h mechanochemical route, with no evolution of HCl gas commonly reported for analogous oxyhalide systems.^{27,30} The optimized composition, $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ (LTSOC), exhibits a high room-temperature lithium-ion conductivity of $\sim 4.2 \text{ mS cm}^{-1}$ —an order of magnitude improvement over the pristine compound, $\text{Li}_2\text{TaS}\text{Cl}_5$ (LTSC). X-ray diffraction confirms its highly disordered nature, while solid-state ⁶Li and ³⁵Cl magic-angle-spinning (MAS) NMR reveal a symmetric Cl^- framework and isotropic Li^+ environment with rapid ion exchange. Raman data show a possible clustering of the S^{2-} , Cl^- , and O^{2-} anions around Ta^{5+} centers. This combination of S^{2-} , Cl^- , and O^{2-} anions disrupts long-range ordering and facilitates percolating Li^+ conduction pathways through dynamically reconfigurable local polyhedral environments. Mechanical testing shows that LTSOC deforms more readily and is elastically softer than LTSC. In an all-solid-state $\text{Li}|\text{LPSC|LTSOC|NMC811}$ cell, LTSOC is utilized as the catholyte, enabling an initial discharge capacity of $\sim 187 \text{ mAh g}^{-1}$ at 0.1 C (25 °C) and 81.8% capacity retention over 100 cycles with a Coulombic efficiency >99%, demonstrating stable and high performance.

To understand the structural properties of the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ samples prepared via rapid 2-h mechanochemical synthesis, XRD and solid-state MAS NMR were conducted. The XRD patterns of the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ series (Figure 1a) display broad diffraction features, indicative of the lack of a well-defined long-range periodic structure. Despite the broad profile, distinct weak features were detected at approximately 15.1°, 32.6°, and 42.5°. Compared to reference patterns of potential secondary phases, including Li_2S , Li_2O , and LiCl , the Li_2S reflections appear slightly shifted to higher angles by less than 0.5°, particularly in $\text{Li}_2\text{TaS}\text{Cl}_5$ and $\text{Li}_2\text{TaS}_{0.3}\text{O}_{0.7}\text{Cl}_5$. Peaks corresponding to Li_2O were observed to be more prominent in $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$.

Additionally, weak reflections near 30° and 50° from LiCl were consistently observed across all samples, likely arising from the intermediate formation of LiCl during the mechanochemical reaction, as similarly reported in related oxyhalide systems.^{27,30} The diffraction pattern of the sintered LTSOC sample is shown in Figure S1. While heating produced sharpened peaks with increased crystallinity, it significantly reduced electrochemical performance. The ⁶Li MAS NMR spectra in Figure 1b provide insight into the local structural environments of Li^+ ions in LTSC and LTSOC. Both samples exhibit near-identical isotropic peaks, -0.67 ppm for LTSC and -0.68 ppm for LTSOC.

Notably, the isotropic ⁶Li peak resonates close to the LiCl chemical shift at -1.1 ppm , which could stem from possible dominant $\text{Li}^+ \text{Cl}^-$ interactions. The line width of LTSOC (20 Hz) is over 1.5 times narrower than that of LTSC, 34 Hz. The narrower peak observed in LTSOC suggests (a) a more homogeneous lithium local environment or (b) motional averaging due to faster Li^+ mobility that reduces anisotropic and isotropic broadening.^{32–34} The broader line width in LTSC suggests restricted Li^+ ion dynamics. Consistent with

the PXRD data, ^6Li NMR revealed the presence of minor LiCl and Li_2S impurities in LTSC, and additional Li_2O in LTSOC.

To probe the local Cl^- environment, ^{35}Cl MAS NMR was performed on LTSC and LTSOC. As shown in Figure 1c, both samples exhibit a single sharp ^{35}Cl resonance centered at ~ 6.4 ppm, indicating that Cl^- ions experience a comparable local coordination environment. The retention of a similar chemical shift between the two materials confirms that Cl^- remains structurally integrated into the mixed-anion framework, likely within the Ta-centered octahedra, and not phase-separated or forming LiCl -like domains, except for the minor LiCl impurities observed. The ^{35}Cl peak in LTSOC is narrower (92 Hz) compared to that in LTSC (108 Hz). Enhanced Li^+ or anion dynamics in LTSOC could contribute to motional averaging of the quadrupolar interaction, yielding narrowed lines. These results complement the ^6Li NMR data discussed previously and reinforce the notion that oxygen substitution modulates the local anion framework, thereby facilitating dynamic environments conducive to fast ion transport.

SEM and EDS analyses were performed on ball-milled LTSOC (Figure 1d) and LTSC (Figure S2) powders to assess their morphology and elemental distribution. SEM images show that LTSOC particles exhibit more aggregation than those of LTSC, indicative of enhanced interparticle contact. EDS maps confirm a uniform distribution of Ta, S, O, and Cl in LTSOC, indicating successful integration of all anions into the amorphous matrix. Similarly, LTSC displays homogeneous mapping of Ta, S, and Cl.

Figure 2a presents the Raman spectra of several tantalum-based compounds and the synthesized LTSOC: TaCl_5 ,

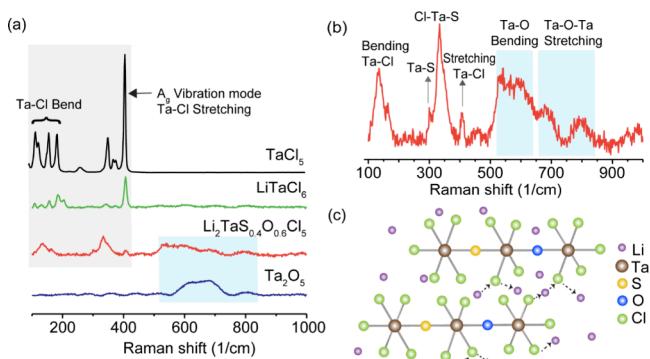


Figure 2. Raman spectra and short-range structural analysis of $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$. (a) Comparison with reference compounds. (b) Raman spectrum of $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ with key vibrational modes assigned. (c) Polymeric-type $[\text{TaS}_x\text{O}_{1-x}\text{Cl}_5]^{2-}$ anionic chains with possible Li^+-Cl^- interactions along the conduction pathways indicated.

LiTaCl_6 , Ta_2O_5 , and $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$. The spectrum of TaCl_5 exhibits sharp peaks, typical of molecular species. The bands between 100 and 200 cm⁻¹ originate from Ta-Cl bending modes; specifically, those below 150 cm⁻¹ arise from out-of-plane bending due to Cl wagging or twisting, and those between 150 and 200 cm⁻¹ from in-plane Cl bending. The peaks between 200 and 400 cm⁻¹ are attributed to Ta-Cl stretching, including the intense, sharp peak at ~ 405 cm⁻¹ from axial Ta-Cl symmetric stretch and the weaker peaks at slightly lower wavenumbers from equatorial Ta-Cl symmetric stretch. LiTaCl_6 shows similar peaks but slightly broader with lower intensity, typical of extended 3D lattices with long-range

interactions, lattice vibrations, and possible site disorder. In contrast, the spectrum of LTSOC (red curve) shows significantly broadened peaks and notable shifts, indicating substantial structural modification due to partial substitution of S^{2-} and O^{2-} for Cl^- within the $[\text{TaS}_x\text{O}_{1-x}\text{Cl}_5]^{2-}$ octahedra. Figure 2b provides a more detailed view of the Raman features for LTSOC. Low-wavenumber peaks at approximately 134 cm⁻¹ and 165 cm⁻¹ are assigned to Ta-Cl and Ta-S bending vibrations.³⁵ The low resolution of these bending modes in the Raman spectrum indicates heterogeneity in the local structure and bonding environments. A broad, prominent peak at ~ 330 cm⁻¹, absent in TaCl_5 and LiTaCl_6 , is attributed to a Cl-Ta-S vibrational mode, potentially arising from a hybrid Cl-Ta-S linkage or structural distortions induced by multianion coordination. This peak is flanked by the Ta-S E_g mode at ~ 300 cm⁻¹,^{36,37} and the octahedral Ta-Cl A_{1g} mode at ~ 405 cm⁻¹. This assignment is consistent with literature reports attributing Ta-S vibrations in the 260–390 cm⁻¹ range for related compounds with Ta-S-Ta linkages.^{38–42} Notably, this 405 cm⁻¹ peak from axial Ta-Cl symmetric stretching—dominant in TaCl_5 —is significantly diminished in LTSOC, indicating that S/O substitution mainly occurs at axial positions, as illustrated in Figure 2b. In the 600–900 cm⁻¹ region (highlighted in blue), the LTSOC spectrum shows broad features associated with oxygen incorporation, specifically Ta-O-Ta stretching and bending modes.⁴³ These features confirm the presence of Ta-O bonds. The broadening across the spectrum reflects increased structural disorder, which is likely to influence ionic transport properties.

The anion arrangement shown in Figure 2c aligns with the basic structural motifs reported for compounds such as LiTaOCl_4 and LiNbOCl_4 , in which the M^{5+} octahedra $[\text{MO}_2\text{Cl}_4]$ are coordinated equatorially by four Cl^- ions and axially by two O^{2-} atoms. These axial oxygens bridge neighboring Ta/Nb centers along the c-direction, forming 1D flexible anionic chains.^{24,30,44} Similarly, the placement of S^{2-} into axial O^{2-} positions is consistent with computational studies on the $[\text{TaSOCl}_4]$ octahedron, where sulfur preferentially occupies axial sites due to its lower substitution free energy (ΔG)—over 26 times lower for S replacing O compared to Cl.²⁷

Figure 2c schematically illustrates short-range ordered motifs and possible Li^+ diffusion pathways. The Ta-centered octahedra are connected into polymeric chains, forming a symmetrical, yet locally disordered anion environment composed of S, O, and Cl. The dashed lines indicate potential Li^+ hopping pathways. Li^+ diffusion in the amorphous LTSOC phase likely proceeds via localized hops between coordination sites in the distorted $[\text{Ta}-\text{Cl}-\text{O}-\text{S}]$ clusters. These clusters generate a dynamic energy landscape, with Li^+ ions interacting weakly, particularly with Cl^- , as supported by ^6Li NMR, which shows proximity to halide-rich environments. This mechanism is reminiscent of the “monkey-bar” transport model proposed for LiTaCl_6 by Lei et al.,⁴⁵ where Li^+ ions transiently coordinate to anionic sites before hopping to the next site. Such sequential hopping creates a percolating network of migration pathways, enabling long-range ionic conduction even in the absence of long-range crystalline order. The bridging O and S links $[\text{TaS}_x\text{O}_{1-x}\text{Cl}_5]$ octahedra, which may facilitate Li^+ hopping between different $[\text{TaS}_x\text{O}_{1-x}\text{Cl}_5]$ octahedron units and promote long-range Li^+ diffusion with high transference numbers.

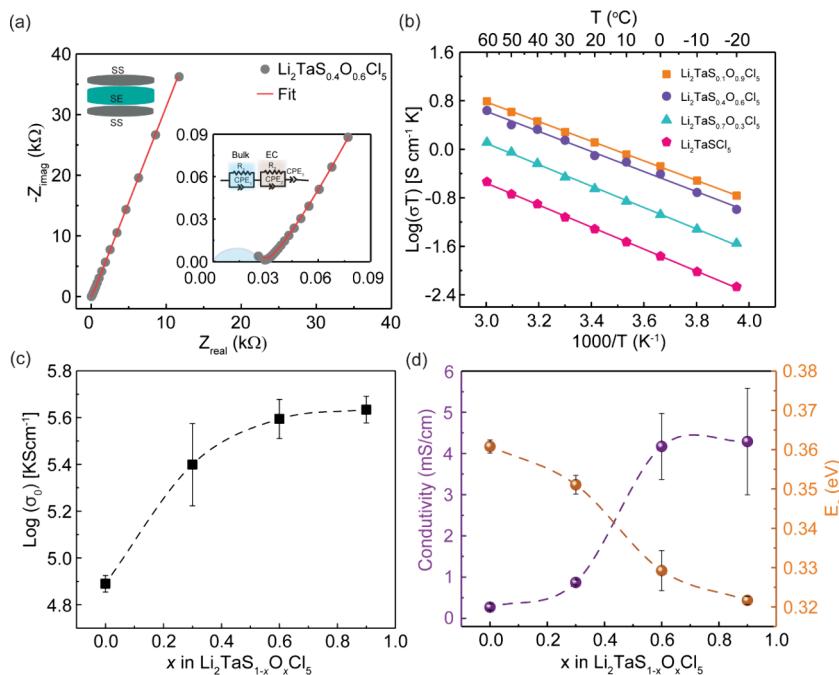


Figure 3. Electrochemical properties of $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$. (a) Nyquist plots and equivalent circuit fit of $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ at 25 °C. (b) Arrhenius plots of the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ series measured from -20 to 60 °C. (c) Arrhenius prefactor (σ_0) distribution in the LTSOC series. (d) Comparison of the ionic conductivities and activation energies trend in the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ series. Data points with error bars represent means and standard deviations ($x \pm 1\sigma$) from measurements conducted on three or more samples.

The room-temperature ^7Li T_1 spin–lattice relaxation times for the LTSOC and LTSC are approximately 0.45 and 0.25 s, respectively (Figure S3). The longer T_1 in the oxygen-rich LTSOC suggests slower relaxation. To relate these relaxation behaviors to ion mobility and determine the motional regime, the Bloembergen–Purcell–Pound (BPP) model⁴⁶ was applied to variable-temperature T_1 measurements (Figure S3a, Tables S1 and S2). Both LTSC and LTSOC exhibit trends characteristic of the slow-to-intermediate motion regime, in which higher Li^+ mobility leads to shorter T_1 times before T_1 reaches the minimum.⁴⁶ Thus, shorter T_1 times in the slow motional regime typically correlate with faster Li^+ motion. However, the observed discrepancy between the absolute T_1 values for LTSOC and LTSC and the conductivity trend (discussed in the next section) implies that Li^+ jump frequency alone may not fully explain the superionic behavior observed in LTSOC. Other factors, may contribute to the observed high ionic conductivity.

Additionally, the ^7Li MAS NMR spectra (Figure S3b) show a clear temperature-dependent line width narrowing in both LTSC and LTSOC, indicative of enhanced Li^+ mobility at elevated temperatures. In LTSC, the line width decreases from 648 Hz at 25 °C to 480 Hz at 85 °C, while in LTSOC, it narrows from 460 to 360 Hz. This behavior reflects thermally activated Li^+ motion and motional averaging of residual ^7Li – $^{7\text{Li}}$ dipolar coupling interactions.

The ionic transport properties of LTSOC were evaluated using temperature-dependent electrochemical impedance spectroscopy. At 25 °C, the fitted Nyquist plot (Figure 3a) yields a total ionic conductivity of ~ 4.2 mS cm $^{-1}$. The data were modeled using an equivalent circuit (inset, Figure 3a), consisting of a parallel resistor and constant phase element (R, CPE) to represent the ionic response of the sample, in series with a second CPE that accounts for the ion-blocking

behavior of the stainless-steel electrodes. The extracted Brug capacitance ($\sim 10^{-11}$ F) suggests the contributions from bulk and grain boundaries cannot be clearly distinguished. Even at -20 °C (Figure S4), the impedance response does not resolve separate features for bulk and grain boundary processes. As such, the reported values reflect total ionic conductivity, consistent with trends observed in related halide-based solid electrolytes.^{47,48}

Arrhenius plots for the $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$ series derived from variable-temperature EIS measurements (Figure S5) are presented in Figure 3b, with the corresponding activation energies and room-temperature ionic conductivities, calculated using eqs S1 and S2, summarized in Figure 3d. The unsubstituted LTSC exhibits the highest activation energy (~ 0.36 eV) and the lowest conductivity (0.27 mS cm $^{-1}$). In contrast, partial substitution of S^{2-} with O^{2-} significantly enhances conductivity, with values reaching 0.87, 4.17, and 4.29 mS cm $^{-1}$ for $x = 0.3$, 0.6, and 0.9, respectively. This increase is accompanied by a corresponding decrease in activation energy to 0.35, 0.33, and 0.32 eV, respectively. Notably, no significant difference in conductivity is observed between the $x = 0.6$ and $x = 0.9$ compositions. Oxygen incorporation likely stabilizes a superionic phase, as $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ exhibits a more than 15-fold improvement in conductivity compared to LTSC. With a total processing time under 2.5 h, this mechanochemically synthesized LTSOC material compares favorably with most reported lithium halide fast-ion conductors synthesized through lengthy, multistep, or high-temperature processes.

In Figure 3c, we observe a systematic increase in $\text{Log}(\sigma_0)$ with increasing oxygen content (x) in $\text{Li}_2\text{TaS}_{1-x}\text{O}_x\text{Cl}_5$, particularly from $x = 0$ to 0.6, before reaching a plateau. This rise in the prefactor σ_0 , which includes an entropic component through the term $\exp(\Delta S_m/k_B)$,⁵⁰ reflects a higher entropy of migration, S_m , driven by increased structural and

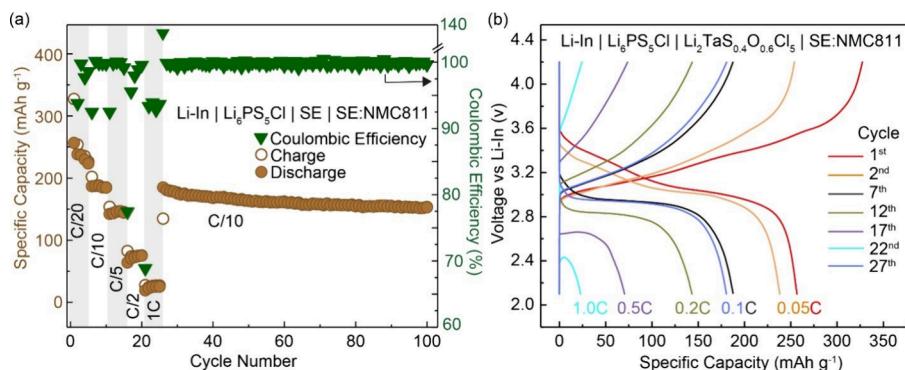


Figure 4. All-solid-state battery cycling performance of $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ in a cell with the configuration of $\text{Li}-\text{In}|\text{Li}_6\text{PS}_5\text{Cl}|\text{SE}|\text{SE}:\text{NMC811}$, where SE is $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$. (a) Rate capability plots. (b) Voltage-capacity profiles showing the charge–discharge curves at 0.05C , 0.1C , 0.2C , 0.5C , and 1C . $\text{C} = 275.5 \text{ mAh g}^{-1}$, referenced to the theoretical capacity of NMC811.⁴⁹

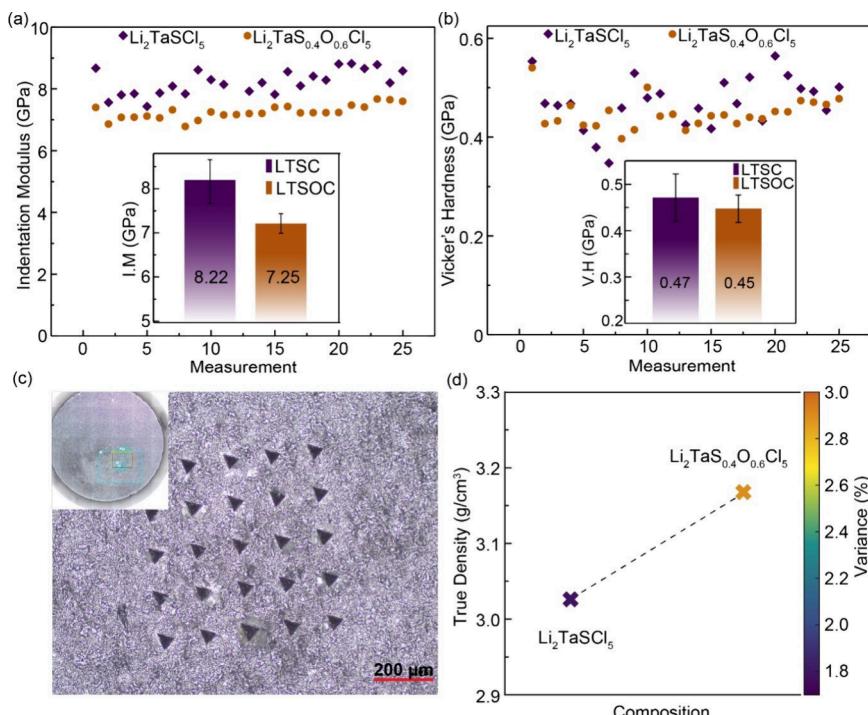


Figure 5. Mechanical properties of $\text{Li}_2\text{TaS}\text{Cl}_5$ and $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$. (a) Indentation modulus measured using the Advanced Dynamic Modulus and Hardness (ADEH) protocol. (b) Vickers hardness values. The error bars represent the standard deviation ($x \pm 1\sigma$) of multiple measurements. (c) Optical surface profilometry images ($10\times$ magnification) of a 5×5 grid of indentation points (depth $4\text{--}5 \mu\text{m}$) on a $\text{Li}_2\text{TaS}_{0.4}\text{O}_{0.6}\text{Cl}_5$ pellet. Inset: Pellet image at $2.5\times$ magnification, with marked areas showing the indented sections. (d) True density measurements by gas pycnometry.

dynamic disorder introduced by oxygen substitution. The resulting entropic enhancement complements the reduction in activation energy and contributes to the overall improvement in ionic conductivity. Figure S6 shows the current–time response at various voltages, with the inset displaying the corresponding I–V plot. From this, the electronic conductivity of LTSOC was determined to be $\sim 3 \times 10^{-9} \text{ S cm}^{-1}$, confirming its negligible contribution to the total conductivity of 4.2 mS cm^{-1} . These results underscore the potential of triple-anion design for enabling high-performance solid electrolytes compatible with large-scale solid-state battery manufacturing.

ASSLBs assembled using the commercially available high-voltage NMC811 cathode and a multilayered configuration NMC811:LTSOC:LTSOC| $\text{Li}_6\text{PS}_5\text{Cl}|\text{Li}-\text{In}$ were cycled be-

tween 2.1 and 4.2 V at 23°C . To evaluate rate capability, initial cycles were performed at 0.05C , 0.1C , 0.2C , 0.5C , and 1C ($\text{C} = 275.5 \text{ mAh g}^{-1}$), delivering discharge capacities of 256.7 , 187.0 , 143.4 , 70.4 , and 23.0 mAh g^{-1} , respectively (Figure 4a). The decline in capacity with increasing C-rate reflects kinetic limitations, including sluggish Li^+ transport and elevated overpotential. At higher current densities, increased ionic flux and restricted charge transfer induce greater polarization, leading to early voltage cutoffs and underutilization of the cathode.⁴⁹ These effects are also evident in the voltage profiles (Figure 4b), which show modest hysteresis that gradually widens with cycling, likely due to an increase in overpotential at high rates.

Despite these challenges, the cell maintains excellent long-term stability, retaining 152.9 mAh g^{-1} over 100 cycles at

0.1C—approximately 81.8% of its initial capacity. This performance is attributed to (i) the chemical and electrochemical compatibility of LTSOC with NMC811, which minimizes parasitic reactions and promotes a stable cathode-electrolyte interphase (CEI), (ii) the high oxidative stability of LTSOC, preventing electrolyte breakdown at elevated voltages, and (iii) maintenance of good cathode-electrolyte interfacial contact due to favorable mechanical properties of the glassy triple-anion conductor. As a result, a high Coulombic efficiency of >99% is maintained throughout long-term cycling. These results demonstrate LTSOC as a promising high-voltage-compatible solid electrolyte for enabling durable solid-state batteries.

Given the role of mechanical features in sustaining cathode-electrolyte contact, we characterized the mechanical response of the glassy triple-anion conductor. Figure 5a and 5b show the indentation modulus and Vickers hardness of both SEs, measured using a Berkovich nanoindenter. Indentation hardness (converted to Vickers hardness, see *methods* in the Supporting Information) quantifies how resistant a material is to permanent plastic deformation under a sharp indenter, whereas indentation modulus reflects the elastic response of a material, derived from the contact stiffness normalized by contact area.⁵¹ At depths of 4–5 μm , we find that LTSC exhibits a higher indentation modulus (~8.2 GPa) compared to LTSOC (~7.3 GPa), a ~12% decrease upon O²⁻ substitution. A lower indentation modulus for LTSOC means that the material is elastically softer or less stiff under load. This could be beneficial for maintaining good interfacial contact with electrodes and for accommodating volumetric changes during battery cycling.⁵² Hardness values, by contrast, were not significantly different, with LTSC having a slightly higher Vickers hardness of ~0.47 and LTSOC ~0.45 GPa (Figure 5b). The overlap within experimental scatter suggests that oxygen incorporation does not significantly alter the plastic resistance to indentation. Their hardness (0.4–0.5 GPa) is consistent with halide SEs, and far lower than garnet oxides such as Li₇La₃Zr₂O₁₂ (>6 GPa hardness, modulus >150 GPa).^{52–54}

Figure 5c presents optical images of the indented LTSOC pellet surface, confirming reproducible indentation imprints without major crack formation or delamination. The absence of severe cracking indicates that these electrolytes can undergo localized deformation under stress. True density measurements (Figure 5d) confirm that the lower modulus of LTSOC is intrinsic and not due to porosity effects. LTSOC exhibits a higher density (3.17 g cm⁻³) than LTSC (3.03 g cm⁻³), with relatively low variance (2.9%).

In conclusion, we investigated the structure and performance of a series of superionic solid electrolytes, Li₂TaS_{1-x}O_xCl₅ (0 ≤ x ≤ 0.9), synthesized via a 2-h mechanochemical route. The composition, Li₂TaS_{0.4}O_{0.6}Cl₅, achieved a room-temperature ionic conductivity of ~4.2 mS cm⁻¹ with a low activation energy of ~0.33 eV, representing a ~15-fold enhancement compared to Li₂TaS₁Cl₅. XRD confirmed their amorphous nature, while ⁶⁷Li NMR revealed fast-moving Li⁺ ions interacting with Cl⁻. Raman spectra showed broadened and weakened Ta–Cl vibrational modes and new modes associated with Ta–O and Ta–S bonding, with O²⁻ and S²⁻ occupying axial sites and bridging [TaSOCl₄] octahedra, facilitating Li⁺ transport across units. An all-solid-state battery employing Li₂TaS_{0.4}O_{0.6}Cl₅ as the catholyte and NMC811 as the cathode delivered an initial specific capacity of ~187 mAh g⁻¹, retained

~81.8% capacity after 100 cycles at 0.1C, and exhibited >99% Coulombic efficiency. Nanoindentation confirmed that oxygen substitution lowers modulus while preserving hardness, yielding mechanical compliance for stable interfaces. These results highlight the critical role of oxygen incorporation in stabilizing a fast-ion-conducting disordered structure, providing a generalizable design strategy for high-performance, scalable halide-based SEs for next-generation ASSLBs.

■ ASSOCIATED CONTENT

Data Availability Statement

Data supporting the findings of this study are available from the corresponding author upon request.

■ Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsmaterialslett.5c01054>.

Detailed experimental procedures, materials, and methods, and additional data, including selected XRD patterns, SEM images, NMR T₁ relaxation data, EIS data, and DC polarization data ([PDF](#))

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Notes

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